SLE + GFF != KPZ

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LIOUVILLE MEASURE AND THE KPZ RELATION







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- Different versions: Q-box counting (DS), Hausdorff (RV), we consider a Minkowski version

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• The expected quantum Minkowski dimension is then defined as: $q_{M,E}(A,n) = \inf_{q} \{ \limsup_{n} \mathbf{E} M_{q}^{Q}(A,n) < \infty \}$

• **PROPOSITION:**

Consider some domain and let A be a fixed compact subset in its interior with Minkowski dimension d_M . Then its expected quantum Minkowski dimension $q_{M,E}$ satisfies the KPZ relation:

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- Proof:
 - Scaling lemma for Liouville balls: $\mathbf{E}(\mu_{\gamma}(S_r)^q) = O(1)r^{\left(2+\frac{\gamma^2}{2}\right)q-\frac{\gamma^2q^2}{2}}$
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- What about more natural counterexamples?

KPZ RELATION AND CONTOUR LINES







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- **PROOF**:
 - Sample the SLE4 and look at scaling of dyadic squares intresecting the SLE4
 - Use Jensen to bound: $\mathbf{E}(\mu_{\gamma}(S)^q) \leq \mathbf{E}(\mu_{\gamma}(S))^q$
 - Regularize the field, use Fubini, control field δ -far from the curve; bound neighbourhood of the line
 - Use our knowledge that the Hausdorff (and Minkowski) dimension of the SLE4 is $\frac{3}{2}$



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There is a coupling of the GFF and chordal SLE κ (κ < 8), such that:

- An instance of the GFF can be constructed by sampling first the SLE κ , and then sampling GFF in the slit domain
- SLE κ is measurable with respect to the GFF
- But boundary conditions more complicated, need to also incorporate winding of the curve, given by $w_T(z) := \arg f_T'(z)$
 - Not defined on the curve, blows up nearing the curve





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• COROLLARY:

The usual KPZ relation does not hold in expected Minkowski nor in almost sure Hausdorff version:



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 - Sum over CR-Whitney squares using results on Euclidean fractal dimension of SLE
 - Use this to determine quantum fractal dimension of SLE



SCALING OF CR-WHITNEY SQUARE

• S some dyadic square, we write

$$\mathbf{E}\big(\mu_{\gamma}(S)^{q} | S \in W\big) =$$

$$\mathbf{E}\left(\left(\lim_{\delta \to 0} \int_{S} \delta^{\frac{\gamma^{2}}{2}} e^{\gamma h_{\delta}(z)} dz\right)^{q} | S \in W\right) =$$
$$\mathbf{E}\left(\left(\lim_{\delta \to 0} \int_{S} \delta^{\frac{\gamma^{2}}{2}} e^{\gamma h_{\delta,H_{t}}(z) + \gamma w(z)} dz\right)^{q} | S \in W\right)$$

SCALING OF CR-WHITNEY SQUARE

- Incorporate winding:
 - CR-Whitney: condition the centre of the square to satisfy $CR(z, H_{SLE}) \cong \epsilon$
 - show that winding is up to an additive error constant in each CR-Whitney square
 - determine exponential moments of winding under this specific conditioning
- Use Kahane convexity inequalities for lower bound

WINDING OF THE SLE CURVES



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- We determined the winding $(\arg f_{\tau}'(z))$ of the chordal SLE κ conditioned to pass ϵ -close of a fixed point z.
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- The notions of windings are different, but (should) agree asymptotically near the curve

• THEOREM:

Fix $z \in \mathbf{H}$; $0 < \kappa < 8$. Let τ be the first time that the SLE κ cuts z from infinity.

For ϵ small enough, conditioned on $CR(z, H_{SLE}) \simeq \epsilon$ the exponential moments of winding $w(z) \coloneqq \lim_{t \to \tau} \arg f_t'(z)$ are given by

 $\mathbf{E}(e^{\lambda w(z)}|\mathrm{CR}(z,H_{SLE}) \asymp \epsilon) \asymp \epsilon^{-\frac{\lambda^2 \kappa}{8}}$

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- This reduces our problem to a diffusion problem (as in **(B)**): Winding w(z) = $\int_0^{\tau} \cot \frac{X_s}{2} ds$ where τ is now the first exit time from $[0,2\pi]$ for the diffusion

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- Same diffusion studied in papers of (L), (SSW)

- **PROOF** the ideal world:
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• If this was the case for our conditioning, we could calculate: $Ee^{\lambda w(z)} | \tau \in [T, T + c] =$ $Ee^{\lambda \int_0^{\tau} \cot \frac{X_s}{2} ds} | \tau \in [T, T + c] =$ $Ee^{-\lambda \frac{\sqrt{\kappa}}{2} B_{\tau} + \lambda X_{\tau}} | \tau \in [T, T + c] = O(1)e^{\frac{\lambda^2 \kappa}{8}T}$

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WINDING OF CHORDAL SLE:

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 - The error for the main contribution is given by: $e(x,s) = \frac{\phi_0'(x)}{\phi_0(x)} - \frac{P_x'(\tau \in [T-s, T-s+c])}{P_x(\tau \in [T-s, T-s+c])}$

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 - To bound the rest term, we use more probabilistic arguments

THANK YOU!

