



The 3rd Soft Matter Summer School :
Polymer Sciences in Biology
June 21 - July 4, 2015

“Neutron and Rheology”

Mitsuhiro Shibayama,
ISSP, U. Tokyo

**Part 1. Structural Analyses of Polymers by
Small Angle Neutron Scattering**

**Part 2. Contrast Variation SANS
- The basics and applications -**

**Part 3. Rheo-SANS Studies on Structure Evolution
in Polymer-particle Aqueous Solutions**



Prof. Tsukamoto

eel

Kashiwa Campus, U. Tokyo



Kashiwa Campus, U. Tokyo

Grad. School of Frontier Sci.

Env. Sci.
Life Sci.
Frontier Mater. Sci.

ISSP

Cosmic Ray Inst

Super Kamiokande

Origin of Universe

Institute for the Physics
and Mathematics of the Universe

Guest House

Prof. Koshiba

Prof. Murayama

Since 2000



One of three major campuses of U. Tokyo



The 3rd Soft Matter Summer School :
Polymer Sciences in Biology
June 21 - July 4, 2015

“Neutron and Rheology”

Mitsuhiro Shibayama,
ISSP, U. Tokyo

Part 1. Structural Analyses of Polymers by Small Angle Neutron Scattering

1. Introduction
2. Neutron and neutron scattering
3. Neutron scattering theory
4. Small angle Scattering
5. Applications



The 3rd Soft Matter Summer School :
Polymer Sciences in Biology
June 21 - July 4, 2015

“Neutron and Rheology”

Mitsuhiro Shibayama,
ISSP, U. Tokyo

Part 1. Structural Analyses of Polymers by Small Angle Neutron Scattering

1. Introduction
2. Neutron and neutron scattering
3. Neutron scattering theory
4. Small angle Scattering
5. Applications



Scattering vs spectroscopy



Spectroscopy
Scattering

The first observation of total ellipse in the South Pole

~ 8 am, November, 24, 2003 by NHK

<http://www.neptune.carina.gr.jp/same/solaec03/>

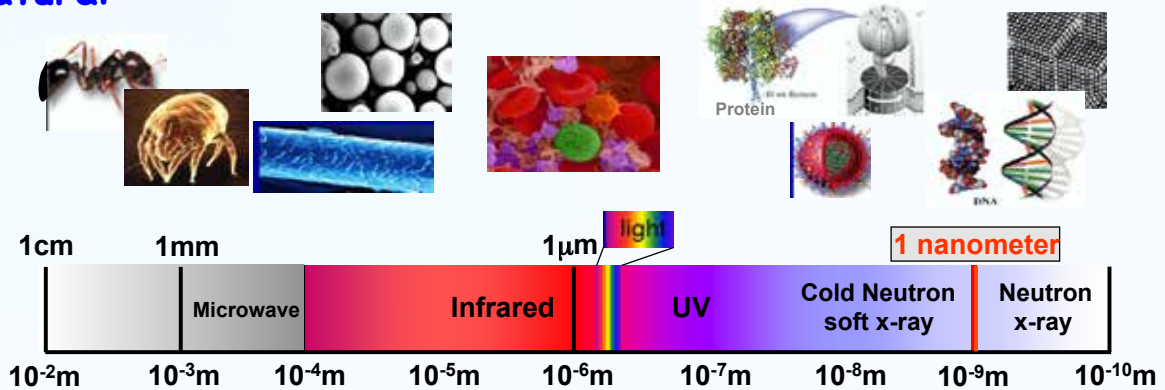
5



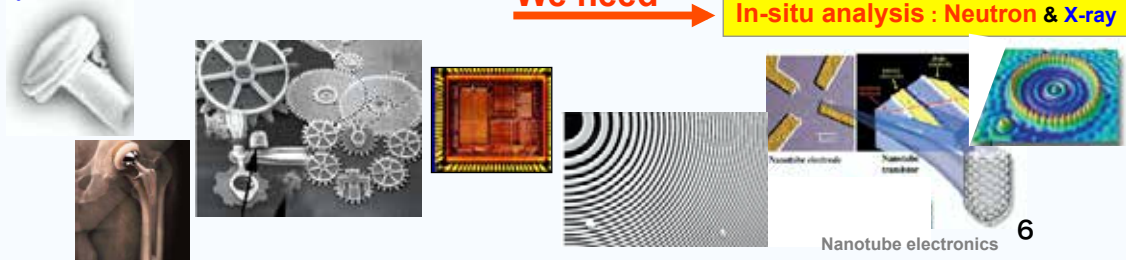
Towards Nanometer Technology

Courtesy of S. Choi, KAIST

Natural



Manmade





Nano-structure characterization



SEM



AFM

**Electron microscopy,
Atomic microscopy:**
*easy, but local structure,
thin film or surface only*



Light scattering
*mesoscopic structure
necessary to be transparent*



SAXS (Lab.)



SAXS (SPRING-8)

X-ray scattering
*(lab): easy, but weak intensity
(large facility): strong beam,
but irradiation damage*



SANS (J-PARC/JRR-3)

Neutron scattering
*(large facility): difficult to get beam time,
low-resolution, difficult for deuteration,
but large contrast*



Neutron

What is Neutron?

$$m_n \sim 1g/N_{\text{Avogadro}}$$

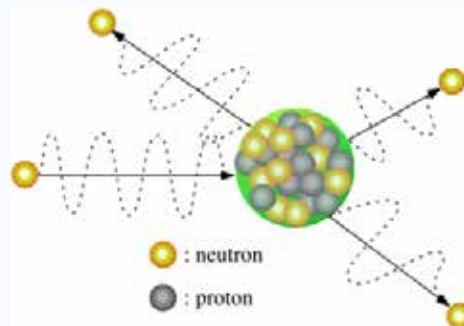
Radius; 1.5×10^{-13} cm (10^{-5} of the radius of hydrogen atom)

Mass; 1.6749×10^{-27} kg (nearly equal to that of proton)

Charge; 10^{-18} e (substantially zero)

Half-life time; 10.3 min ($n \rightarrow p + \text{meson}$)

Quantum spin number; 1/2



Generation of neutrons:

Atomic reactor or accelerator

Kinds of neutrons

Cold neutrons; $E \leq 0.002$ eV

Thermal neutrons; $0.002 \leq E \leq 0.5$ eV

Epithermal; $0.5 \leq E \leq 500$ eV

Fast neutrons; $500 \text{ eV} \leq E$

*Similar to
the electromagnetic wave,
i.e., **g**ay, X-ray, UV, VL, IR, ...*

History of neutron scattering:

Discovery: Chadwick (1932)

Observation of diffraction (1936)

Polymer research by neutron scattering (1972)



Chadwick, Nobel winner, 1935



Brockhouse & Shull, Nobel winner, 1994



Sizes of neutron and atom

Outstanding contribution to due to the three characteristic properties

Structure
(charge = 0)

magnetism
(spin = 1/2)

Where atoms are
原子の位置



What atoms do
原子の運動

motion
(mass = 1u)

1.8Å ~ 293K




question:

If the diameter of Tokyo Dome (~ 100m), how large are the atom and neutron?

- (1) apple (~ 10cm),
- (2) orange (~ 5cm),
- (3) Japanese pinball (11mm),
- (4) Jintan ball (1mm)



Why Neutrons ?

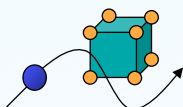
 **Mass** **No Charge** **Spin 1/2** ↑



No charge



Deep penetration



Wavelength ~ Å, nm
(thermal & cold neutron)



**Atomic length scale
& Nano length scale**



Energy ~ meV



**Same magnitude as
basic excitations in solids**

(solid state physics)



Spin = 1/2



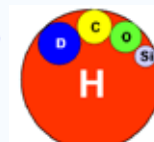
**Magnetic structure &
dynamics**



Interacts with nuclei



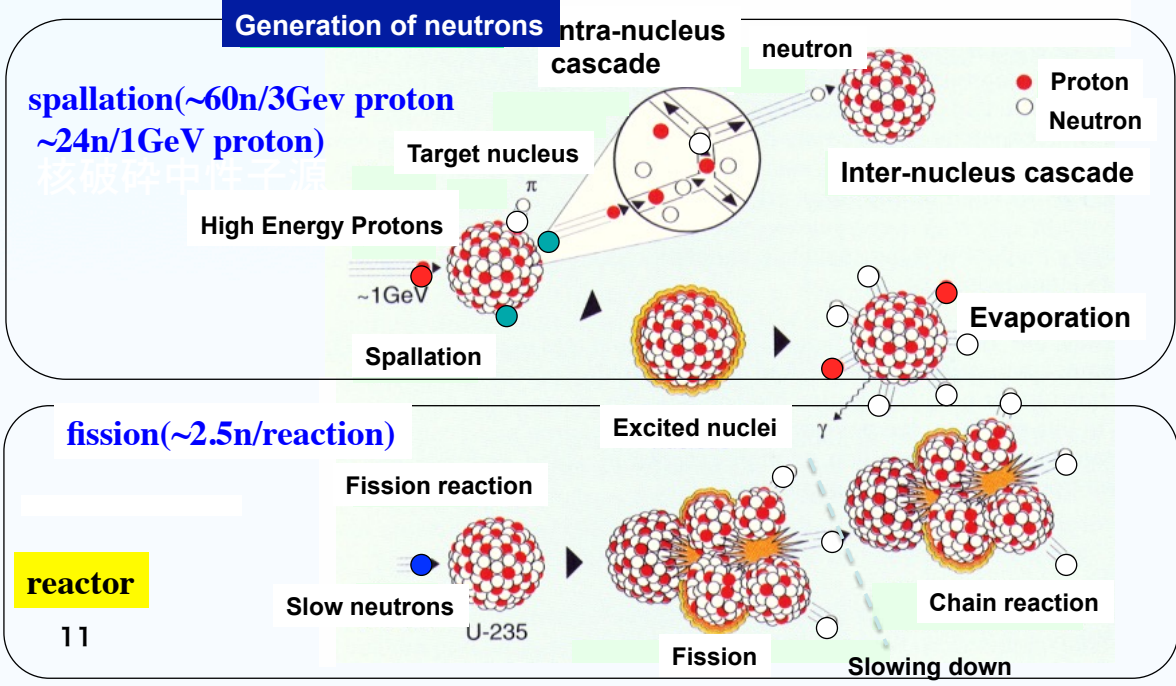
(solid state physics)
Contrast variation
(soft matter)





Generation of neutrons

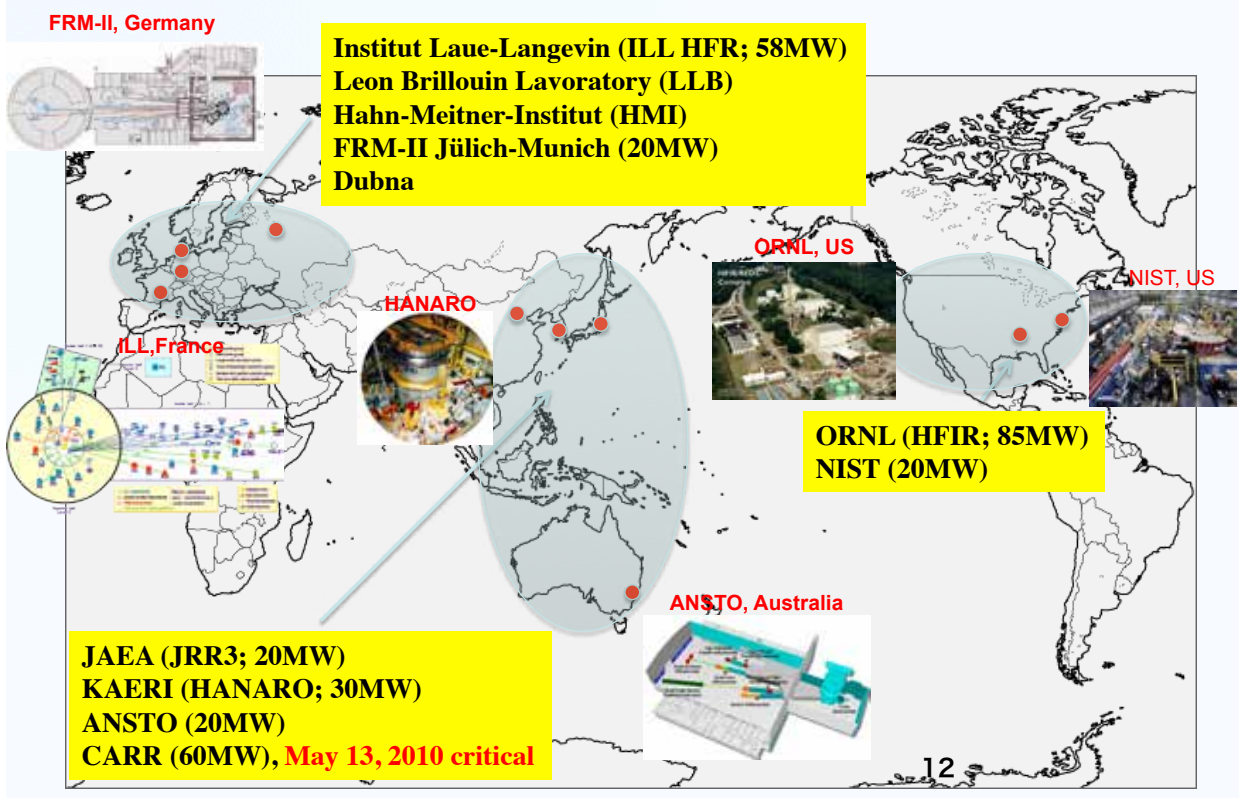
very effective (no. neutrons \propto proton power)
 (spallation 1MW \sim reactor 15MW)
 low heat generation (\sim proton power)



11



Research reactors in the world

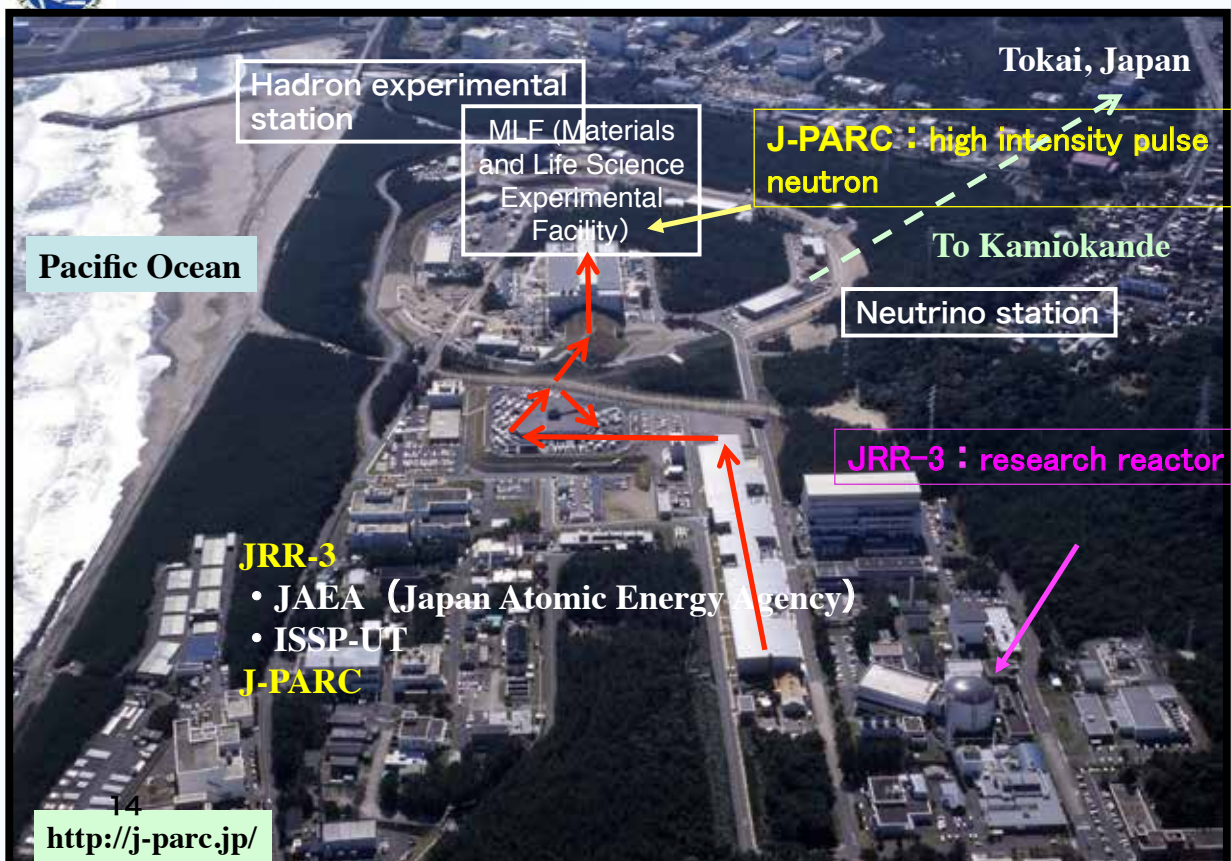


12

Pulse Neutron Source in the world



Neutron Science at JRR-3 and J-PARC





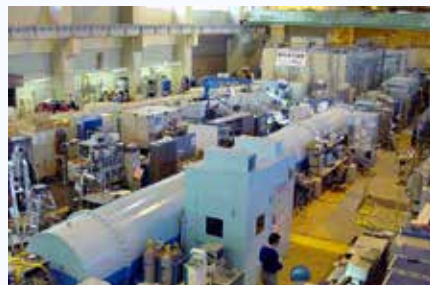
Neutron Science Lab., ISSP, U. Tokyo NSL-ISSP : SANS-U



JRR-3@Japan Atomic Energy Agency
(20MW Research Reactor)



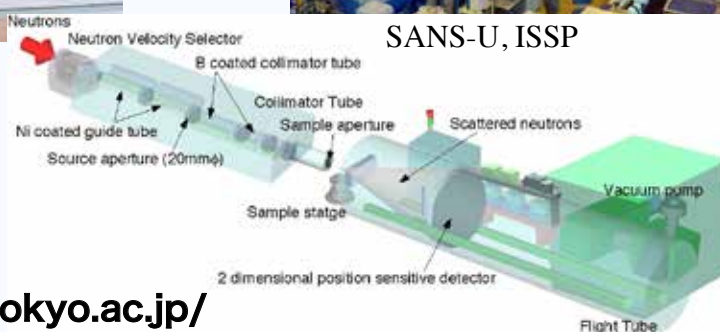
ISSP,
Kashiwa



SANS-U, ISSP



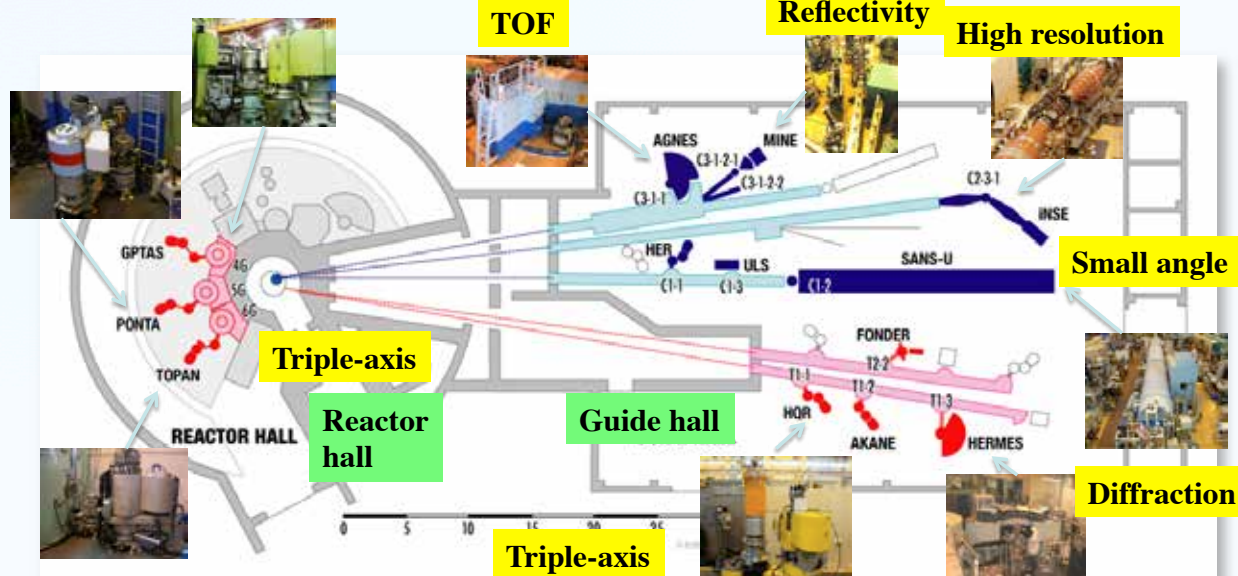
U. Tokyo



<http://neutrons.issp.u-tokyo.ac.jp/>



University-owned Instruments at JRR-3



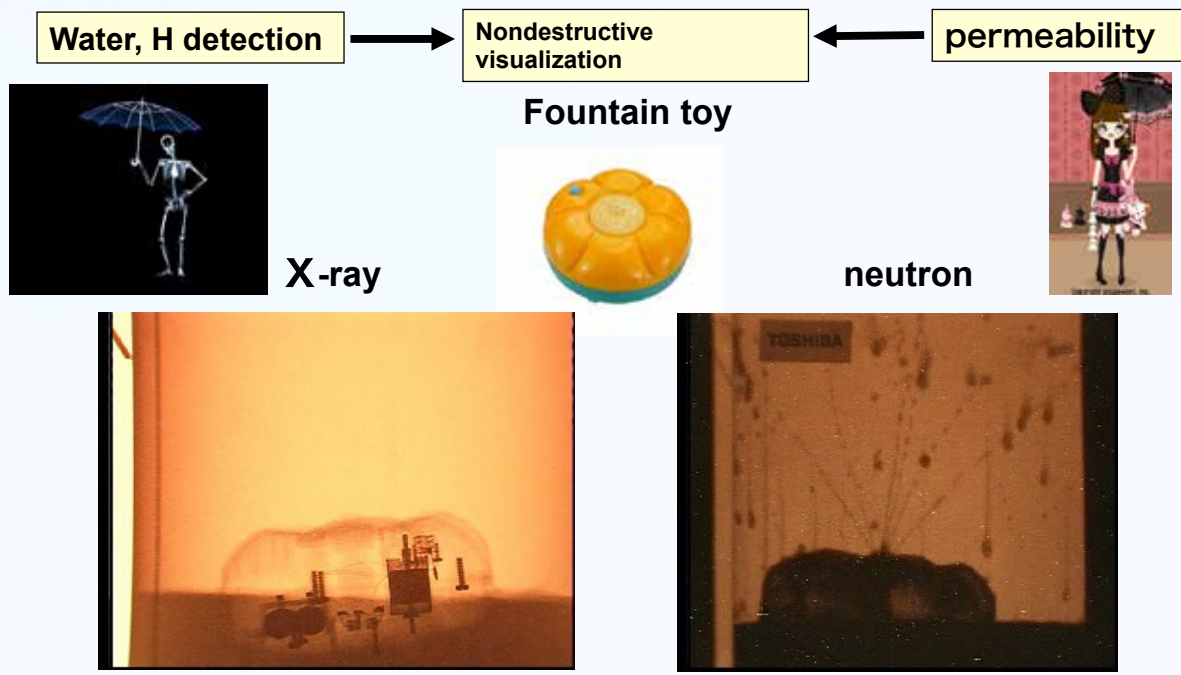
University-owned instruments: 14, **ISSP 9, Tohoku U. 3, Kyoto U. 2**
 No. proposals: ~300
 No. users (man.day) : in-house 2000, outside 5000, total 7000
 No. papers : ~100 /y



小角中性子散乱装置 SANS-U

17

Comparison of X-ray and neutron



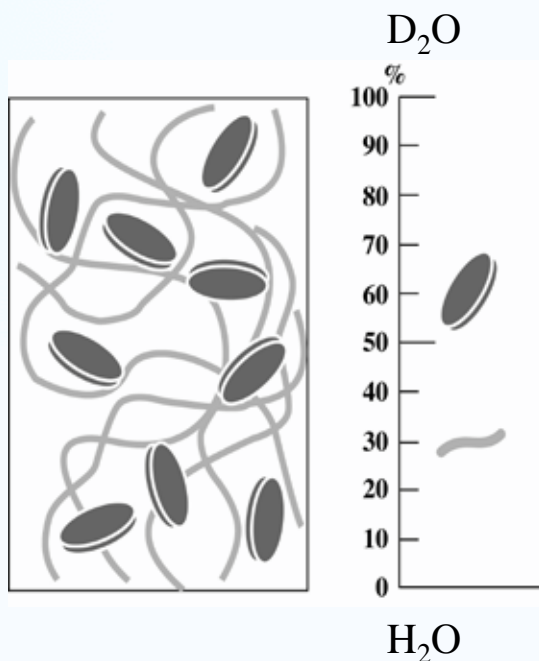


Neutron contrast



19

labeling



Contrast matching



The 3rd Soft Matter Summer School :
Polymer Sciences in Biology
June 21 - July 4, 2015

“Neutron and Rheology”

Mitsuhiro Shibayama,
ISSP, U. Tokyo

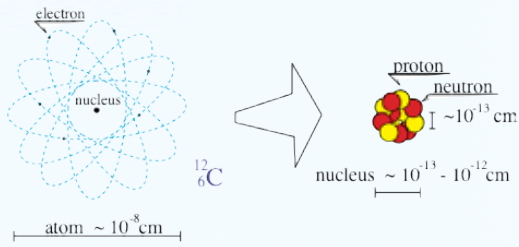
Part 1. Structural Analyses of Polymers by Small Angle Neutron Scattering

1. Introduction
2. Neutron and neutron scattering
3. Neutron scattering theory
4. Small angle Scattering
5. Applications



1. Atom and Neutron

atom and nucleus

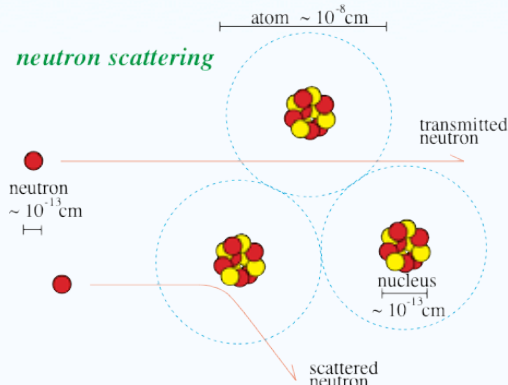


Size of an atom
($\approx 0.1 \text{ nm} = 10^{-8} \text{ cm}$)

A neutron is about $1/10^5 (\approx 10^{-13} \text{ cm})$ as large as an atom.

(A neutron is 1 cm large if an atom is 1 km large.)

neutron scattering



With an eye of neutron, the nuclei in materials are so dilute that most of neutrons pass through the materials without scattering.

When a neutron passes near a nucleus, nuclear scattering takes place.



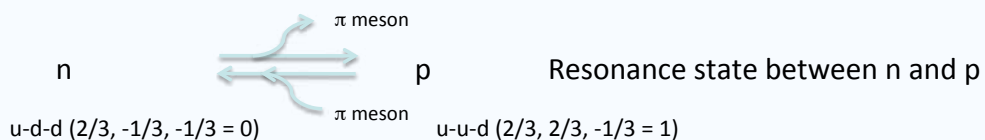
2. Properties of neutron

		Three Generations of Matter (Fermions)			
		I	II	III	
mass	$m_n = 1.675 \times 10^{-27} \text{ kg}$	2.4 MeV	1.27 GeV	171.2 GeV	0
Spin quantum number	$s = 1/2 \text{ } (-1/2)$; Fermion	$1/2$	$1/2$	$1/2$	1
Mag. moment	$\mu_n = -1.913 \mu_N$ μ_N : nuclear mag. Moment, $3.152 \times 10^{-14} \text{ MeV/T}$	$2/3$	$2/3$	$2/3$	0
Lamor freq.	29.16 (MHz/Tesla)	$1/2$	$1/2$	$1/2$	1
Life time	$885.9 \pm 0.9 \text{ s}$ (ca15min)	$1/2$	$1/2$	$1/2$	0
Quark comp.	u-d-d	up	charm	top	photon
		down	strange	bottom	gluon
		electron neutrino	muon neutrino	tau neutrino	weak force
		electron	muon	tau	weak force

Annihilation of neutron (β^- annihilation)

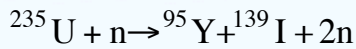
$$n \rightarrow p^+ + e^- + \bar{\nu} (+0.77 \text{ MeV})$$

Wikipedia



3. Generation of neutron (1)

1. reactor: fission of ^{235}U



$$235 + 1 > 95 + 139 + 2 \quad \text{mass defect} \quad \Delta E = mc^2$$

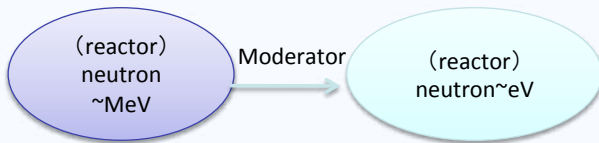
(236) (244)

- * Generation of 2~3 neutron by 1 fission
- * Energy of ca. 200MeV $\approx 3.2 \times 10^{-11}$ J
(8.2×10^{10} J/1g U)

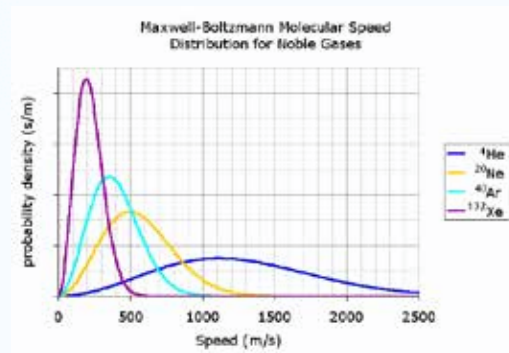
<ref.> velocity distribution of noble gas

$$f(\mathbf{v}) = \left(\frac{m}{2\pi kT}\right)^{3/2} \exp\left(-\frac{m\mathbf{v}^2}{2kT}\right), \quad \mathbf{v} = (v_x, v_y, v_z)$$

$$f(v)dv = 4\pi\left(\frac{m}{2\pi kT}\right)^{3/2} \exp\left(-\frac{mv^2}{2kT}\right)v^2dv$$



Maxwell distribution characterized by the temperature of the moderator

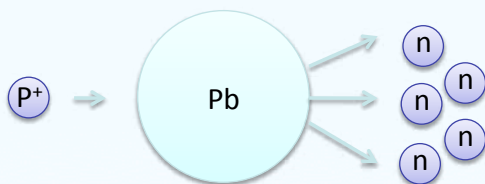


The velocity distribution of neutrons is the Same as that of noble gas.

Q: Calculate the most probable velocity of argon gas at $T=300\text{K}$.

3. Generation of neutron (2)

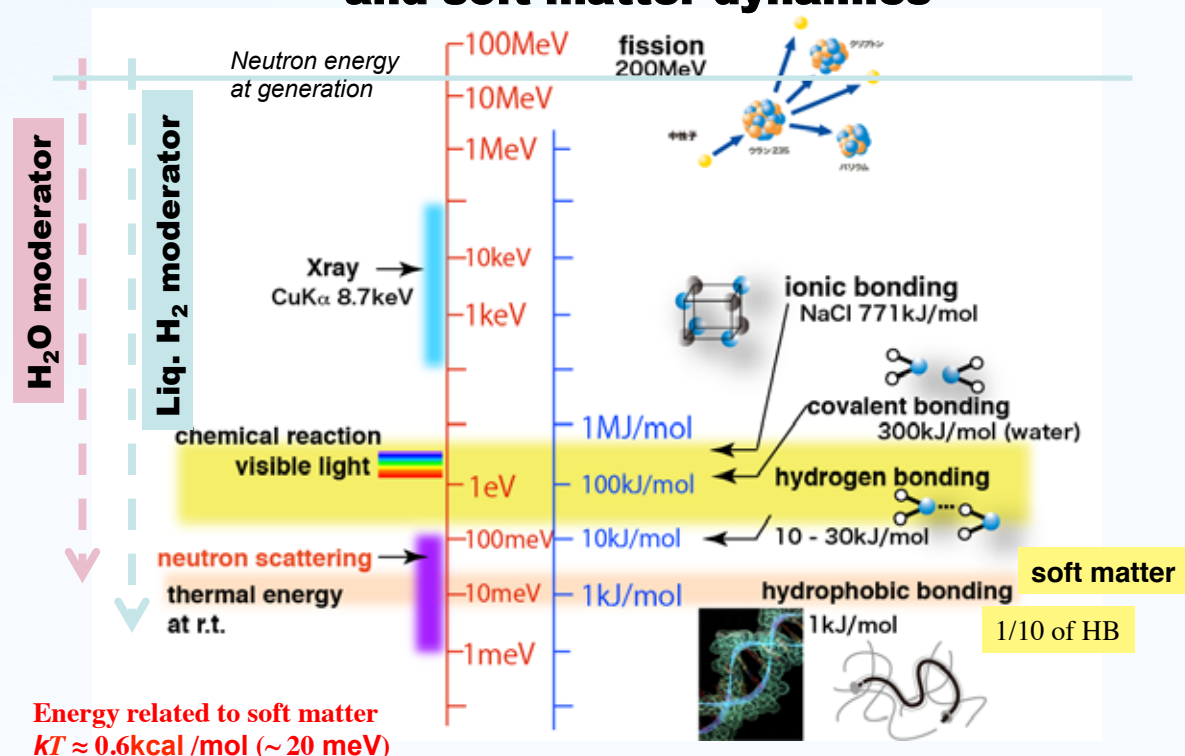
2. spallation



By a collision of a proton with 1GeV with a target nucleus, 20 ~ 25 of high energy neutrons come out.



4. Energies of neutrons, chemical reactions, and soft matter dynamics



25

Soft and flexible



Part 1

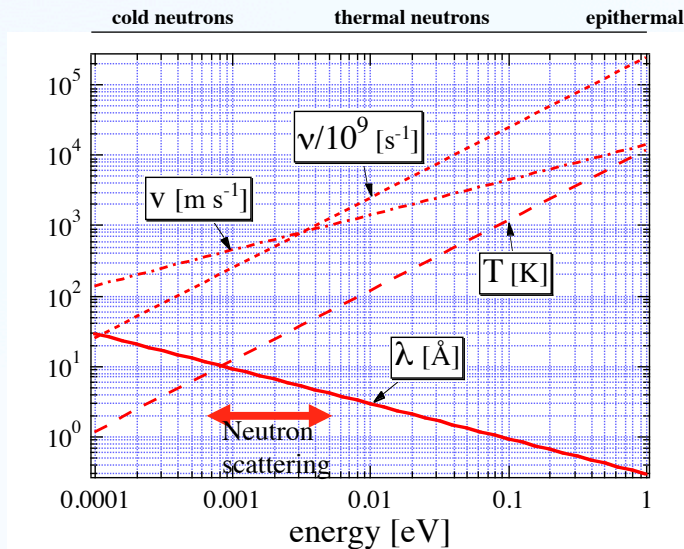
5. Properties of neutrons

	suchness
energy	$E = mv^2/2 = p^2/2m$; (Einstein, particle wave)
wavelength	$\lambda = h/mv = h/p$; (de Broglie wave)
temperature	$E = kT$
velocity	$v = (2E/m)^{1/2}$
flux	$\Phi(v) \sim v^3 \exp(-mv^2/2kT_{\text{mod}})$ (T_{mod} : moderator temperature)

	category
10^{-7} eV	ultra cold neutron
0.1 - 10 meV	cold neutron (moderator: liquid H_2)
10 - 100 meV	thermal neutron ($T_{\text{mod}} \approx$ room temp.)
100 - 500 meV	hot neutron
> 500 meV	epithermal neutron



6. Velocity, wavelength, and wave number of neutrons



$$E = \frac{1}{2}mv^2 \text{ (Einstein), } \textbf{particle} \quad (1)$$

$$E = h\nu; \quad p = mv = \frac{h}{\lambda} \quad \text{(de Broglie)} \quad \textbf{wave} \quad (2)$$

$$E = kT \quad \textbf{energy} \quad (3)$$

$$\lambda_{\text{mean}} = \frac{h}{\sqrt{3m_n kT}} \sim T^{-1/2} \quad \textbf{Maxwell distribution} \quad (4)$$

$$\lambda_{\text{max}} = \sqrt{\frac{3}{2}} \lambda_{\text{mean}} = \frac{30.81}{\sqrt{T}} [\text{\AA}] \quad (5)$$

$$T[\text{K}] = 11.605 \times 10^3 E[\text{eV}] \quad (6)$$

$$\lambda[\text{\AA}] = \frac{0.2860}{\sqrt{E[\text{eV}]}} \quad (7)$$

$$v[\text{m/s}] = 1.383 \times 10^4 \sqrt{E[\text{eV}]} \quad (8)$$

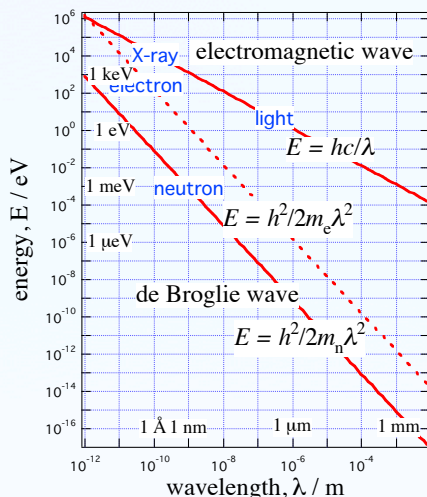
Neutron has **wave-particle duality**.

The velocity, wavelength, and wave number of neutrons depend on temperature.

Only cold neutrons and thermal neutrons are used for small angle neutron scattering. 27



7. Energy dispersion



Scattering by photon:

For photons, the relationship between energy ϵ and wavenumber $k = 2\pi/\lambda$ is given by

$$\epsilon = \hbar\omega = \frac{hc}{\lambda} \quad h = 6.626 \times 10^{-34} [\text{J}\cdot\text{s}]$$

For visible light $\epsilon \approx 1\text{eV}$, $\lambda = 0.4 \sim 0.7 \times 10^4 \text{\AA}$

Hence, light is a suitable probe for μm -ordered structures.

For \AA -ordered structures, photons with

$\epsilon \approx 10^4 \text{eV} = 10\text{keV}$ with are necessary and X-ray is the Best means.

Scattering by electrons:

Electrons with mass m_e has the following dispersion relationship.

$$\epsilon = \frac{\hbar^2 k^2}{2m_e} = \frac{h^2}{2m_e \lambda^2} \quad m_e = 9.109 \times 10^{-31} \text{ kg}$$

$\lambda = 1 \text{\AA}$, $\epsilon \approx 100\text{eV}$

Scattering by neutrons:

The same dispersion eq. as for electrons

$$\epsilon = \frac{\hbar^2 k^2}{2m_n} = \frac{h^2}{2m_n \lambda^2}$$

Note that the mass of neutron mass is very different from that of electron.
 $m_n = 1.675 \times 10^{-27} \text{ kg}$ (ca 1800 times larger than e)

$\lambda = 1 \text{\AA}$, $\epsilon \approx 0.05\text{eV}$

~ thermal energy

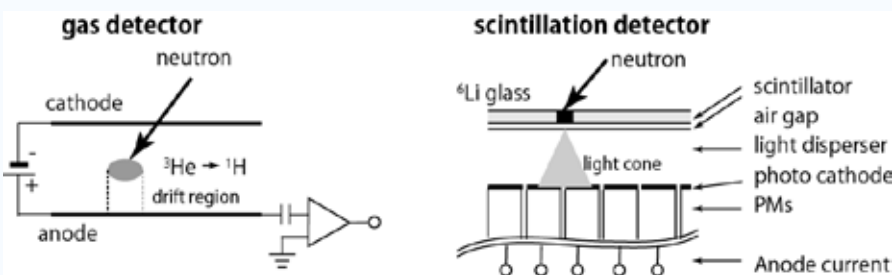
8. Detection of neutrons

Neutron: electroneutral

By generating electric charges via nuclear reaction, and counting them

	Cross section, (25meV)	Generated particles	energy— [MeV]	Total energy[MeV]
$n + {}^3\text{He}$	5333b	p, ${}^3\text{T}$	0.57, 0.2	0.77
$n + {}^6\text{Li}$	941b	${}^3\text{T}$, ${}^4\text{He}$	2.74, 2.05	4.79
$n + {}^{10}\text{B}$	3838b	${}^4\text{He}$, ${}^7\text{Li}$, γ	1.47, 0.83, 0.48	2.30
$n + {}^{235}\text{U}$	681b	fission		1 - 2

b ; a unit of scattering **cross section** b = barn(10^{-24} cm^2)



“Neutron and Rheology”

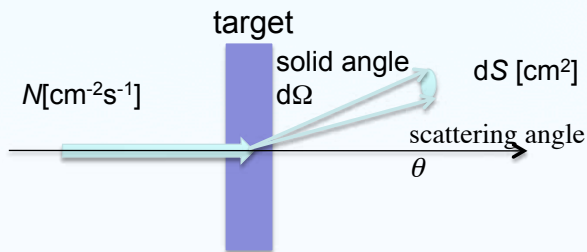
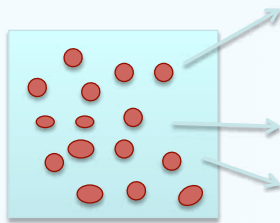
Mitsuhiro Shibayama,
ISSP, U. Tokyo

Part 1. Structural Analyses of Polymers by Small Angle Neutron Scattering

1. Introduction
2. Neutron and neutron scattering
- 3. Neutron scattering theory**
4. Small angle Scattering
5. Applications

1. What is scattering?

Stones in a box



Bullet (probe)	Stones (target)
photon	atom
neutron	nucleus
electron	molecule
proton	elemental particle
....

By randomly shooting a gun, you may guess what in the box are.

- numbers of bullets per unit area and time $N[\text{cm}^{-2}\text{s}^{-1}]$
 - number of bullets scattered in the solid angle $dn[\text{s}^{-1}]$
- $$dn \propto NdS/r^2 = Nd\Omega$$
- $$dn = N\sigma(\theta)d\Omega$$

$\sigma(\theta)$: proportional const.
 dimension $[\sigma] = \frac{[dn d\Omega]}{[N]} = \frac{[\text{s}^{-1}]}{[\text{cm}^{-2}\text{s}^{-1}]} = [\text{cm}^2]$

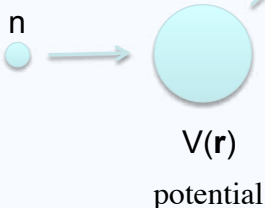
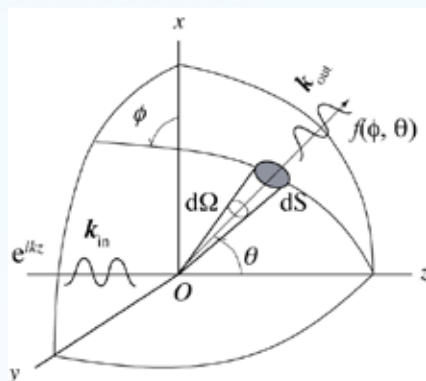
$\sigma(\theta)$: **scattering cross section**

Total cross section $\sigma_{\text{tot}} = \int \sigma(\theta) d\Omega$

Scattering by an atom	Scattering by a nucleus
Bohr radius $a_0 = 5.29 \times 10^{-9} \text{ cm}$	$\sigma_{\text{tot}} \sim 10^{-24} \text{ cm}^2$
$\rightarrow \sigma_{\text{tot}} \approx a_0^2 = 2.80 \times 10^{-17} \text{ cm}^2$	

2. Scattering amplitude

Scattered wave



$$\psi(\mathbf{r}) = \frac{1}{r} f(\Omega) e^{ik \cdot \mathbf{r}} \quad f(\mathbf{r}) : \text{scattering amplitude}$$

Stationary wave

$$H = H_0 + V(\mathbf{r}) \quad V(\mathbf{r}) : \text{potential}$$

$$H_0 = \frac{p^2}{2\mu} \quad \mu: \text{reduced mass} \quad \mu = \frac{m_1 m_2}{m_1 + m_2}$$

p : momentum

Schrödinger eq.

$$H\psi = E\psi$$

$$\psi(\mathbf{r}, t) = \phi(\mathbf{r}) e^{-iEt/\hbar} \quad \phi(\mathbf{r}) : \text{a solution of the eigen function}$$

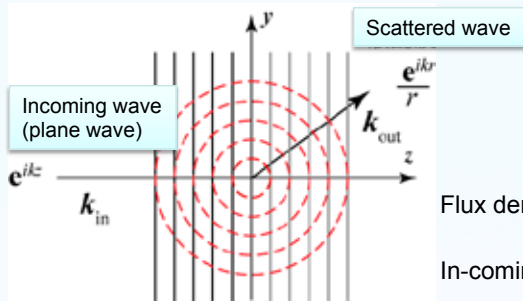
$$\left[-\frac{\hbar^2 \Delta}{2\mu} + V(\mathbf{r}) \right] \phi(\mathbf{r}) = E\phi(\mathbf{r}), \quad E = \frac{p^2}{2\mu} = \frac{\hbar^2 k^2}{2\mu}$$

Let $V(\mathbf{r}) \equiv (\hbar^2 / 2\mu) U(\mathbf{r})$, and

$$\left[\Delta + k^2 - U(\mathbf{r}) \right] \phi(\mathbf{r}) = 0$$

3. Scattering cross section and amplitude

The wave equation for scattering is a sum of plane wave e^{ikz} and $\sim e^{ikr}/r$.



$$\varphi_k(\mathbf{r}) \underset{r \rightarrow \infty}{\sim} e^{ikz} + f_k(\theta, \varphi) \frac{e^{ikr}}{r}$$

Plane wave Scattering amp.

Calculation of scattering cross section

Flux density: $J(r) = \frac{1}{\mu} \text{Re} \left[\varphi^*(r) \frac{\hbar}{i} \nabla \varphi(r) \right]$

In-coming wave $J_{in} = \frac{\hbar k}{\mu}$

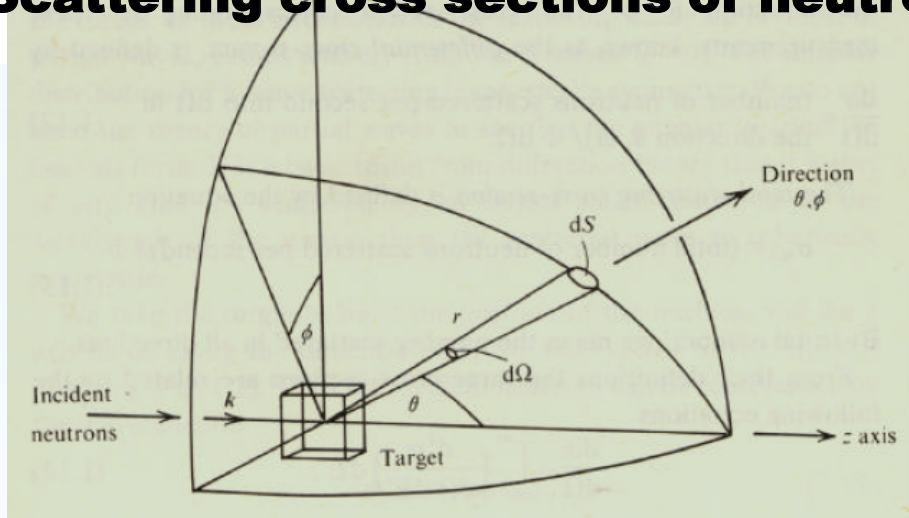
Scattered wave: $J_{out}(r) = \frac{\hbar k}{\mu} \frac{1}{r^2} |f_k(\theta, \varphi)|^2$

Hence, $dn = J_{out} dS = J_{out} r^2 d\Omega = \frac{\hbar k}{\mu} |f_k(\theta, \varphi)|^2 d\Omega$

From the definition of cross section, $dn = J_{in} \sigma d\Omega = \frac{\hbar k}{\mu} \sigma d\Omega$

$\sigma = |f_k(\theta, \varphi)|^2$ Note: the cross section is a square of the scattering amplitude

4. Scattering cross sections of neutrons



Φ = No. of incident neutrons per cm^2 per second

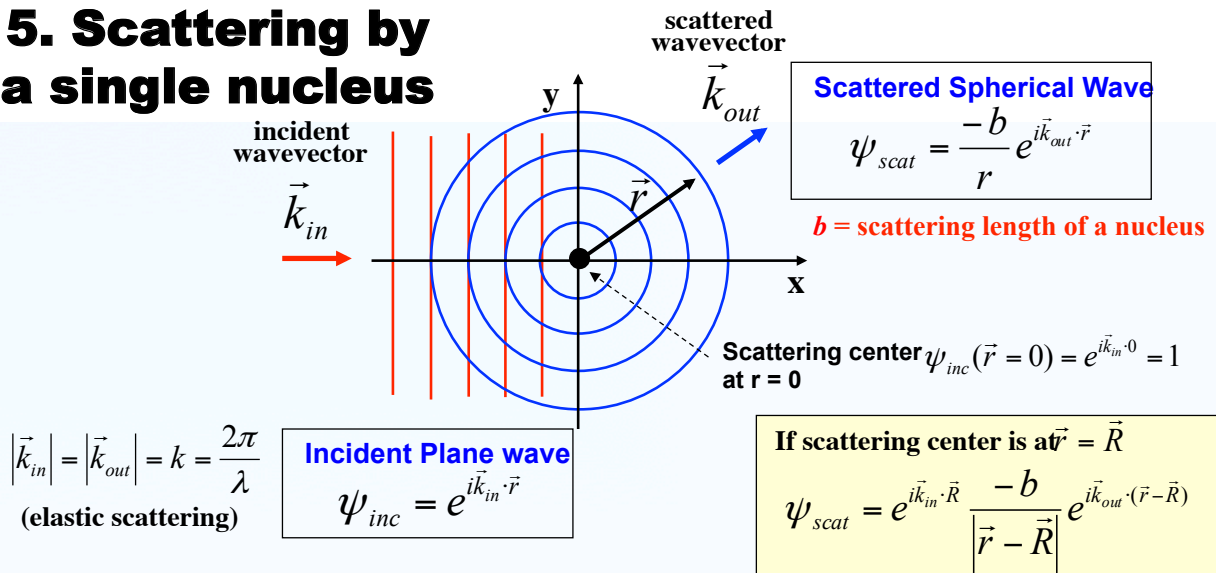
σ = Total No. of neutrons scattered per second / Φ

σ measured in barns
1 barns = 10^{-24} cm^2

Elastic $\rightarrow \frac{d\sigma}{d\Omega} = \frac{\text{No. of neutrons scattered per second into } d\Omega}{\Phi d\Omega}$

Inelastic $\rightarrow \frac{d^2\sigma}{d\Omega dE} = \frac{\text{No. of neutrons scattered per second into } d\Omega \text{ and } dE}{\Phi d\Omega dE}$

5. Scattering by a single nucleus



No. of incident neutrons passing through unit area per second

$$\Phi = v |\psi_{inc}|^2 = v \quad v: \text{Neutron velocity}$$

No. of scattered neutrons passing through unit area per second

$$\sigma_{total} = \int \frac{d\sigma}{d\Omega} d\Omega = 4\pi b^2$$

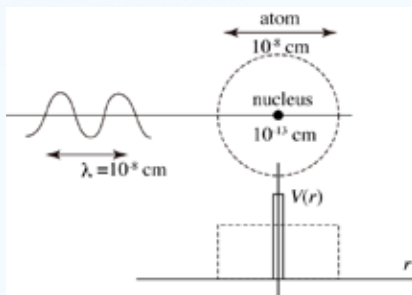
$$v |\psi_{scat}|^2 dS = v \left(\frac{b^2}{r^2} \right) dS = v b^2 d\Omega$$

$$\frac{d\sigma}{d\Omega} = \frac{\text{No. of Neutrons Scattered per sec. into } d\Omega}{\Phi d\Omega} = \frac{v |\psi_{scat}|^2 dS}{v d\Omega} = \frac{v b^2 d\Omega}{v d\Omega} = b^2$$



Part 2

6. Fermi's pseudo potential



$$V(r) = \frac{2\pi\hbar^2}{m} b \delta(\mathbf{r} - \mathbf{R})$$

The interaction potential for nuclear scattering is given by A delta function.

b : the scattering length

$$b = b_1 + i b_2 \quad ; \quad \text{complex number} \\ \text{(imaginary part: absorption)}$$

Differential cross section (Born approx.)

$$\begin{aligned} \sigma^{(B)}(\theta, \varphi) &= \frac{\mu^2}{4\pi^2 \hbar^4} \left| \int d^3 r e^{-i\mathbf{Q} \cdot \mathbf{r}} V(r) \right|^2 \\ &= \frac{\mu^2}{4\pi^2 \hbar^4} \left| \int d^3 r e^{-i\mathbf{Q} \cdot \mathbf{r}} \frac{2\pi\hbar^2}{m} b \delta(\mathbf{r} - \mathbf{R}) \right|^2 \\ &= |b|^2 \end{aligned}$$

Note: the scattering cross section does not depend on Q .
(A Fourier tr. of a delta fn. is a constant.)

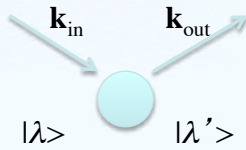
$$U(r) = \frac{2\mu}{\hbar^2} V(r) = 4\pi b \delta(\mathbf{r} - \mathbf{R})$$

The scattering amplitude

$$f(\theta, \varphi) = -\frac{1}{4\pi} \int d^3 r' e^{-i\mathbf{Q} \cdot \mathbf{r}'} U(r') = -b$$

b : the scattering length

7. Inelastic scattering



In the case of inelastic scattering, the wavelength is also changed as follows,

$$\left(\frac{d\sigma}{d\Omega} \right)_{k,\lambda \rightarrow k',\lambda'} = \frac{k'}{k} \left(\frac{m}{2\pi\hbar^2} \right)^2 \langle k', \lambda' | V | k, \lambda \rangle^2$$

On the other hand, the energy and momentum transfer are preserved.

$$\left(\frac{\partial^2 \sigma}{\partial \Omega \partial \omega} \right)_{k,\lambda \rightarrow k',\lambda'} = \frac{k'}{k} \left(\frac{m}{2\pi\hbar^2} \right)^2 \sum_{\lambda} P_{\lambda} \langle k', \lambda' | V | k, \lambda \rangle^2 \delta(\hbar\omega + E_{\lambda} - E_{\lambda'})$$

P_{λ} : distribution of λ

37

8. Scattering by many nuclei

The scattered wave from many nuclei located at \vec{R}_j

$$\psi_{scat} = \sum_j e^{i\vec{k}_{in} \cdot \vec{R}_j} \frac{-b_j}{|\vec{r} - \vec{R}_j|} e^{i\vec{k}_{out} \cdot (\vec{r} - \vec{R}_j)} = e^{i\vec{k}_{out} \cdot \vec{r}} \sum_j \frac{-b_j}{|\vec{r} - \vec{R}_j|} e^{-i(\vec{k}_{out} - \vec{k}_{in}) \cdot \vec{R}_j}$$

Therefore

$$\frac{d\sigma}{d\Omega} = \frac{v |\psi_{scat}|^2 dS}{vd\Omega} = \frac{dS}{d\Omega} \left| e^{i\vec{k}_{out} \cdot \vec{r}} \sum_j \frac{b_j}{|\vec{r} - \vec{R}_j|} e^{-i(\vec{k}_{out} - \vec{k}_{in}) \cdot \vec{R}_j} \right|^2$$

If we measure far enough away so that $r \gg R_i$, then $|\vec{r} - \vec{R}_j| \approx r$ $d\Omega = \frac{dS}{r^2}$

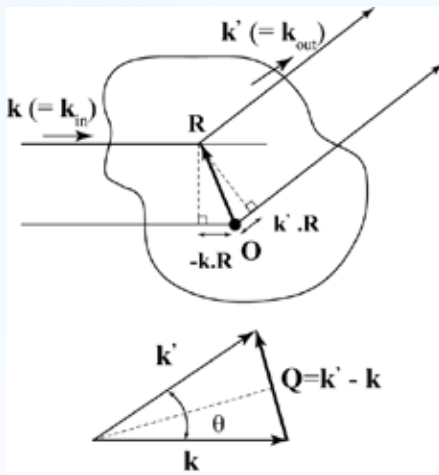
$$\frac{d\sigma}{d\Omega} = \left| \sum_j b_j e^{-i\vec{Q} \cdot \vec{R}_j} \right|^2 = \sum_{i,j} b_i b_j e^{-i\vec{Q} \cdot (\vec{R}_i - \vec{R}_j)} \quad \left| e^{i\vec{k}_{out} \cdot \vec{r}} \right|^2 = 1$$

where the wavevector transfer \vec{Q} is defined as

$$\vec{Q} = \vec{k}_{out} - \vec{k}_{in}$$

38

9. Scattering vector



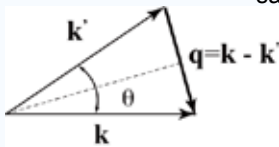
$$\mathbf{Q} \equiv \mathbf{k}_{out} - \mathbf{k}_{in} = \mathbf{k}' - \mathbf{k}$$

$$\begin{aligned} \text{Path difference} &= \mathbf{k}' \cdot \mathbf{R} - \mathbf{k} \cdot \mathbf{R} = (\mathbf{k}' - \mathbf{k}) \cdot \mathbf{R} \\ &= \mathbf{Q} \cdot \mathbf{R} \end{aligned}$$

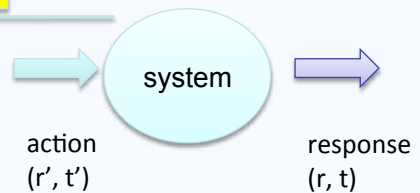
$$|\mathbf{k}| = |\mathbf{k}'| = k = \frac{2\pi}{\lambda}$$

$$\mathbf{Q} = 2k \sin\left(\frac{\theta}{2}\right)$$

Cf. From the viewpoint of **momentum transfer** from incoming particle to sample, the following definition is better.



$$\begin{aligned} \hbar \mathbf{q} : \quad \hbar \mathbf{k} &\rightarrow \hbar \mathbf{k}' \\ \mathbf{q} &\equiv \mathbf{k} - \mathbf{k}' \end{aligned}$$



response = <after | action | before>

10. Coherent and incoherent scattering

Neutron scattering is also dependent on isotopes and the relative orientation of spins. Here, we consider scattering from an assembly of isotopes with $\{R_l, b_l\}$.

Fermi potential

$$V(r) = \frac{2\pi\hbar^2}{m} \sum_l b_l \delta(\mathbf{r} - \mathbf{R}_l)$$

The scattering amplitude

$$\begin{aligned} \langle k' | V | k \rangle &= \frac{2\pi\hbar^2}{m} \sum_l b_l \int d^3r e^{-ik'\cdot r} \delta(\mathbf{r} - \mathbf{R}_l) e^{ik\cdot r} \\ &= \frac{2\pi\hbar^2}{m} \sum_l b_l e^{i\mathbf{Q}\cdot \mathbf{R}_l} \end{aligned}$$

The scattering Cross section

$$\frac{d\sigma}{d\Omega} = \left(\frac{m}{2\pi\hbar^2} \right)^2 \left| \langle k' | V | k \rangle \right|^2$$

If the scattering is not dependent on the spin states and is dependent only on random distribution of isotopes,

$$\begin{aligned} \frac{d\sigma}{d\Omega} &= \sum_{l,l'} e^{i\mathbf{Q}\cdot(\mathbf{R}_l - \mathbf{R}_{l'})} \left| \langle b_l^* b_{l'} \rangle \right|^2 \\ \langle b_l^* b_{l'} \rangle &= (1 - \delta_{ll'}) \bar{b}^2 + \delta_{ll'} \bar{b}^2 \end{aligned}$$

Hence,

$$\frac{d\sigma}{d\Omega} = N \left(\bar{b}^2 - \bar{b}^2 \right) + \bar{b}^2 \sum_{l,l'} e^{i\mathbf{Q}\cdot(\mathbf{R}_l - \mathbf{R}_{l'})}$$

incoherent coherent

Now, we need to take care of the change of spin states,

$$s \rightarrow s'$$

$$\frac{d\sigma}{d\Omega} = \sum_{s,s'} P_s \sum_{l,l'} e^{i\mathbf{Q}\cdot(\mathbf{R}_l - \mathbf{R}_{l'})} \left| \langle s' | b_l^* b_{l'} | s \rangle \right|^2$$



11. Scattering Length

Neutron Interaction Potentials

Nuclear Interaction
(Neutron-Nucleus)

$$V_N(\mathbf{r}) = \frac{2\pi\hbar^2}{m_n} b_N \delta(\mathbf{r})$$

Magnetic Interaction
(Neutron-Unpaired Electron)

$$V_M(\mathbf{r}) = -\boldsymbol{\mu} \cdot \mathbf{B}(\mathbf{r})$$

Magnetic moment of neutron
B-field induced unpaired spin

Scattering length, b

$$b = \frac{m_n}{2\pi\hbar^2} V(\mathbf{Q})$$

Fourier Transform of $V(\mathbf{r})$

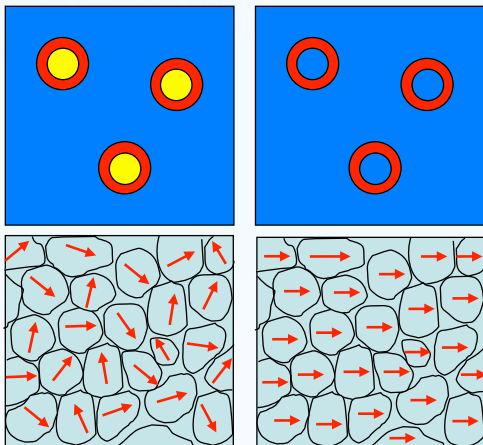
$$b = b_N + b_M = b_N + \gamma_e \boldsymbol{\sigma} \cdot \mathbf{S}_{\perp} f(Q)$$

Pauli operator for neutron
Magnetic form factor

Nuclear Magnetic
Spin component₄₁ perpendicular to Q



12. Scattering Length Density

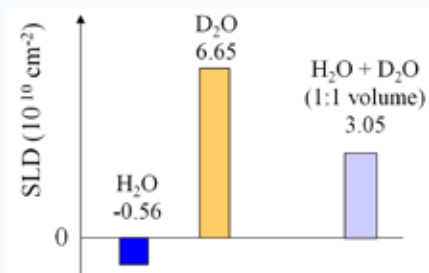


Scattering length density, ρ

$$\rho = \frac{\sum_j^n b_j}{\bar{V}}$$

b_j = bound coherent scattering length of atom j

\bar{V} = volume containing the n atoms



Contrast variation

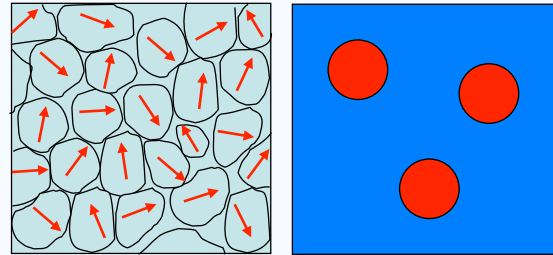
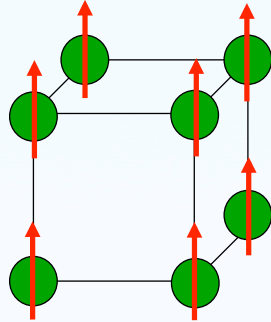
- bound coherent scattering length (10^{-13} cm)
 $b_H = -3.749 \text{ fm}$ $b_D = 6.671 \text{ fm}$



What Contrast Do We Use in Neutron Scattering ?

□ Thermal Neutron Diffraction

□ SANS/REF



- Use atomic properties
: **Scattering length**

- Use material properties
: **Scattering length density**

43



Part 2

13. Calculation of scattering lengths

<http://www.ncnr.nist.gov/resources/n-lengths/>

$$b \equiv b_{molecule} = \sum_i r_i b_{atom,i}$$

Ex. benzene C_6H_6

$$\begin{aligned} b_{benzene} &= 6b_H + 6b_C \\ &= 6 \times (-3.739 \times 10^{-13}) + 6 \times (6.646 \times 10^{-13}) \\ &= 17.442 \times 10^{-13} [\text{cm}] \end{aligned}$$

Isotope	conc	Coh b	Inc b	Coh xs	Inc xs	Scatt xs	Abs xs
	%	fm (=10 ⁻¹³ cm)	fm	barn(=10 ⁻²⁴ cm ²)	barn	barn	barn
isotope	Conc.	Coh. Scatt. length	Inc. scatt. length	Coh. Cross section	Inc. cross section	Scattering cross section	Absorption cross section
H	---	-3.739	---	1.7568	80.26	82.02	0.3326
¹ H	99.985	-3.7406	25.274	1.7583	80.27	82.03	0.3326
² H	0.015	6.671	4.04	5.592	2.05	7.64	0.000519
C	---	6.646	---	5.551	0.001	5.551	0.0035
N	---	9.36	---	11.01	0.5	11.51	1.9
O	---	5.803	---	4.232	0.0008	4.232	0.00019

b

σ_{coh}

σ_{inc}

σ_s

σ_a

Q: Calculate the scattering lengths of light (H₂O) and heavy (D₂O) waters.



Neutron scattering lengths and cross sections

Isotope	conc	Coh b	Inc b	Coh xs	Inc xs	Scatt xs	Abs xs
H	---	-3.7390	---	1.7568	80.26	82.02	0.3326
1H	99.985	-3.7406	25.274	1.7583	80.27	82.03	0.3326
2H	0.015	6.671	4.04	5.592	2.05	7.64	0.000519
3H	(12.32 a)	4.792	-1.04	2.89	0.14	3.03	0

Column	Unit	Quantity
1	---	isotope
2	---	Natural abundance (For radioisotopes the half-life is given instead)
3	fm	bound coherent scattering length
4	fm	bound incoherent scattering length
5	barn	bound coherent scattering cross section
6	barn	bound incoherent scattering cross section
7	barn	total bound scattering cross section
8	barn	absorption cross section for 2200 m/s neutrons

NOTE: The above are only thermal neutron cross sections. For dependent cross sections please go to the National Nuclear Energy Research Institute website.

Select the element, and you will get a list of scattering lengths and cross sections. Feature section of neutron scattering lengths and cross sections, No. 3, 1992, pp. 29-37.

The scattering lengths and cross sections only go through the element.

Note: 1fm=1E-15 m, 1barn=1E-24 cm², scattering lengths and cross sections in parenthesis are uncertainties. A long table with the complete list of elements and isotopes is also available.



14. Neutron Scattering : Fourier Transform

□ Differential scattering cross-section

$$\frac{d\sigma}{d\Omega}(\vec{Q}) = \left\langle \left| \sum_j b_j e^{-i\vec{Q}\cdot\vec{R}_j} \right|^2 \right\rangle$$

Dirac delta function

$$\int \delta(\vec{r}) d\vec{r} = 1$$

$$\int f(\vec{r}) \delta(\vec{r} - \vec{R}) d\vec{r} = f(\vec{R})$$

$$n(\vec{r}) = \sum_j \delta(\vec{r} - \vec{R}_j) : \text{Atomic number density}$$

$$\rho_{sld}(\vec{r}) = \sum_j b_j \delta(\vec{r} - \vec{R}_j) : \text{Scattering length density}$$

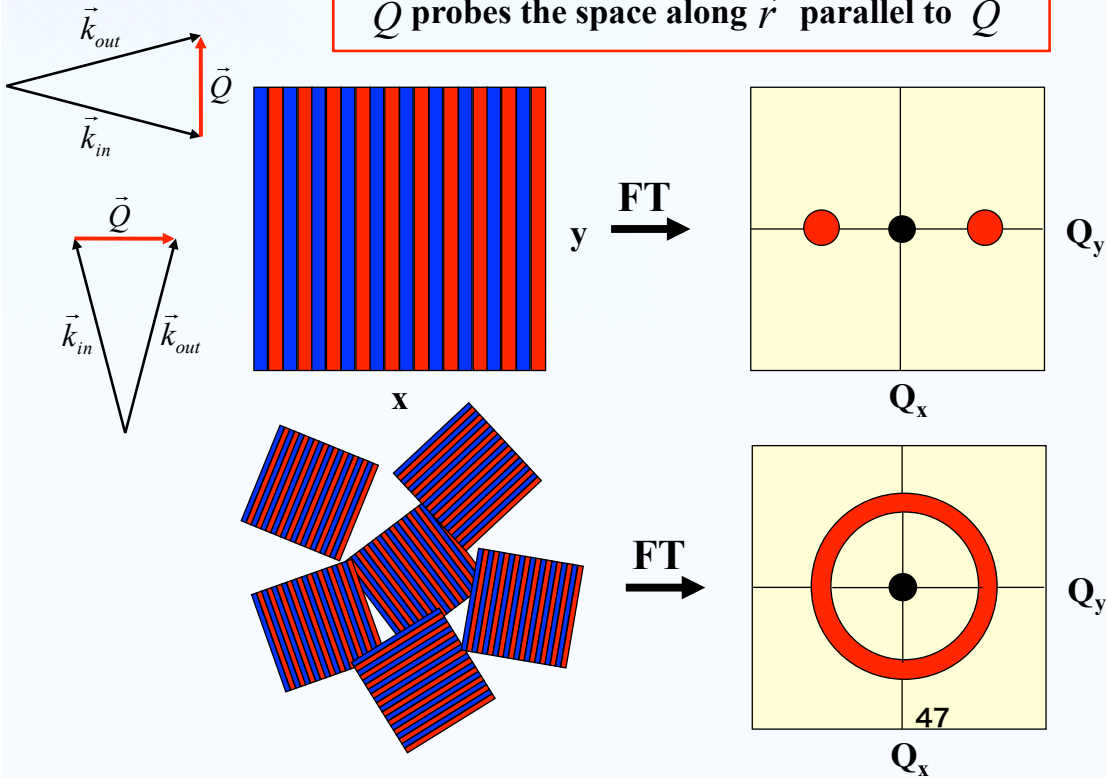
$$\text{F.T.} \{ \rho_{sld}(\vec{r}) \} = \int \rho_{sld}(\vec{r}) e^{-i\vec{Q}\cdot\vec{r}} d\vec{r} = \int \sum_j b_j \delta(\vec{r} - \vec{R}_j) e^{-i\vec{Q}\cdot\vec{r}} d\vec{r} = \sum_j b_j e^{-i\vec{Q}\cdot\vec{R}_j}$$

$$\frac{d\sigma}{d\Omega}(\vec{Q}) = \left\langle \left| \int \rho_{sld}(\vec{r}) e^{-i\vec{Q}\cdot\vec{r}} d\vec{r} \right|^2 \right\rangle$$



15. Fourier Transform

\vec{Q} probes the space along \vec{r} parallel to \vec{Q}



Part 3 16. From scattering cross section to absolute intensity

Total cross section: the sum of scattering cross section σ_s and absorption cross section σ_a .
 The scattering cross section: **coherent scattering cross section σ_{coh}** and **incoherent scattering cross section σ_{inc}** .

$$\sigma_{tot} = \sigma_s + \sigma_a$$

$$\sigma_s = \sigma_{coh} + \sigma_{inc}$$

Scattering intensity (differential scattering cross section)

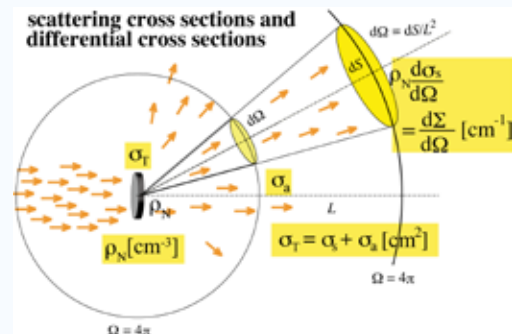
The differential scattering cross section is given by normalizing the scattering cross section by the number density of the scatterers in the sample.

$$\frac{d\sigma}{d\Omega} [cm^2] \xrightarrow{\times N/V} \frac{N}{V} \frac{d\sigma}{d\Omega} [cm^{-1}] = \rho_N \frac{d\sigma}{d\Omega} [cm^{-1}] = \frac{d\Sigma}{d\Omega} [cm^{-1}]$$

(one particle)

$$\rho_N = \frac{N_A d [g/cm^3]}{m [g/mol]} [cm^{-3}] \quad \text{Number density}$$

$\frac{d\Sigma}{d\Omega} (Q) [cm^{-1}]$: commonly-used unit for scattering intensity



$$\frac{d\Sigma}{d\Omega} (Q) [cm^{-1}] \xrightarrow{\int_{\Omega}} \Sigma [cm^{-1}] \quad \text{: total cross section of the sample}$$



Summary:

Comparison of scattering amplitudes

neutron

$$f^{(1)}(\theta) = -b$$

constant

X-ray

$$f^{(1)}(\theta) = \frac{e^2}{m_e c^2} A(\theta), \quad A(\theta = 0) = Z$$

angular dependent

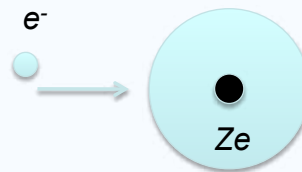
electron

$$f^{(1)}(\theta) = \frac{2m_e e^2}{\hbar^2} \frac{Z - A(\theta)}{Q^2}, \quad A(\theta = 0) = Z$$

Two contributions; from the nucleus, and the electron clouds

$$V(r) = -\frac{Ze^2}{r} + \int d^3r' \frac{(-e)(-e)^2 \rho(r')}{|r-r'|}$$

Atom and Incoming electron Electron cloud and incoming electron



$$\int d^3r \rho(r) \equiv Z$$

49



Part 2

Summary: scattering variables

Total cross section

$$\sigma_{tot} = \int \sigma(\theta, \varphi) d\Omega$$

Differential scattering cross section

$$\sigma = |f_k(\theta, \varphi)|^2$$

Fermi's pseudo potential

$$V(r) = \frac{2\pi\hbar^2}{m} b \delta(\mathbf{r} - \mathbf{R})$$

Scattering length

$$f(\theta, \varphi) = -\frac{1}{4\pi} \int d^3r' e^{-i\mathbf{Q}\cdot\mathbf{r}'} U(r') = -b$$

Scattering vector

$$\mathbf{Q} \equiv \mathbf{k}_{out} - \mathbf{k}_{in} = \mathbf{k}' - \mathbf{k}$$

Scattered wave

$$\psi(r) = -\frac{b}{r} e^{ikr}$$

Scattering cross section for a particle

$$\frac{d\sigma}{d\Omega} = b^2$$

Coherent scattering cross section

$$\sigma_{coh} \equiv 4\pi \bar{b}^2$$

Incoherent scattering cross section

$$\sigma_{inc} \equiv 4\pi (\overline{b^2} - \bar{b}^2)$$

50



“Neutron and Rheology”

Mitsuhiro Shibayama,
ISSP, U. Tokyo

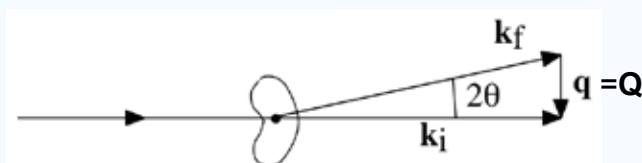
Part 1. Structural Analyses of Polymers by Small Angle Neutron Scattering

1. Introduction
2. Neutron and neutron scattering
3. Neutron scattering theory
4. **Small angle Scattering**
5. Applications



Part 4

1. what is small-angle scattering?

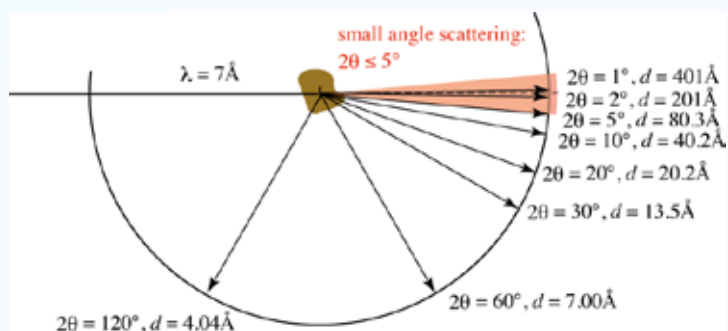


- Constructive interference from structures in the direction of q

- Diffraction length scale

$$d \approx \frac{2\pi}{q}, \quad 2\theta \approx \frac{\lambda}{d} \approx \frac{7 \text{ \AA}}{40 \sim 800 \text{ \AA}}$$

$$2\theta \approx 0.5^\circ \sim 10^\circ$$



- Scattering is at small angles - non-zero but smaller than classical diffraction angles

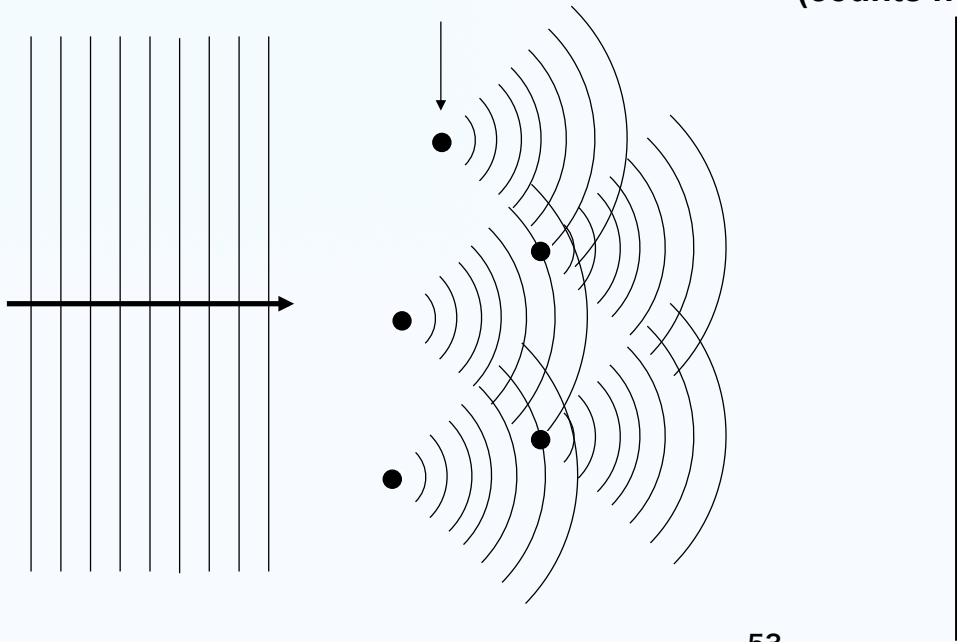


2. Young's Double Slit Experiment (1)

Incident Neutron Wave

Atoms

Detector
(counts neutrons)



53

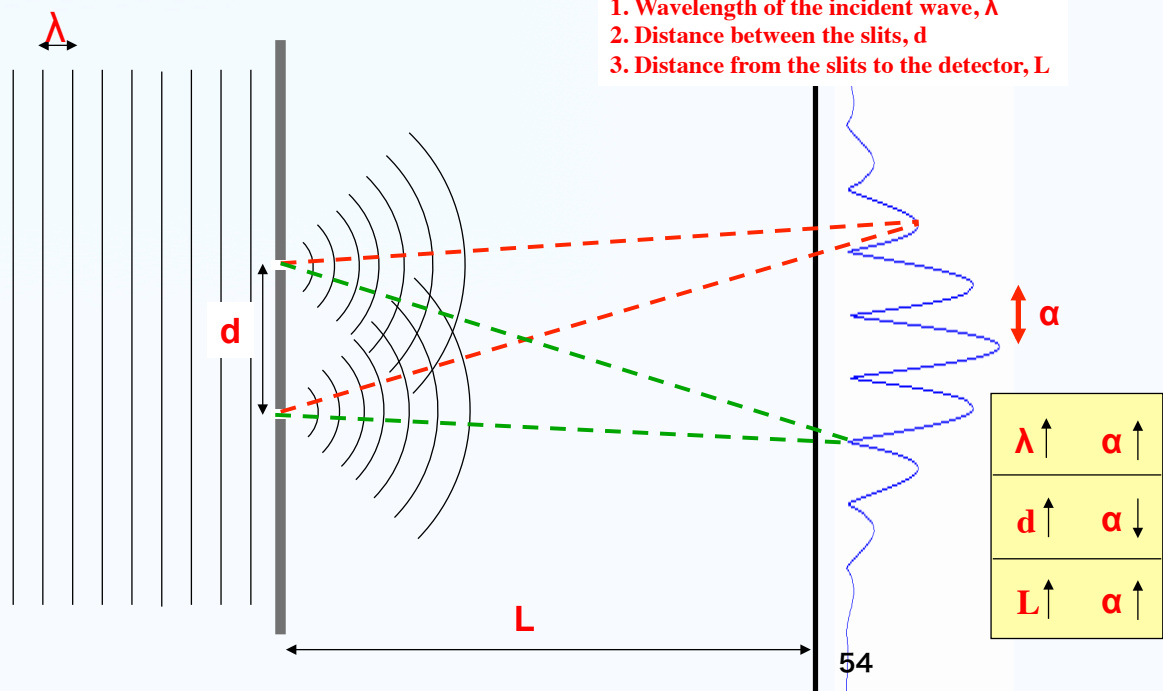


2. Young's Double Slit Experiment (2)

Incoming plane wave

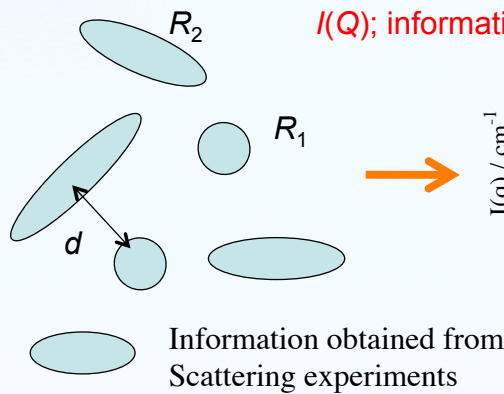
Scattered spherical wave

Interference Pattern

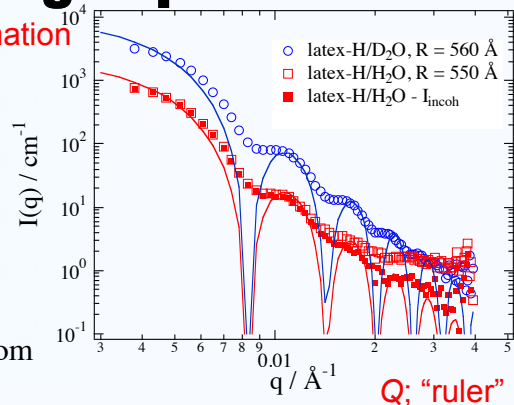


54

3. Information obtained by small-angle scattering experiments

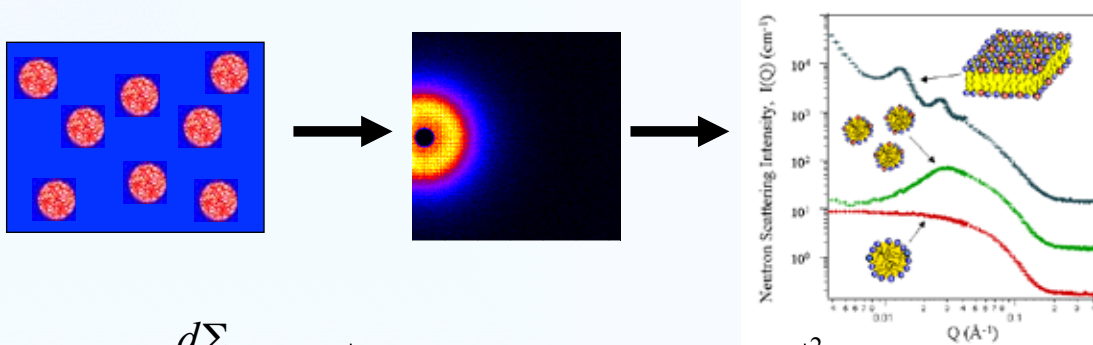


Structural Information
 size, R_1 , R_2 ,
 shape,
 volume fraction, ϕ
 orientation,
 domain distance, d
 fractal dimension, D
 miscibility,
 specific surface, S/V



Ex. SANS function from a polystyrene latex (PS)

4. What Information from SANS ? : Particulate Systems



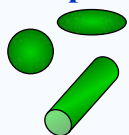
$$\frac{d\Sigma}{d\Omega}(\mathbf{Q}) = |\text{Fourier Transform of } \rho(\mathbf{r})|^2$$

$$= \frac{N_p}{V} |F(Q)|^2 \frac{1}{N_p} \left\langle \sum_i \sum_j e^{i\mathbf{Q} \cdot (\mathbf{r}_i - \mathbf{r}_j)} \right\rangle$$

Number density

$$= n_p P(Q) S(Q)$$

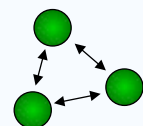
Shape and dimensions of particles



Intra-Particle interference
: Form factor

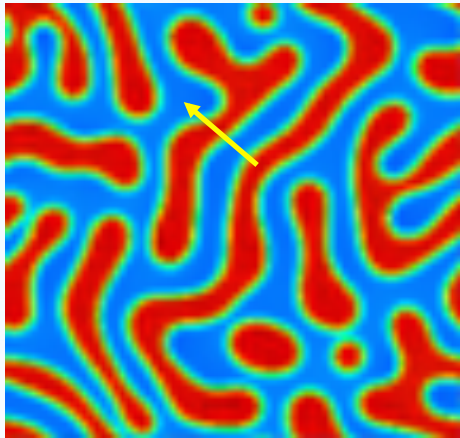
Inter-Particle interference
: Structure factor

Interaction between particles





4. What Information from SANS ? : Non-Particulate Systems



$$\frac{d\Sigma}{d\Omega}(Q) = 4\pi \langle (\Delta\rho)^2 \rangle \int_V \gamma(r) \frac{\sin(Qr)}{Qr} r^2 dr$$

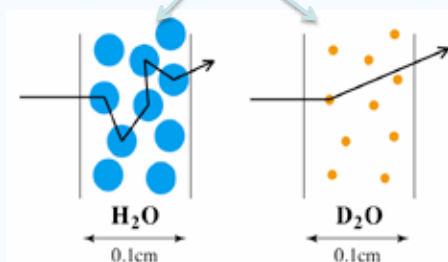
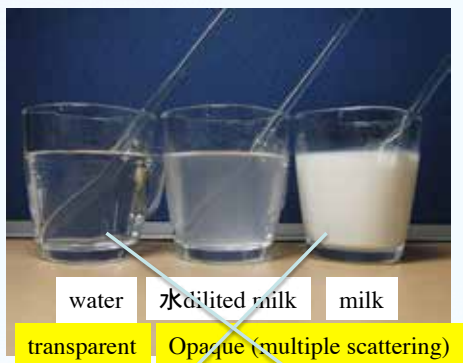
Contrast
Correlation function
Orientation average

$$\gamma(r) = \frac{\int \langle \Delta\rho(\mathbf{r}') \Delta\rho(\mathbf{r}'+\mathbf{r}) \rangle d\mathbf{r}'}{\int \langle \Delta\rho(\mathbf{r}') \Delta\rho(\mathbf{r}') \rangle d\mathbf{r}'}$$

57

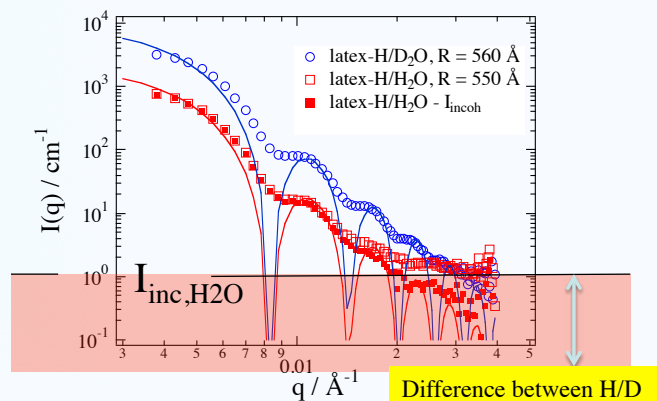


5. Multiple scattering by H



MFP \approx 0.15cm MFP \approx 1.6cm

$$\text{mean free path (MFP)} = \frac{1}{\rho\alpha_T} = \frac{1}{\Sigma_T} [\text{cm}] = \frac{1}{\text{total cross section}}$$



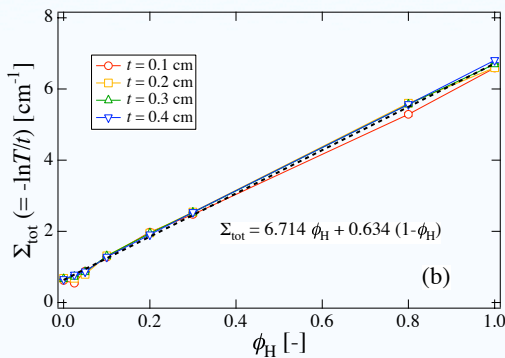
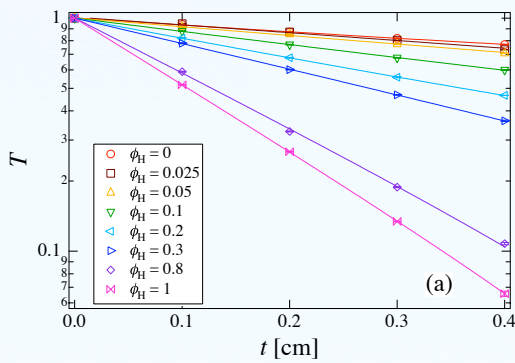
Ex. Polystyrene latex

If incoherent scattering is large, Information at large q cannot be obtained.

Need to design to lower incoherent scattering

58

6. Transmission of H₂O/D₂O



mean free path (MFP) = $\frac{1}{\rho\alpha_T} = \frac{1}{\Sigma_T}$ [cm] = $\frac{1}{\text{total cross section}}$

$$T \equiv \frac{I_{tr}}{I_0} = \exp(-\Sigma_{tot} t)$$

The Lambert-Beer law holds even $T \approx 0.05$.

$$\Sigma_{tot} = 6.714\phi_H + 0.634(1 - \phi_H)$$

A linear relationship holds all fractions irrespective of sample thickness.

Shibayama, et al., J. Appl. Crystallogr. 2009, 42, 621.

7. Scattering function from a polymer chain



the radius of gyration

$$R_G = \frac{1}{N} \sum_{n=1}^N \mathbf{R}_n$$

$$R_g^2 = \frac{1}{N} \sum_{n=1}^N \langle (\mathbf{R}_n - \mathbf{R}_G)^2 \rangle$$

$$R_g^2 = \frac{1}{2N^2} \sum_{n=1}^N \sum_{m=1}^N \langle (\mathbf{R}_m - \mathbf{R}_n)^2 \rangle$$

N ; the degree of polymerization

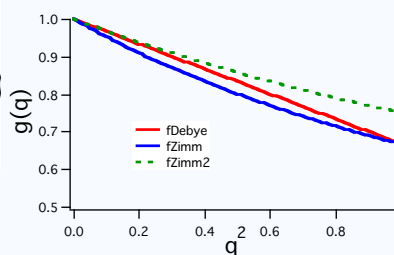
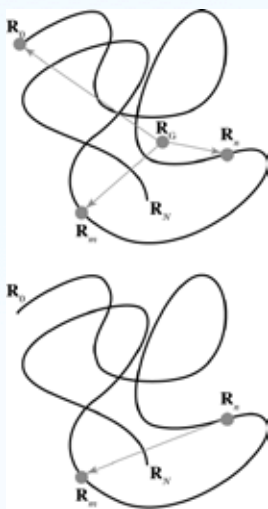
the Debye fn.

$$g(\mathbf{r}) = \frac{1}{N} \sum_n \sum_m \langle \exp[i\mathbf{Q} \cdot (\mathbf{R}_m - \mathbf{R}_n)] \rangle$$

$$= \frac{1}{N} \sum_n \sum_m \exp\left[-\frac{|m-n|}{6} b^2 Q^2\right]$$

$$= N f_D \left((QR_g)^2 \right) \quad b; \text{ the segment length}$$

$$f_D \left((QR_g)^2 \right) \equiv f_D(x) = \frac{2}{x^2} (e^{-x} - 1 + x)$$



$$g(x) = \frac{2N}{x^2} (e^{-x} - 1 + x), \quad x \equiv R_g^2 Q^2$$

$$R_g = \frac{N^{1/2} b}{\sqrt{6}}$$

Q: Discuss the asymptotic behavior of the Debye function near $x=0$ and $x=\text{large}$.

Summary:

correlation functions and scattering intensity for various systems

$$g(r) \Leftrightarrow I(Q) = \int g(r) \exp(i\mathbf{Q} \cdot \mathbf{r}) d\mathbf{r} = \int g(r) \frac{\sin Qr}{Qr} 4\pi r^2 dr$$

corr. fns.

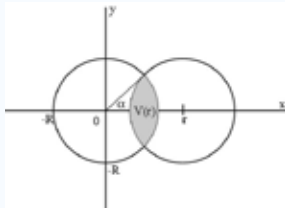
scatt. fns.

Gaussian fn. (scattering from an assembly of non-interacting particles)

$$g(r) = \exp\left[-\frac{r^2}{2(R_g^2/3)}\right] \quad \longleftrightarrow \quad I(Q) \sim \exp\left(-\frac{R_g^2}{3} Q^2\right)$$

Scattering from isolated particles

$$g(r) = \frac{V(r)}{4\pi R^3/3} = 1 - \frac{3}{4} \frac{r}{R} + \frac{1}{16} \left(\frac{r}{R}\right)^3 \quad \longleftrightarrow \quad I(Q) \sim \Phi^2(QR) = \left\{ \frac{3[\sin(QR) - QR \cos(QR)]}{(QR)^3} \right\}^2$$



Lorentz fn. (semi-dilute polymer solutions)

$$g(r) = \frac{1}{r} \exp\left(-\frac{r}{\xi}\right) \quad \longleftrightarrow \quad I(Q) \sim \frac{1}{1 + Q^2 \xi^2}$$

Polymer chains

The volume $V(r)$ of the shaded part

$$g(r) = \frac{1}{N} \sum_{n=1}^N \sum_{m=1}^N \langle \delta\{\mathbf{r} - (\mathbf{R}_m - \mathbf{R}_n)\} \rangle \quad \longleftrightarrow \quad f_D(x) = \frac{2}{x^2} (e^{-x} - 1 + x) \quad x = (QR_g)^2$$

Summary: Scattering functions

Differential
Cross section

Scattering is a Fourier transform of
the contrast variation square

$$\frac{d\Sigma}{d\Omega}(q) = |\text{Fourier Transform of } \rho(\mathbf{r})|^2 = (\Delta\rho)^2 \times$$

$$\left\{ \begin{array}{l} n_p P(q) S(q) = \begin{cases} \frac{n_p P(q)}{S(q)} & \text{Particle dispersion (dilute) } (S(q) = 1) \\ n_p P(q) S(q) & \text{Particle dispersion (semi-dilute)} \end{cases} \\ S(q) & \text{Non-particulate system} \\ S(q) = \frac{S(0)}{1 + \xi^2 q^2} & \text{Concentration fluctuations (one phase)} \end{array} \right.$$

$\Delta\rho$ Scattering contrast

$P(q)$

form factor

n_p Number density
of the particles

$S(q)$

structure factor

ξ

correlation length



“Neutron and Rheology”

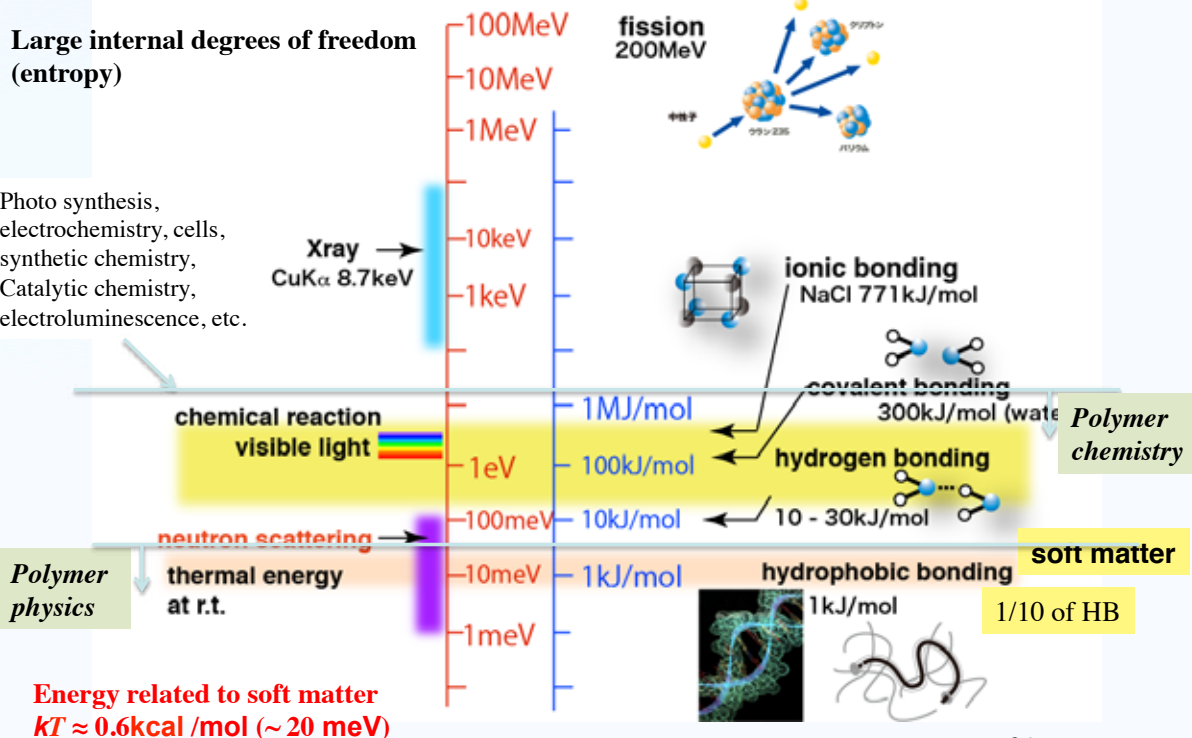
Mitsuhiro Shibayama,
 ISSP, U. Tokyo

Part 1. Structural Analyses of Polymers by Small Angle Neutron Scattering

1. Introduction
2. Neutron and neutron scattering
3. Neutron scattering theory
4. Small angle Scattering
5. Applications

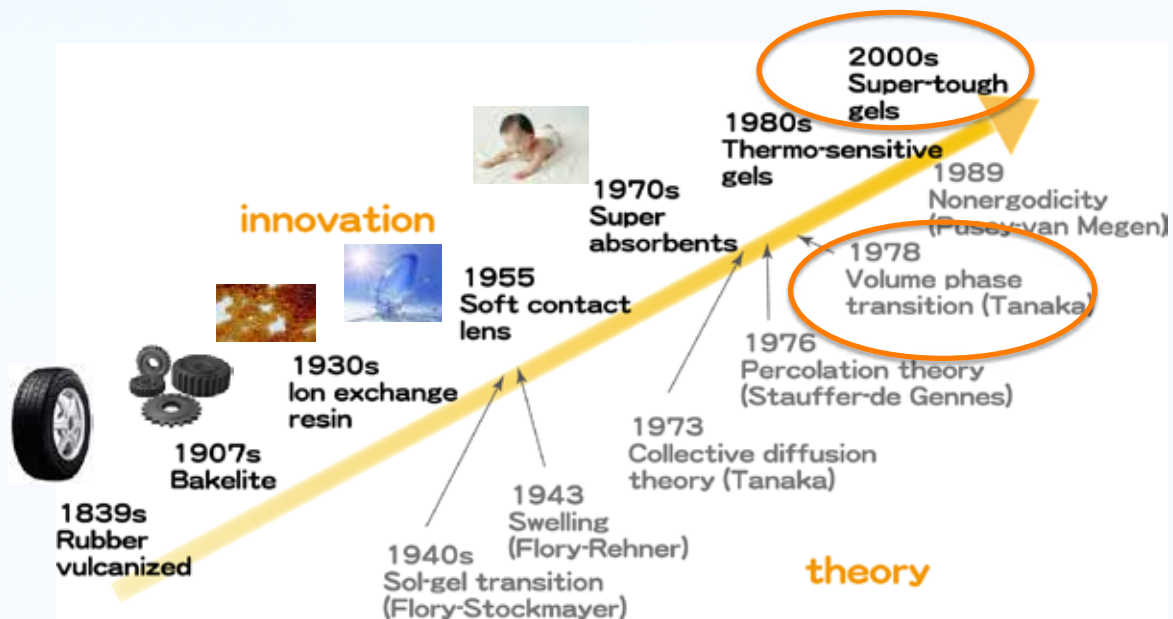


Energies governing soft matter dynamics





History of network polymer and gels



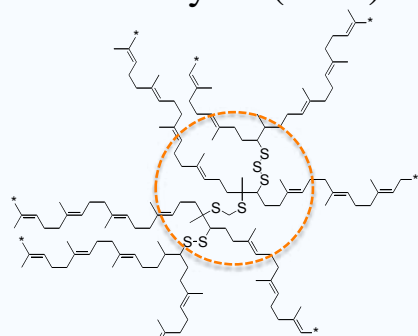
65



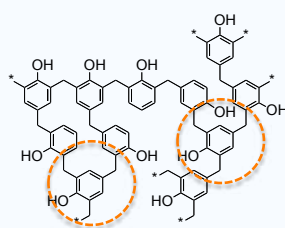
Historical innovations in network polymers



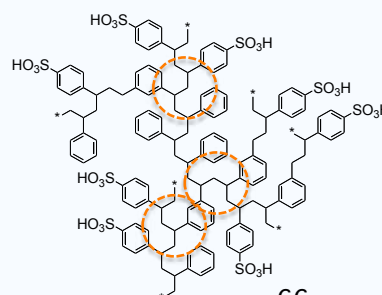
Rubber, Ebonite,
Goodyear (1839)



Bakelite
Dr. Leo Bakeland (1907) (1935)



Ion exchange resin



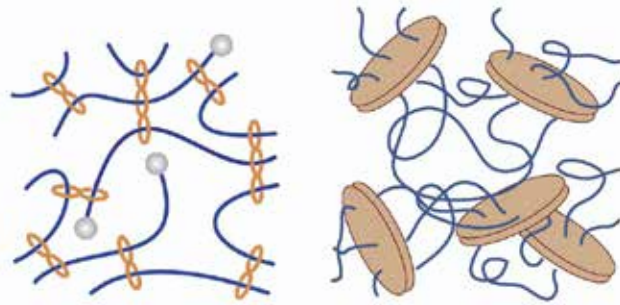
Still keep their significance as industrial materials since their discovery.

66



Tough polymer gels

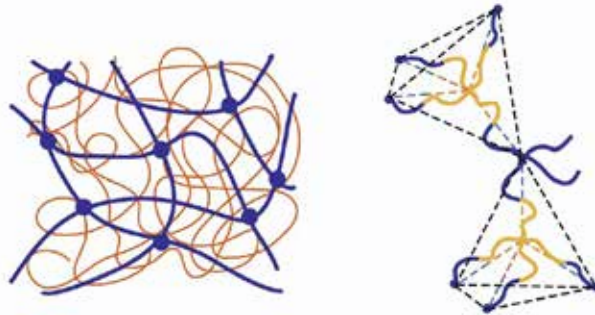
Itoh
(2001)



(a) slide ring (SR) gel (b) nanocomposite (NC) gel

Haraguchi
(2002)

Gong
(2003)



(c) double network (DN) gel (d) tetra-PEG gel

Sakai
(2008)



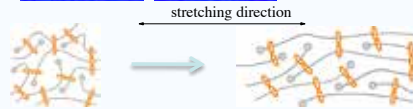
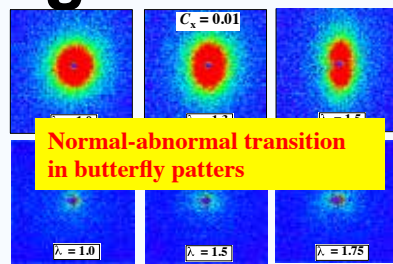
Super tough gels



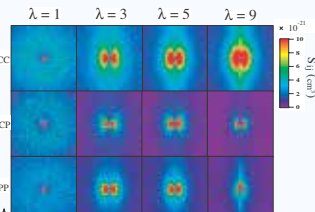
Slide-ring gel Okumura, Adv. Mater., 2001



Nanocomposite gel Haraguchi, Adv. Mater., 2002

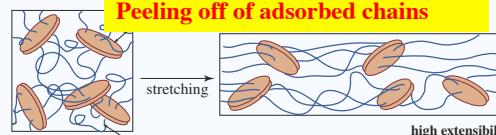


Karino et al., Macromolecules, 2005.



(a) NC gel

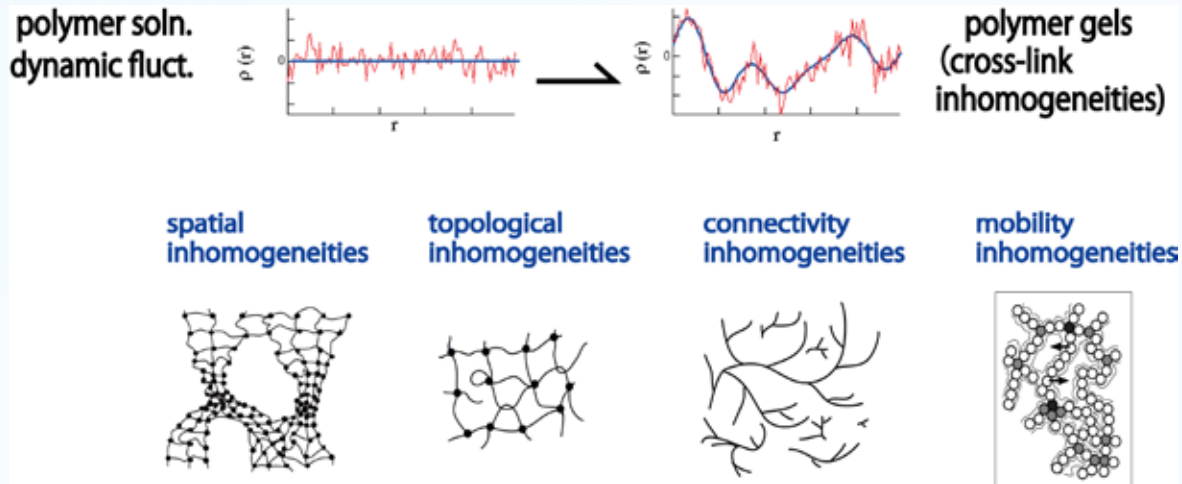
CV-SANS;
Peeling off of adsorbed chains



large $\langle R^2 \rangle^{1/2}$ high extensibility
68
Shibayama et al., Macromolecules, 2005,
Nishida et al., PRE, 2009.



Various types of inhomogeneities in polymer gels

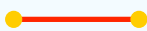


M. Shibayama, Bull. Chem. Soc. Jpn. 2002, 75, 641-659.

69

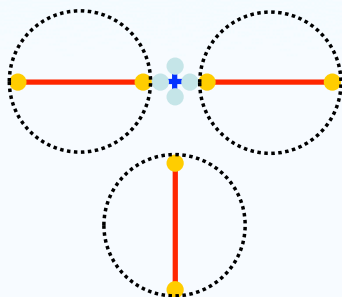


Model Network



• Buried reaction point

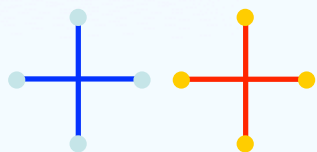
• Self-biting



 Inhomogeneous structure

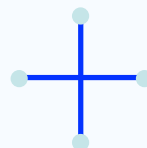
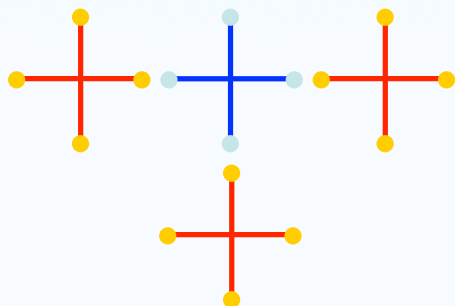


Tetra Network



- Reaction front is always on the surface

- No self-biting

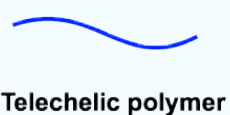


Homogeneous structure



Exploration of realization of ideal-polymer network

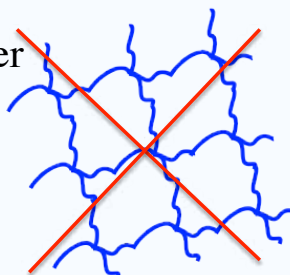
Cross-linking of telechelic polymer



+



multi-arm crosslinker

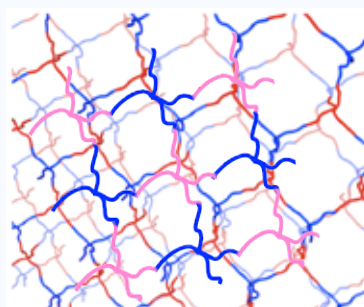


Cross-end coupling



Symmetric multi-functional polymers

+



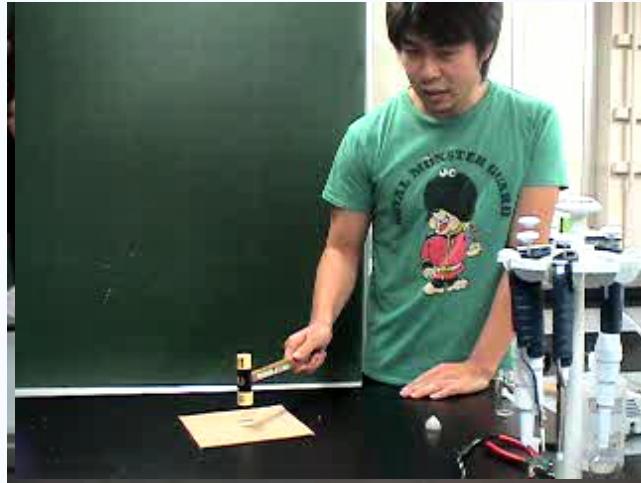
New paradigm of Gel formation
Module assembling by "cross-end-coupling"



Tetra-PEG gel

Advanced physical, chemical, and biological properties

1. high compressive toughness (superior to cartilage)
 2. high transparency
 3. biocompatible and nontoxic
 4. easy and quick preparation
- Etc.



3x speed
~ 1 min

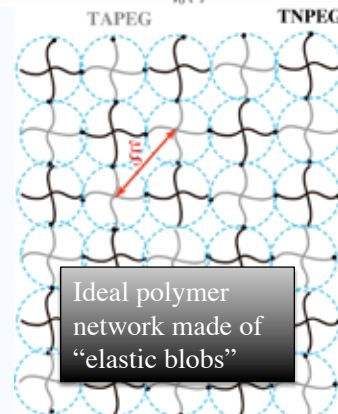
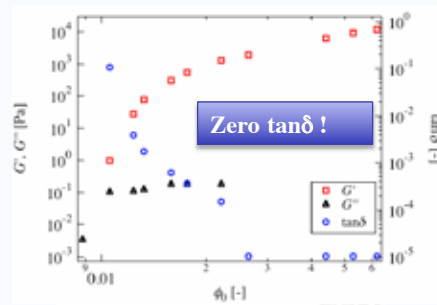


Remarkable properties of Tetra-PEG gel

Tetra-PEG ball vs power ball



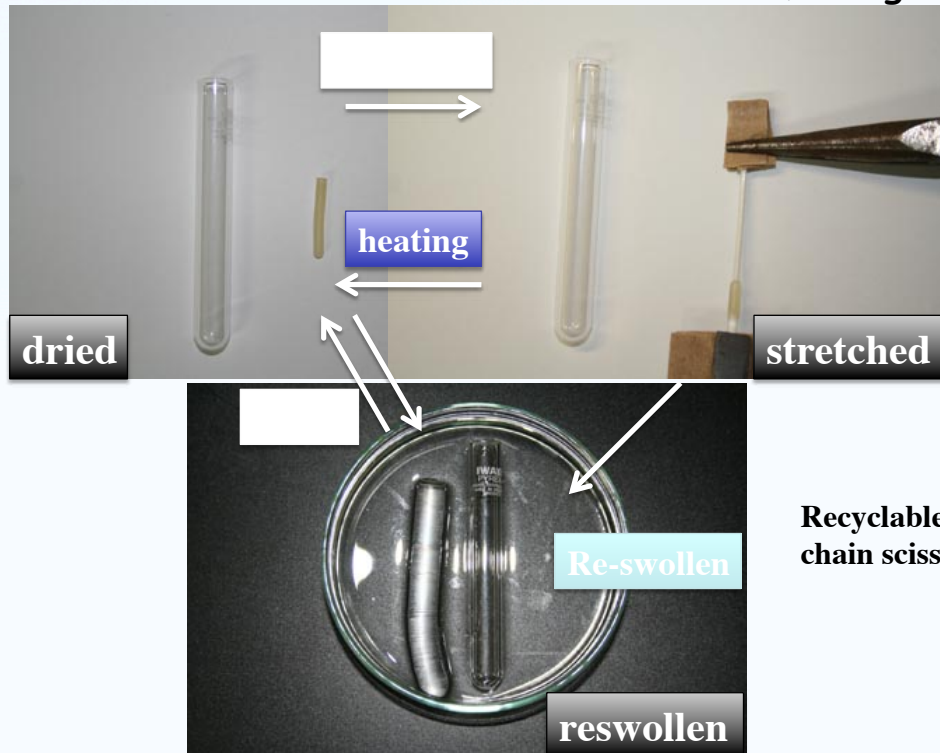
the reflection coefficient ≈ 0.84





Tetra-PEG gel: dried, stretched, re-swollen

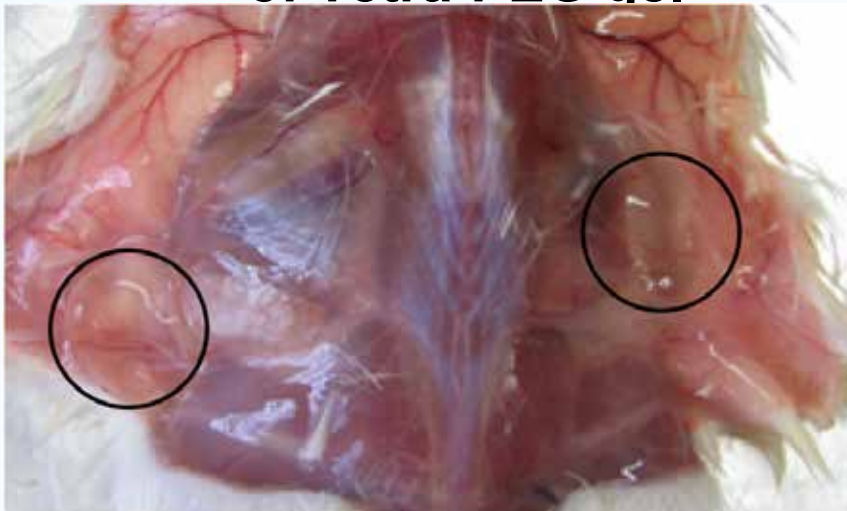
10,000 g/mol



77



subcutaneous implantation of Tetra-PEG gel

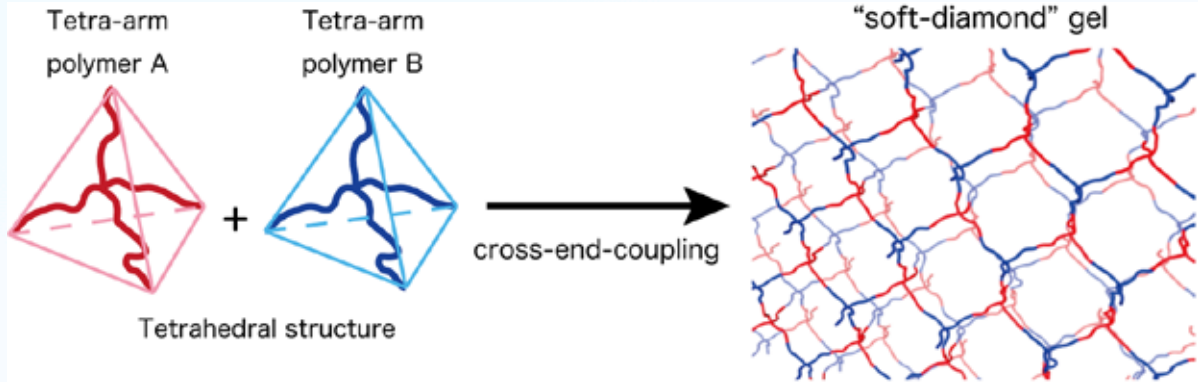


The back of immunocompetent mice
100mL of Tetra-PEG was implanted under anesthesia.
One week after implantation.

78



As-prepared gels



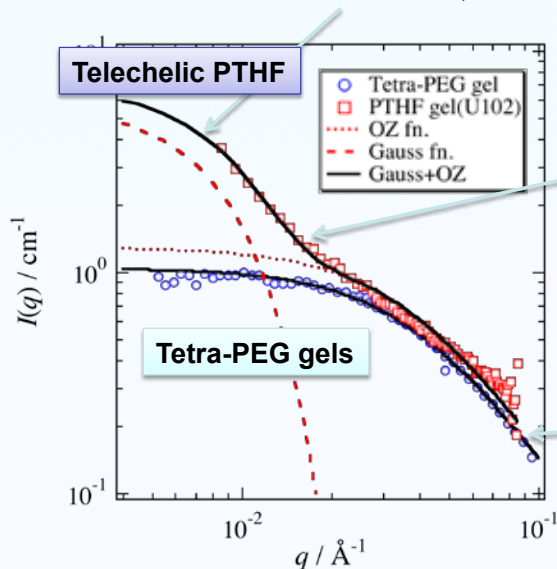
TAPEG	TNPEG	
5k	5k	
10k	10k	
20k	20k	
40k	40k	<i>Symmetric condition</i>

79



Comparison of SANS functions of gels prepared by crosslinking of telechelic polymer vs Tetra-PEG gel

PTHF; made by end-coupling of telechelic PTHF (Takahashi, 1995)



For ordinary gels,

$$I(q) = \frac{I(0)_{SL}}{(1 + \xi^2 q^2)^2} + \frac{I(0)_{OZ}}{1 + \xi^2 q^2}$$

$$I(q) = I_G(0) \exp\left[-\frac{\xi^2 q^2}{3}\right] + \frac{I(0)_{OZ}}{1 + \xi^2 q^2}$$

Gauss **OZ**

OZ function

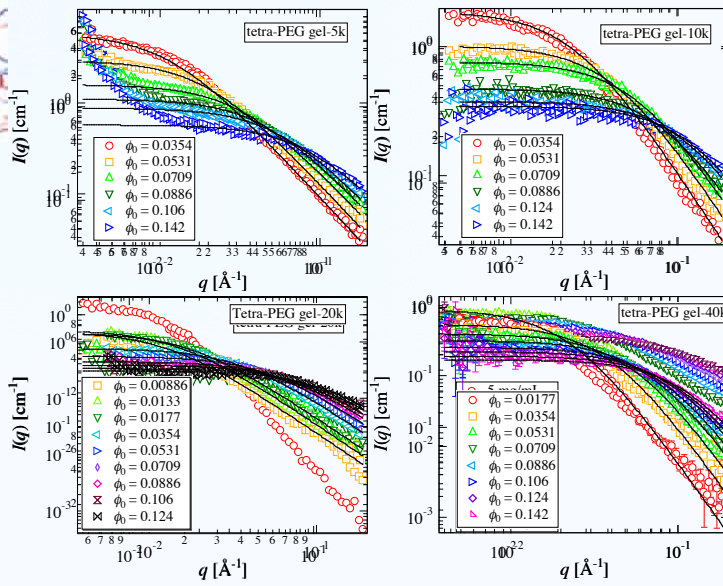
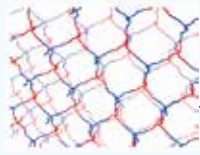
$$I(q) \sim \frac{I(0)_{OZ}}{1 + \xi^2 q^2}$$

Note: Tetra-PEG gels can be described by only OZ fn.

80



Concentration and M_w dependence



As-prepared gels

$$I(q) \sim \frac{I(0)_{OZ}}{1 + \xi^2 q^2}$$

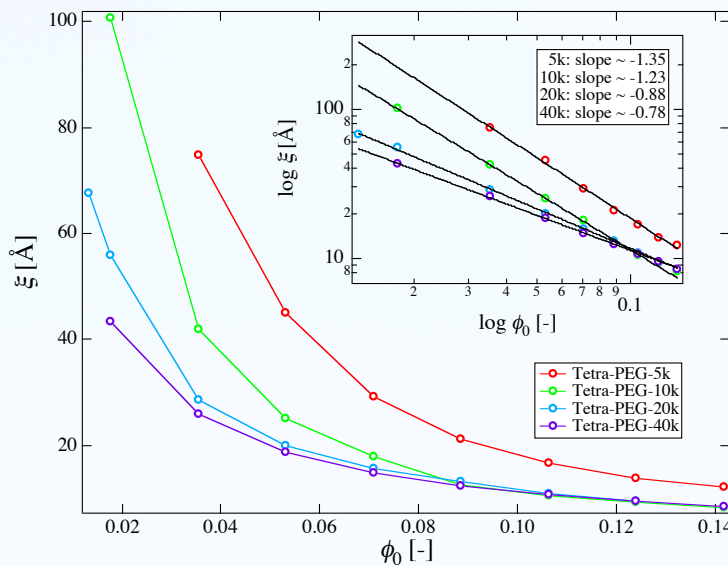
Messages:

1. SANS functions are well fitted by the OZ equation.
2. There appears an upturn at low q 's, particularly for low M_w 's.

81



M_w dependence of ξ for Tetra-PEG gel



$$\xi \sim \phi_0^{-3/4}$$

(Exclusively for -40k gel)

Note: $\xi(M_w=40k) < \xi(M_w=5k)$!!

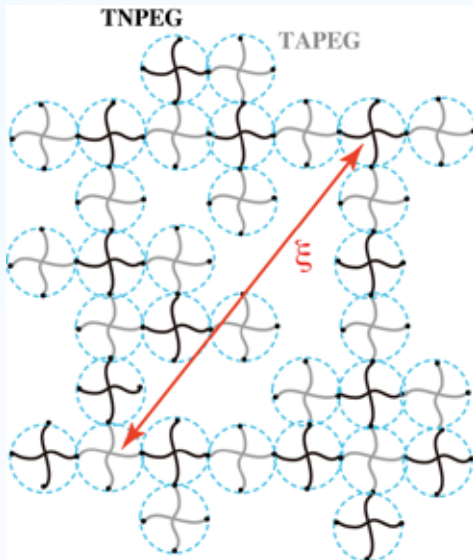
Messages:

Tetra-PEG-40k recovers a proper behavior of polymer solutions In semidilute solutions, while Tetra-PEG-5k is far from it due to its short arm length.

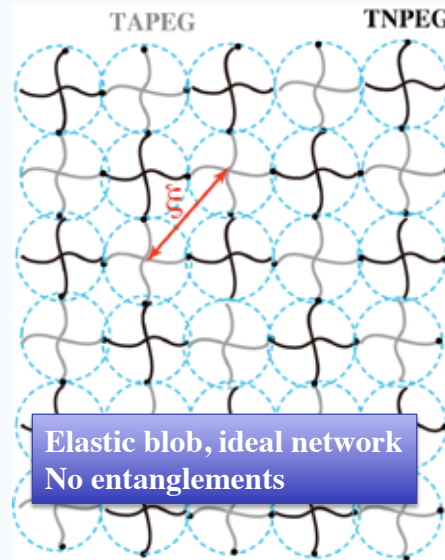
82



Model of inhomogeneities in as-prepared gels



5k-Tetra-PEG gel

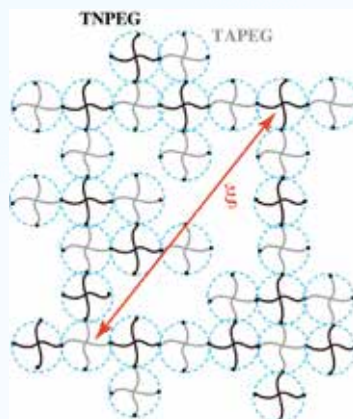


40k-Tetra-PEG gel

83



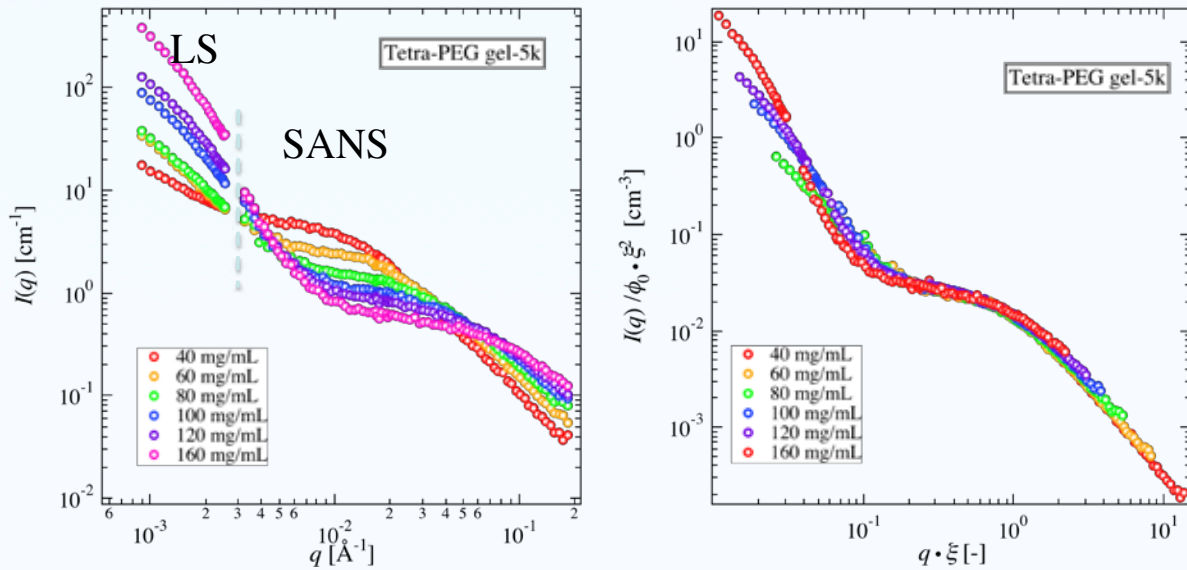
Inhomogeneities in Tetra-PEG gels



84



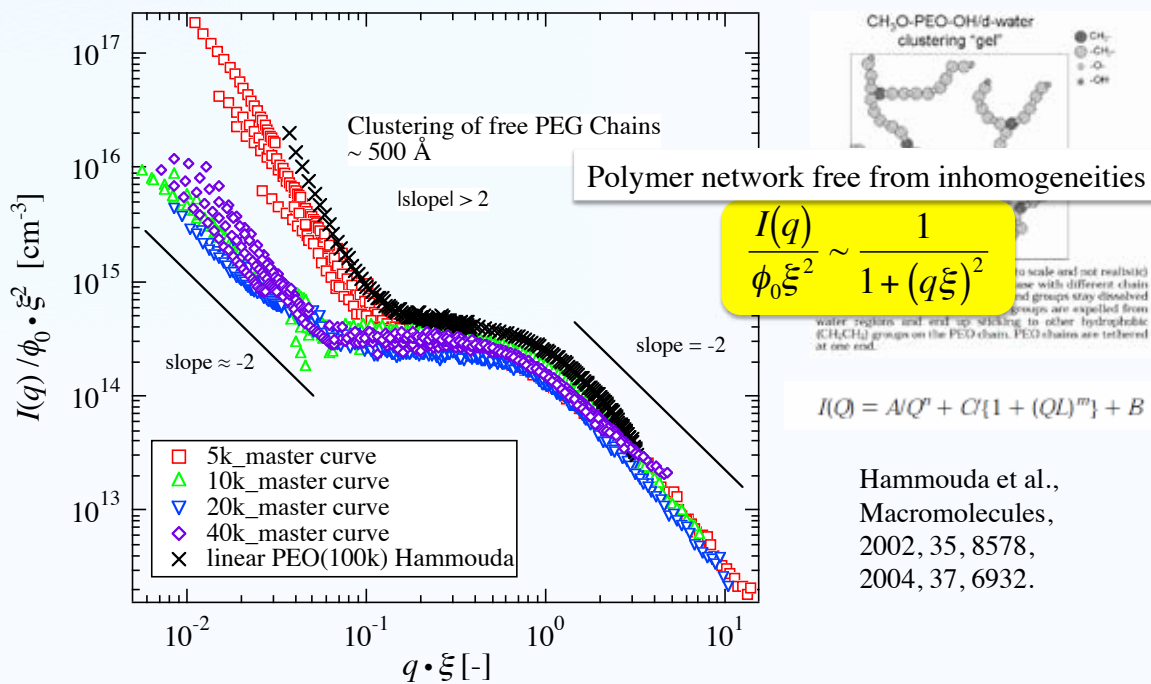
Scattering curves and Master curves of Tetra-PEG gel-5k



85



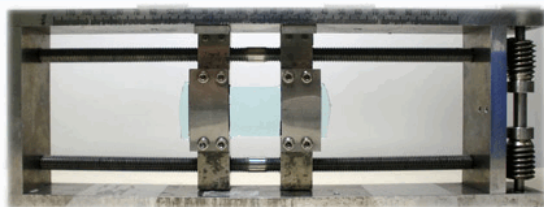
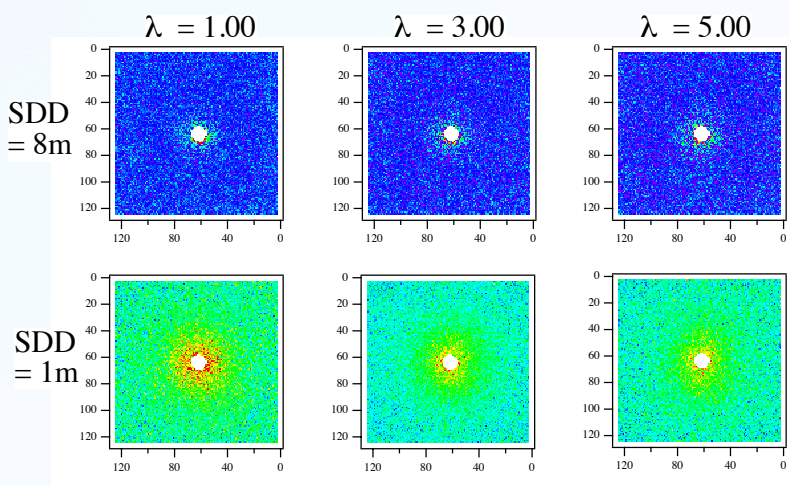
Master curves



86



Deformation SANS for Tetra-PEG



Deformation mechanisms of super-tough gels

M. Shibayama, J. Phys. Soc. Jpn., 78, 041008 (2009)

	Conv. Gel (chem. Gel)	Slide ring gel	Nanocomposite gel	Tetra-PEG gel
Deformation model				
Mech. Prop.	brittle	High elasticity High deformability	High elasticity High deformability	High elasticity High deformability
inhomogeneities	Very large	Low C_x : decreasing High C_x : increasing	small	none
Scatt. pattern	Abnormal	Normal / abnormal liq. like/gel like	Normal liq. like/clay scatt.	isotropic
Charac. feature	Inhomog. str.	Movable cross-links	Plane crosslink, large molec. wt.	Elastic blob