

# Fermions in low dimensions and non-Hermitian random matrices

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#### Random Matrices and Related Topics in Jeju May 6 2024

with Sungsoo Byun, Markus Ebke and Grégory Schehr [arXiv:2206.08815 = J Stat Phys '23, arXiv2308.05519 = J Phys A '23]

#### Outline

#### Part I: Map to Random Matrix Theory (RMT)

- d-dimensional harmonic trap: GUE & beyond
- 2*d* with magnetic field / rotating trap: complex **Ginibre** ensemble (GinUE)

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#### Part II: Universality and holography in 2d

- variance of number of Fermions  $\mathcal{N}_A$  in A $\boxed{\operatorname{Var}\mathcal{N}_A \sim \partial A}$  for a large class of non-Hermitian RMT

- entropy  $\mathcal{S}_q(\textit{N},\textit{A}) \propto \text{Var}\mathcal{N}_{\textit{A}}$  in bulk scaling

Summary and open questions

- ▶ *N* Fermions at  $\vec{r}_j \in \mathbb{R}^d$ , momenta  $\vec{p}_j = -i\vec{\nabla}_j, j = 1, ..., N$
- ► Harmonic oscillators in *d*-dimensions:

$$\mathcal{H} = \sum_{j=1}^{N} \mathcal{H}_{j}(\vec{p}_{j}, \vec{r}_{j}), \quad \mathcal{H}_{j}(\vec{p}_{j}, \vec{r}_{j}) = \frac{1}{2}(\vec{p}_{j}^{2} + \vec{r}_{j}^{2}) \text{ single particle,}$$
  
non-interacting

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• **example 1d:**  $\mathcal{H}_j(p, x)u_n(x) = E_n u_n(x)$ Hermite eigenfunctions  $u_n(x) \sim H_n(x)e^{-\frac{1}{2}x^2}$ , n = 0, 1, ...

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• example 1d: 
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N-particle Fermionic ground state: Slater determinant

$$\Psi_0(x_1,\ldots,x_N) \sim \det \left[ H_{j-1}(x_i) e^{-x_i^2/2} \right]_{i,j=1}^N \sim e^{-\frac{1}{2}\sum_{i=1}^N x_i^2} \Delta_N(\{x\})$$

► Vandermonde determinant 
$$\Delta_N(\{x\}) = \det \left[x_i^{j-1}\right]_{i,j=1}^N$$

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■ amplitude |Ψ<sub>0</sub>|<sup>2</sup> ~ |Δ<sub>N</sub>({x})|<sup>2</sup>e<sup>-∑<sup>N</sup><sub>i=1</sub>x<sup>2</sup><sub>i</sub> = GUE joint eigenvalue density, determinantal point process</sup>

•  $\mathcal{H}_j(\vec{p}, \vec{r}) = \frac{1}{2} \sum_{l=1}^d (p_l^2 + r_l^2)$  single particle in d dimensions with eigenfunction:

$$\Psi_{\{n\}}(\vec{r}) = \prod_{l=1}^{d} u_{n_l}(r_l), \, \vec{r} = (r_1, \ldots, r_N)$$

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determinental point process in d-dimensions

 $\begin{aligned} |\Psi_0(\vec{r}_1, \dots, \vec{r}_N)|^2 &\sim \det \left[\mathcal{K}_N^{\text{Fermi}}(\vec{r}_1, \vec{r}_2)\right]_{i,j=1}^N \text{ joint density} \\ R_k(\vec{r}_1, \dots, \vec{r}_k) &= \det \left[\mathcal{K}_N^{\text{Fermi}}(\vec{r}_i, \vec{r}_j)\right]_{i,j=1}^k k \text{ -point function} \end{aligned}$ 

$$\mathcal{K}_{N}^{\text{Fermi}}(\vec{r}_{1},\vec{r}_{2}) = \sum_{0 \le j_{1} + \ldots + j_{d} \le n-1} \Psi_{\{j_{l}\}}(\vec{r}_{1})\Psi_{\{j_{l}\}}(\vec{r}_{2})$$
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does not correspond to eigenvalues of RMT for d > 1, but same techniques apply





- ► global density: 1*d* semi-circle (left) vs. 2*d* cup (right)  $\begin{bmatrix}
  \frac{1}{N^{d/2}}\mathcal{K}_{N}^{\text{Fermi}}(\sqrt{N}\vec{r},\sqrt{N}\vec{r}) \sim (2-|\vec{r}|^{2})^{d/2} \\
  \equiv \rho(|\vec{r}|)$
- local correlations: bulk kernel at  $\vec{r}_0$

$$\frac{1}{\rho(\vec{r}_0)N}\mathcal{K}_N^{\text{Fermi}}\left(\vec{r}_0 + \frac{\vec{u}}{N\rho(\vec{r}_0)^{\frac{1}{d}}}, \vec{r}_0 + \frac{\vec{v}}{N\rho(\vec{r}_0)^{\frac{1}{d}}}\right) \sim \frac{J_{d/2}(2|\vec{u}-\vec{v}|)}{(|\vec{u}-\vec{v}|)^{d/2}}$$

e.g. 1d:  $\sim$  sine-kernel



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► 2*d* complex coordinates z = x + iy: eigenfunctions  $(m = \omega = \hbar = 1)$  $\Phi_{n_1,n_2}(x,y) \sim e^{|z|^2/2} \overline{\partial_z}^{n_1} \partial_z^{n_2} e^{-|z|^2}$ ,

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- ▶ bound states:  $0 \le \Omega < \omega = 1$ , no degeneracy
- ► <u>N</u> Fermions, fill up energies  $E_{n_1,n_2}$ : choose  $\Omega$  s.th. first N levels  $E_{n_1,0} < E_{0,1}$  for  $n_1 = 0, 1, \dots, N - 1 \Leftrightarrow 1 - \frac{2}{N} < \Omega < 1$

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- higher LL [Smith et al. '22], Cf. [Haimi, Hedenmalm '13]

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#### - Renyi entanglement entropy

 $S_q(N,A) \equiv rac{1}{1-q} \ln[\operatorname{Tr}((
ho_A)^q)] \ q>1 \Rightarrow ext{von Neumann}$ 

 $\rho_A = \operatorname{Tr}_{A^C}[\rho]$  partial trace over complement  $A^C = \mathbb{C} \setminus A$ , or express with overlap  $\mathbb{A}_{kl} = \int_A d^2 z \Phi_k(z) \overline{\Phi_l(z)}$ 

 $S_q(N, \mathcal{A}) = rac{1}{1-q} \mathrm{Tr}[\ln(\mathbb{A}^q + (I-\mathbb{A})^q)]$  [Klich '06]



**global density** cirular law on disc of radius  $\sqrt{N}$  ( $\neq$  cup)



▶ global density cirular law on disc of radius  $\sqrt{N}$  (≠ cup) choice area  $A = D_R$  centred disc of radius  $R \sim \sqrt{N}$ 

• mean 
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- variance  $Var \mathcal{N}_A \sim \partial A \sim R$  "holography"
- entropy  $S_q(N, A) = \alpha_q \text{Var} \mathcal{N}_A$ ,  $q \ge 1$  (bulk only!)
- hyperuniformity:  $\lim_{N \to \infty} \operatorname{Var} \mathcal{N}_A / \mathbb{E} \mathcal{N}_A = 0$



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- Q1 How general is the variance in non-Gaussian RMT universal? (e.g. anharmonic trap?)
- **Q2** Var  $\sim \partial A$  holds for general A, general RMT?

#### Complex, quaternion and real Ginibre ensemble



GinUE (left) vs. GinSE (middle) vs. GinOE (right): uniform vs. depletion vs. accumulation on ℝ

#### spectral correlations

 $R_k(z_1,\ldots,z_k)$  determinantal vs. Pfaffian

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- ▶ limiting local *R<sub>k</sub>*:
  - agree in bulk & at edge = universal
  - differ along  ${\mathbb R}$

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- limiting local R<sub>k</sub>:
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- **•** study further examples: **global density**  $\neq$  **const.** 
  - normal matrices  $\exp[-V(|z|)]$
  - truncated U(N + M) Haar distributed  $\rightarrow U(N)$
  - products of *m* Ginibre matrices

# Mean and Variance - finite N & general V(|z|)

mean number of eigenvalues

 $\mathbb{E}\mathcal{N}_R \equiv \int_{D_R} d^2 z \; R_{1,N}(z)$ 

#### variance

$$\operatorname{Var}\mathcal{N}_R\equiv\int_{D_R}d^2z\int_{\mathbb{C}\setminus D_R}d^2u|K_N(z,u)|^2$$
 GinUE/SE class

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**Proposition** [LMS '18; A, Byun, Ebke '22]: For V(|z|) admissible:  $C^2$ ,  $V(|z|) \gg \ln |z|$  at  $|z| \gg 1$ ,  $\Delta_z V > 0$ 

 $\mathbb{E}\mathcal{N}_{R} = \sum_{j=0}^{N-1} h_{j}(R)h_{j}^{-1} \text{ GinUE (GinSE } j \rightarrow 2j+1)$   $\text{Var}\mathcal{N}_{R} = \sum_{j=0}^{N-1} h_{j}(R)h_{j}^{-1}(1-h_{j}(R)h_{j}^{-1}) \text{ ditto}$ 

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#### example GinUE:

$$\frac{K_N(z,u) = e^{-\frac{1}{2}|z|^2 - \frac{1}{2}|u|^2} \sum_{j=0}^{N-1} \frac{1}{j!} (z\bar{u})^j}{\Rightarrow h_j(R)/h_j = \gamma(j+1, NR^2)/j!} \Rightarrow \text{circular law}$$

## Full counting statistics - finite N & general V(|z|)

higher order cumulants: centred generating function

$$\ln\left[\mathbb{E}\exp\left[-\mu(\mathcal{N}_{R}-\mathbb{E}\mathcal{N}_{R})\right]\right] \equiv \sum_{p=2}^{\infty} \frac{(-1)^{p}}{p!} \mu^{p} \mathbb{E}_{c} \mathcal{N}_{R}^{p}$$

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**Proposition** [Lacroix-a-chez-Toine, Majumdar, Schehr '18 GinUE; A, Byun, Ebke, Schehr '23 GinSE]

For V(|z|) admissable the cumulants read:

$$\mathbb{E}_{c}\mathcal{N}_{R}^{\rho} = (-1)^{\rho+1}\sum_{j=0}^{N-1} \text{Li}_{1-\rho}\left(1 - h_{j}/h_{j}(R)\right)$$

for GinUE and index  $j \rightarrow 2j + 1$  for GinSE class,

where  $\operatorname{Li}_q(x) = \sum_{k=1}^{\infty} k^{-q} x^k$  is the polylogarithm.

The large-N Limit: Global density and scaling limits



eigenvalues of GinUE

▶ rescaled weight  $w(z) = e^{-NV(|z|)}$ , V(|z|) admissible:

global density  $\rho(z) \equiv \lim_{N \to \infty} R_{1,N}(z) = \frac{2}{\beta} \Delta_z V(|z|) \cdot \chi_D$ 

Frostman's equilibrium measure on droplet *D*, for Ginibre  $V(|z|) = |z|^2$  circular law on disc  $D_{R=1}$ 

**scaling regimes** for variance of particles in *D<sub>R</sub>*, radius *R* 

• origin limit:  $R = t/N^{\delta}$  non-universal

• edge: 
$$R = 1 - \frac{S}{c\sqrt{N}} \rightarrow 1$$
, S fixed

#### Mean and Variance - 3 scaling regimes



origin (left) vs. bulk (middle) vs. edge (right) scaling

**Theorem** [LMS '19 / A, Byun, Ebke '22] For GinUE / GinSE class ( $\beta = 2 / 4$ ) with *V* admissible:

(i) (Bulk) For  $R \in (0, 1)$  fixed, we have

$$\lim_{N \to \infty} \frac{\beta}{\sqrt{N\rho(R)}} \operatorname{Var} \mathcal{N}_R = \frac{2R}{\sqrt{\pi}}$$

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   expect same for GinOE: origin at large R, numerics

# Comparison GinOE/UE/SE: limiting mean at origin

	GinOE	GinUE	GinSE
$\mathbb{E}\mathcal{N}_R$	$R^{2} + \frac{1}{2} - \frac{1}{2}e^{2R^{2}} \operatorname{erfc}(\sqrt{2}R)$	R <sup>2</sup>	$R^2 - \frac{1}{4} + \frac{1}{4}e^{-4R^2}$
R  ightarrow 0	$\sqrt{rac{2}{\pi}} R + O(R^3)$	R <sup>2</sup>	$2R^4 + O(R^6)$
$R  ightarrow \infty$	$R^2 + \frac{1}{2} + O(\frac{1}{R})$	R <sup>2</sup>	$R^2 - rac{1}{4} + O(e^{-4R^2})$



► top: global density for product of m = 1, 3, 10 GinUE  $\rho(z) = \frac{1}{m} |z|^{\frac{2}{m}-2}$ divergence at origin



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- origin behaviour ~  $(logT)^{m-1}T$  not visible,  $R = T/N^{m/2}$



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- similar plots for quaternionic matrices

#### Universality: Truncated Unitary Symplectic Matrices



▶ top: global density for truncated unitary symplectic matrices  $USp(2N(1+c)+1) \rightarrow USp(2N)$ , c = 0.2, 0.8, 2.0

$$ho(z) = rac{c(1+c)}{(1+c-|z|^2)^2}$$
 spike at edge

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- ▶ bottom: lim VarN<sub>R</sub> in bulk (line) & edge regime (dotted) vs. finite N = 250 (red) and 10 k samples (blue stars)
- origin behaviour  $\sim T^4$  is visible,  $R = T/N^{1/2}$
- similar plots for unitary matrices

#### Summary & Open Questions

- N non-interacting Fermions in harmonic trap in 1*d* / in 2*d* in rotating trap ⇔ with B mapped to Gaussian RMT with ℝ / C eigenvalues
- ► universal asymptotic ∀ cumulants in GinUE class & GinSE class for V admissible in bulk and edge limit
- bulk relation entropy variance beyond Gaussian RMT
- ► real matrices GinOE: mean and variance at origin  $\sqrt{}$ 
  - $\rightarrow$  bulk (and edge) conjectured to be universal
- variance  $\sim \partial A$  for more general A in elliptic GinUE

 $\longrightarrow$  talk by Leslie Molag