# Dynamics of gases of particles with singular repulsion

#### Djalil CHAFAÏ

Paris-Dauphine / PSL

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# Outline

#### Introduction

Coulomb gases

Concentration of measure

Dynamics for planar case

# High dimensional phenomenon : random matrix spectrum



plot(eig(randn(n, n)+i\*randn(n, n))/sqrt(2\*n)))

## High dimensional phenomenon : random matrix spectrum



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Law of large numbers: orthonormal rows/columns for large n

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 Jean Ginibre (1938 – ) Statistical Ensembles of Complex, Quaternion, and Real Matrices Journal of Mathematical Physics (1965)

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Terence Tao (1975 – ), Sylvia Serfaty (1975 – ), Robert Berman (1976 – ), ...

Non-Hermitian (Ginibre) 2D

$$e^{-n\operatorname{Trace}(MM^*)} = \prod_{j,k=1}^{n} e^{-n|M_{jk}|^2}$$

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- Ginibre is not log-concave (contrary to Dyson)

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# Coulomb kernel in mathematical physics

Coulomb kernel in  $\mathbb{R}^d$ ,  $d \ge 2$ ,

$$x \in \mathbb{R}^d \mapsto g(x) = egin{cases} \log rac{1}{|x|} & ext{if } d=2, \ rac{1}{|x|^{d-2}} & ext{if } d\geq 3. \end{cases}$$

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$$\Delta g = -c_d \, \delta_0 \quad ext{where} \quad c_d = egin{cases} 2\pi & ext{if } d = 2, \ (d-2)|\mathbb{S}^{d-1}| & ext{if } d \geq 3. \end{cases}$$

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For both Dyson and Ginibre: two dimensional repulsion

Coulomb energy of probability measure  $\mu$  on  $\mathbb{R}^d$ :

$$\mathscr{E}(\mu) = \iint g(x-y)\mu(\mathrm{d} x)\mu(\mathrm{d} y) \in \mathbb{R} \cup \{+\infty\}.$$

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Coulomb energy with confining potential (external field)

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Equilibrium probability measure (electrostatics)

$$\mu_* = \operatorname{arg\,inf} \mathscr{E}_V$$

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If V is stronger than g at infinity then μ<sub>\*</sub> is compactly supported
 If V is smooth then μ<sub>\*</sub> has density

$$\frac{\Delta V}{2c_d}$$

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# Examples of equilibrium measures

d	Interaction g	Confinement V	Equilibrium $\mu_*$
1	2	$\infty 1_{interval^c}(x)$	arcsine
1	2	x <sup>2</sup>	semicircle (Dyson)
2	2	$ x ^2$	uniform on a disc (Ginibre)
$\ge$ 3	d	$  x  ^2$	uniform on a ball
$\ge$ 2	d	radial	radial in a ring

Energy of *n* Coulomb charges  $\frac{1}{n}$  at positions  $x_1, \ldots, x_n$  in  $\mathbb{R}^d$ :

$$H(x_1,...,x_n) = \frac{1}{n} \sum_{i=1}^n V(x_i) + \frac{1}{n^2} \sum_{i \neq j} g(x_i - x_j)$$

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 Gibbs measure on (ℝ<sup>d</sup>)<sup>n</sup>:

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 $\blacksquare \ \propto e^{-\beta_n(\langle V, \mu_n \rangle + \langle \Delta^{-1} \mu_n, \mu_n \rangle)} \rightsquigarrow \text{CLT with Gaussian Free Field}$ 

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Laplace method ~> Large Deviation Principle (Gozlan-C.-Zitt)

$$\frac{\log \mathbb{P}\Big(\operatorname{dist}(\mu_n,\mu_*)\geq r\Big)}{n^2} \xrightarrow[n\to\infty]{} -\frac{\beta}{2} \inf_{\operatorname{dist}(\mu,\mu_*)\geq r} \big(\mathscr{E}_V(\mu)-\mathscr{E}_V(\mu_*)\big).$$

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Quantitative estimates? How to relate dist and  $\mathscr{E}_V(\cdot) - \mathscr{E}_V(\mu_*)$ ?

- Concentration of measure

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Coulomb divergence (Large Deviations rate function)

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$$\sqrt{\mathscr{E}(\mu-\nu)}$$

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$$W_{p}(\mu, \nu) = \inf_{\substack{(X,Y)\\X\sim\mu,Y\sim\nu}} \mathbb{E}(|X-Y|^{p})^{1/p}.$$

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#### Theorem (Transport type inequality – C.-Hardy-Maïda)

$$W_1(\mu, v)^2 \leq C_D \mathscr{E}(\mu - v).$$

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$$D \subset \mathbb{R}^d$$
 compact, supp $(\mu + v) \subset D$ ,  $\mathscr{E}(\mu) < \infty$  and  $\mathscr{E}(v) < \infty$ ,

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Extends free transport inequality to any d

#### Idea of proof of Coulomb transport inequality

Potential: if  $U^{\mu}(x) = g * \mu(x)$  then  $\Delta U^{\mu}(x) = -c_{d} \mu$ 

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Electric field:  $\nabla U^{\mu}(x)$ . "Carré du champ":  $|\nabla U^{\mu}|^2$ 

# Idea of proof of Coulomb transport inequality Potential: if $U^{\mu}(x) = g * \mu(x)$ then $\Delta U^{\mu}(x) = -c_d \mu$

Integration by parts & Schwarz's inequality in  $\mathbb{R}^d$  and  $\mathrm{L}^2$ 

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$$\int |\nabla U^{\mu-\nu}(x)|^2 \mathrm{d}x = -\int U^{\mu-\nu}(x) \Delta U^{\mu-\nu}(x) \mathrm{d}x$$
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Finally  $\mathrm{W}_1(\mu, v)^2 \leq |\mathsf{D}| c_{\mathsf{d}} \mathscr{E}(\mu - v).$ 

## Coulomb transport inequality for equilibrium measures

Corollary (Transport type inequality – C.-Hardy-Maïda)

$$\mathrm{d}_{\mathrm{BL}}(\mu,\mu_*)^2 \leq C_{\mathrm{BL}}\Big(\mathscr{E}_V(\mu) - \mathscr{E}_V(\mu_*)\Big).$$

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Growth condition is optimal for W<sub>1</sub>

Theorem (Concentration inequality – C.-Hardy-Maïda)

If V has reasonable growth then for every  $\beta$ , n, r

$$\mathbb{P}\Big(\mathrm{d}_{\mathrm{BL}}(\mu_n,\mu_*)\geq r\Big)\leq \mathrm{e}^{-aeta\,n^2r^2}$$

Moreover if V has at least quadratic growth then  $W_1$  instead of  $d_{BL}$ .

Optimal order in n

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Implies speed of convergence:

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See also Rougerie & Serfaty

Starting point

$$\mathbb{P}(\mathbf{W}_1(\mu_n,\mu_*)\geq r)=\frac{1}{Z}\int_{\mathbf{W}_1(\mu_n,\mu_*)\geq r}e^{-\frac{\beta}{2}n^2\mathscr{E}_V^{\neq}(\mu_n)}\mathrm{d}x.$$

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Regularization: g superharmonic,  $\mu_n^{(arepsilon)}=\mu_n*\lambda_arepsilon,$ 

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Coulomb transport 
$$-\mathscr{E}_V(\mu_n^{(\varepsilon)}) + \mathscr{E}_V(\mu_*) \leq -\frac{1}{c} W_1^2(\mu_n^{(\varepsilon)}, \mu_*).$$

#### Corollary (Ginibre Random Matrices)

If *M* is  $n \times n$  with iid Gaussian entries of variance  $\frac{1}{n}$  in  $\mathbb{C}$ 

## Eigenvalues of $M \propto \exp(-n\sum_{i=1}^n |x_i|^2) \prod_{i < j} |x_i - x_j|^2$

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Bernoulli ±1 random matrices (universality)  $\mu_n \rightarrow \mu_{\bullet}$  (Tao-Vu 2010) but  $\mathbb{P}(W_1(\mu_n, \mu_{\bullet}) \ge r)$  is not known

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## Outline

Introduction

Coulomb gases

Concentration of measure

Dynamics for planar case
Ginibre energy  $H(x_1, ..., x_n) = \frac{1}{n} \sum_{i=1}^n |x_i|^2 + \frac{1}{n^2} \sum_{i \neq j} \log \frac{1}{|x_i - x_i|}$ 

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Mean-field interacting particle system  $X_t^n = (X_t^{i,n})_{1 \le i \le n}$ 

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  - Optimal Poincaré and log-Sobolev constants (C.-Lehec)

# Second moment dynamics

#### Theorem (Second moment dynamics - Bolley-C.-Fontbona)

 $(R_t)_{t\geq 0} = (\frac{1}{n} ||X_t||^2)_{t\geq 0}$  is an ergodic Cox–Ingersoll–Ross process:

$$\mathrm{d}R_t = \sqrt{\frac{8\alpha_n}{n\beta_n}R_t}\,\mathrm{d}B_t + 4\frac{\alpha_n}{n}\left[\frac{n}{\beta_n} + \frac{n-1}{2n} - R_t\right]}\,\mathrm{d}t.$$

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In particular, with  $\gamma_n = \text{Gamma}(n + \frac{n-1}{2n}\beta_n, \beta_n)$ , for any  $t \ge 0$ 

$$W_1(\operatorname{Law}(\mathcal{R}_t),\gamma_n) \leq e^{-4\frac{\alpha_n}{n}t} W_1(\operatorname{Law}(\mathcal{R}_0),\gamma_n).$$

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Furthermore for any  $x \in D$  and  $t \ge 0$ , we have

$$\mathbb{E}(R_t \mid R_0 = r) = r \mathrm{e}^{-\frac{4\alpha_n}{n}t} + \left(\frac{1}{2} + \frac{n}{\beta_n} - \frac{1}{2n}\right) \left(1 - \mathrm{e}^{-\frac{4\alpha_n}{n}t}\right).$$

Eigenvectors:  $\sum_{i=1}^{N} \Re(z_i)$ ,  $\sum_{i=1}^{N} \Im(z_i)$ ,  $\sum_{i=1}^{N} |z_i|^2 + c_N$ .

#### Theorem? (MKV Mean-field limit – Bolley-C.-Fontbona)

If 
$$\sigma = \lim_{n \to \infty} \frac{lpha_n}{eta_n} \in [0,\infty)$$
 then  $\lim_{n \to \infty} \mu_t^n = \mu_t$  with

$$\partial_t \mu_t = \sigma \Delta \mu_t + \nabla \cdot ((\nabla V + \nabla g * \mu_t) \mu_t).$$

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Some sort of law of large numbers

#### Theorem? (MKV Mean-field limit – Bolley-C.-Fontbona)

If 
$$\sigma = \lim_{n \to \infty} \frac{\alpha_n}{\beta_n} \in [0,\infty)$$
 then  $\lim_{n \to \infty} \mu_t^n = \mu_t$  with

$$\partial_t \mu_t = \sigma \Delta \mu_t + \nabla \cdot ((\nabla V + \nabla g * \mu_t) \mu_t).$$

$$dX_t^{i,n} = \sqrt{2\frac{\alpha_n}{\beta_n}} dB_t^{i,n} - 2\frac{\alpha_n}{n} X_t^{i,n} dt - 2\frac{\alpha_n}{n} \sum_{j \neq i} \frac{X_t^{i,n} - X_t^{j,n}}{|X_t^{i,n} - X_t^{i,n}|^2} dt.$$

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Regimes: 
$$(\alpha_n, \beta_n) = (n, n^2)$$
 and  $(\alpha_n, \beta_n) = (n, n)$ 

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Carrillo–McCann–Villani, Fournier–Hauray–Mischler, Serfaty–Duerinckx, ...

Overdamped Langevin dynamics

$$\mathrm{d}X_t = -lpha 
abla \mathcal{H}(X_t) \mathrm{d}t + \sqrt{rac{lpha}{eta}} \mathrm{d}B_t$$

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 $\rightarrow$  Metropolis Adjusted Langevin Algorithm (MALA)

Underdamped Langevin dynamics (adding momentum/inertia)

$$\begin{cases} \mathrm{d}X_t = \nabla U(Y_t)\mathrm{d}t \\ \mathrm{d}Y_t = -\nabla H(X_t) - \gamma \nabla U(Y_t)\mathrm{d}t + \sqrt{\frac{\gamma}{\beta}}\mathrm{d}B_t. \end{cases}$$

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ightarrow Geometric ergodicity via Lyapunov (Lu–Mattingly)

## Eight trajectories for a Dyson Ornstein-Uhlenbeck HMC











## Universal Gumbel fluctuation for edge of beta Ginibre?

