



Multiscale studies of macromolecular systems: concepts and applications

An Introduction

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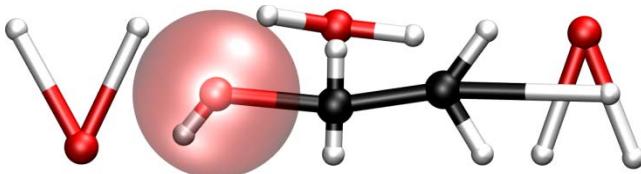


Seoul, June 22 – 26, 2015

Software



Versatile Object-oriented Toolkit for Coarse-graining Applications



votca.org

Apache license (free)

C++, scripting, test suite, hg, wiki pages, bug tracker

votca-csg modules (coarse-graining)
release 1.2
actively developed for 2 years
Google code: votca project

votca-ctp modules (charge transport)
Release planned for August this year
(you can try theta-version)
Google code: votca-ct project



charge transport: V. Ruehle, A. Lukyanov, F. May, M. Schrader, T. Vehoff, J. Kirkpatrick, B. Baumeier, D. Andrienko, submitted, 2011

coarse-graining: V. Ruehle, C. Junghans, A. Lukyanov, K. Kremer, D. Andrienko, J. Chem. Theor. Comp., 5, 3211-3223, 2009

Software



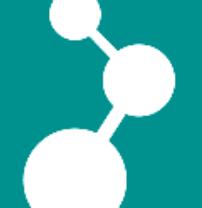
J. D. Halverson, T. Brandes, O. Lenz, A. Arnold, S. Bevc, V. Starchenko, K. Kremer, T. Stuehn, D. Reith, "*ESPResSo++: A Modern Multiscale Simulation Package for Soft Matter Systems*", Computer Physics Communications, 184 (2013), pp. 1129-1149
DOI: 10.1016/j.cpc.2012.12.004
Online access: <http://dx.doi.org/10.1016/j.cpc.2012.12.004>



Outline

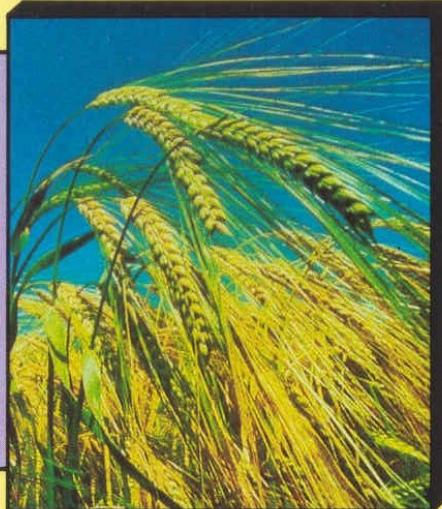
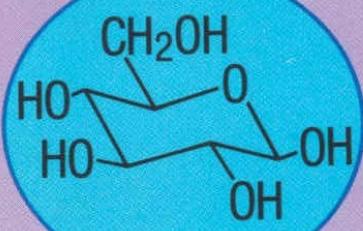
- **Soft and Nanostructured Matter**
- Coarse Graining
 - General Aspects
 - Specific Methods - Examples
- Dynamic Properties
- Adaptive Resolution Simulation: AdResS
- Conclusion/Outlook

Soft Matter - What Systems?

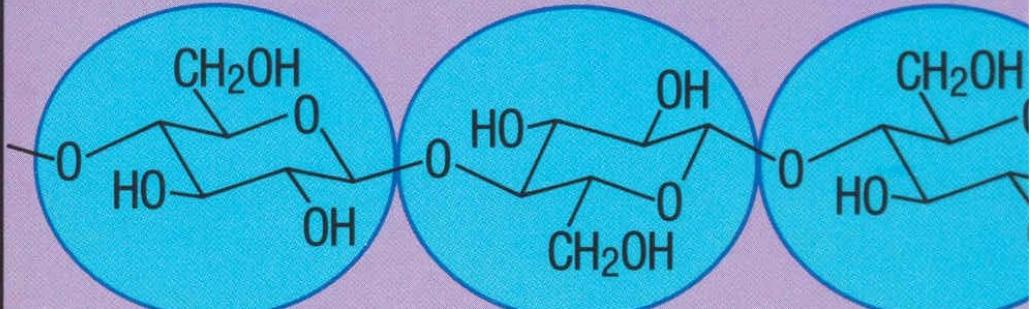


- **Organic Materials:**
 - Polymers
 - Colloids
 - Membranes
 - Hard-Soft Composites (organic-anorganic)
 - Synthetic
 - Biological

natürlich



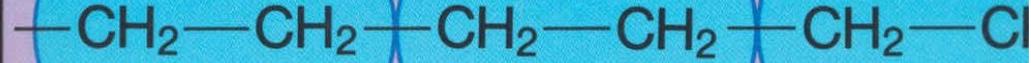
Cellulose



synthetisch



Polyethylen



High Tech Commodity: Polycarbonate (PC)

Bayer Materials



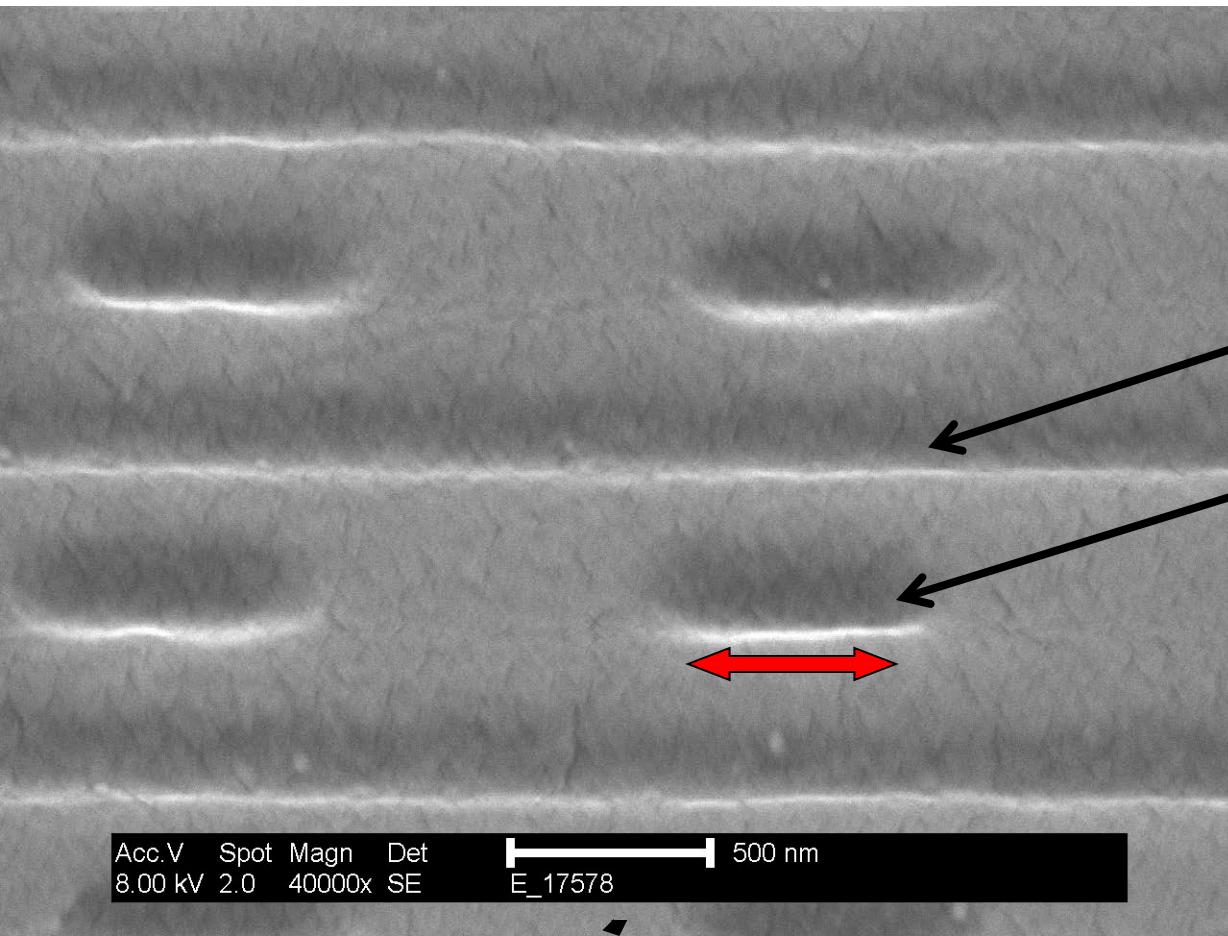
Yearly production more than 500000t/p.a.

Polycarbonate (PC) for

CDs/DVDs...

PC/Ni interface during dye cast process

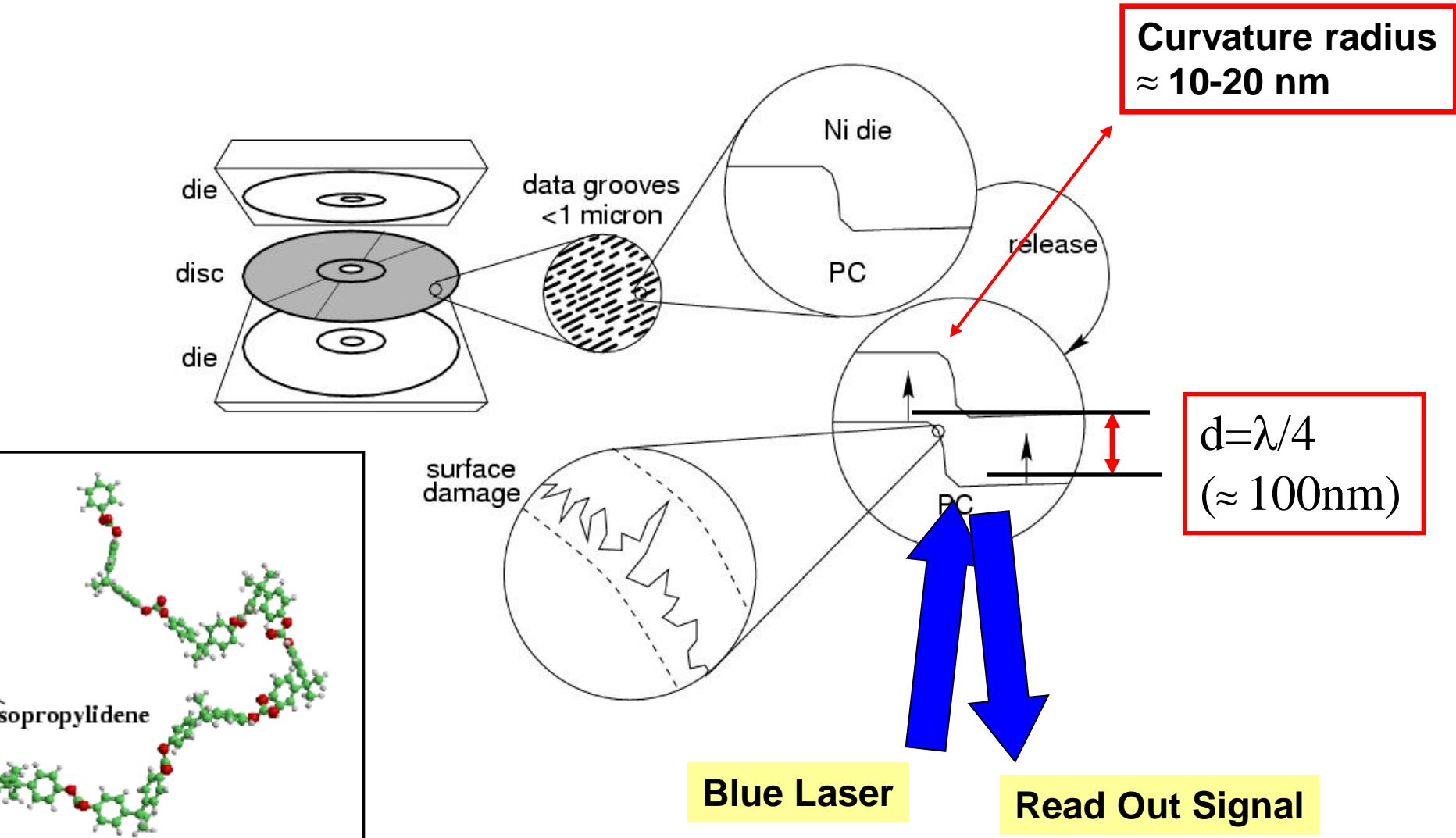
Data pits and guiding grove of a
PC “blue laser optical disc”



Bayer Materials

500 nm

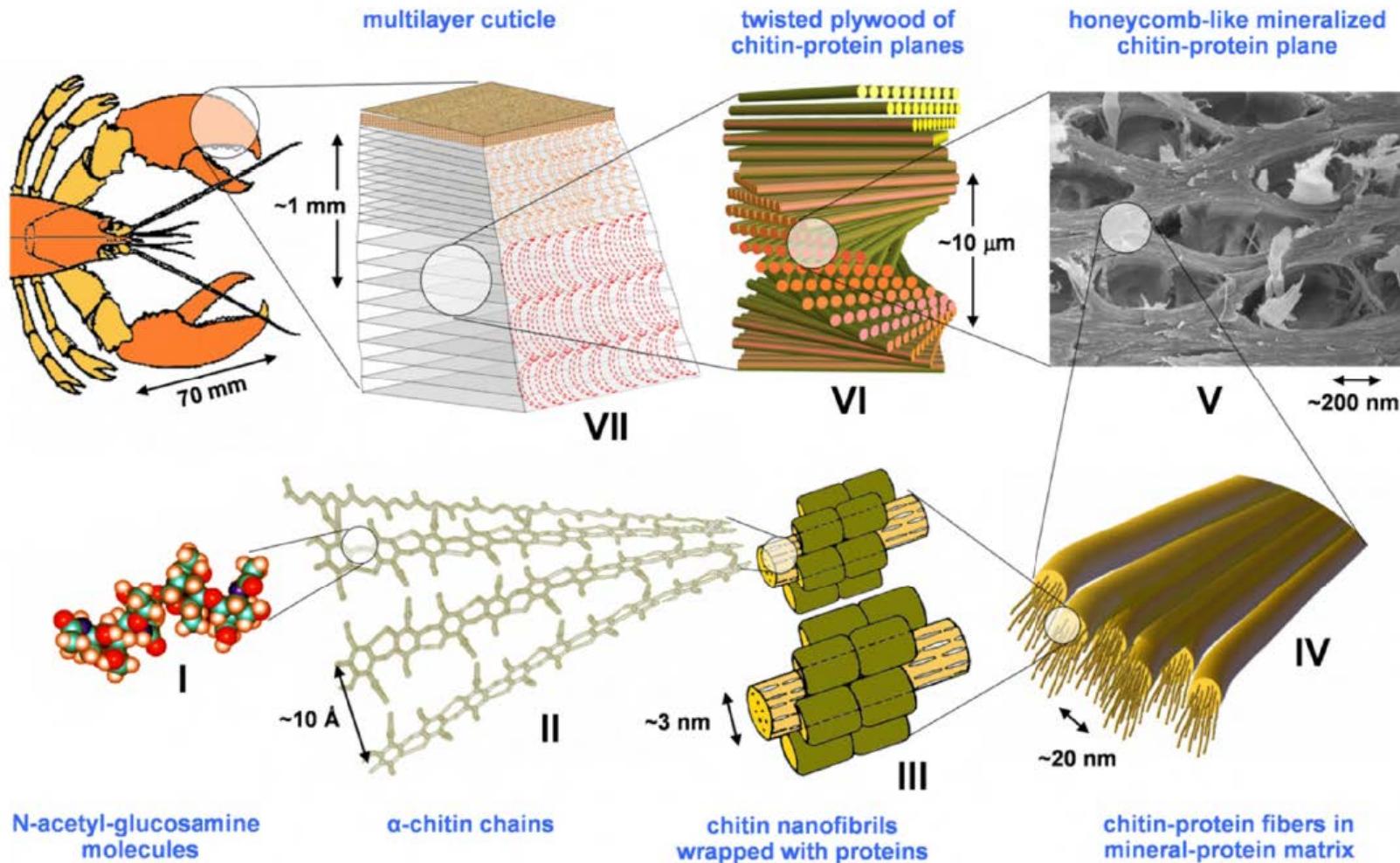
Why study Polycarbonate and the PC/Ni interface?



Hard-soft, organic-inorganic composite

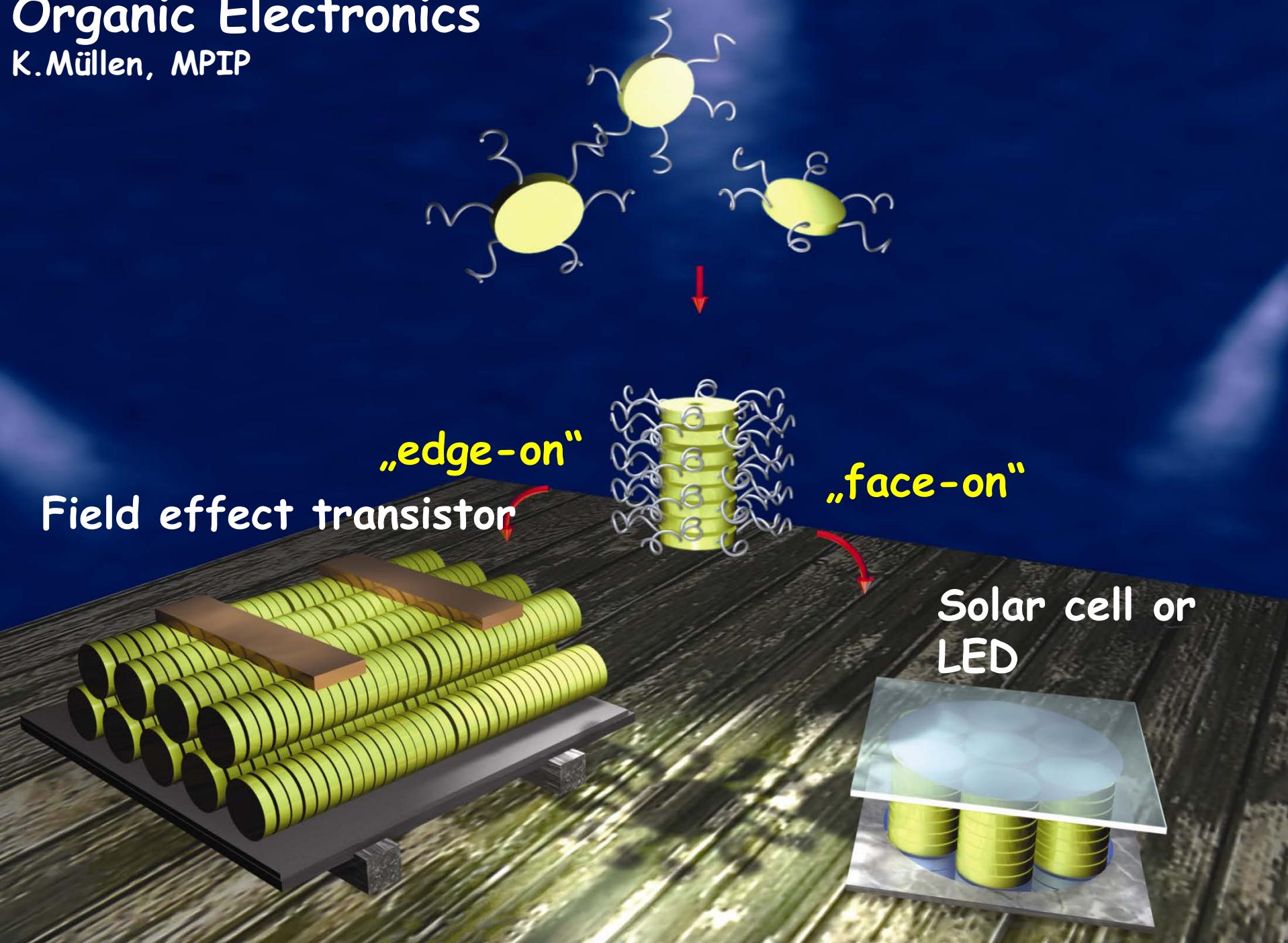


Introduction - Hierarchical organization



Organic Electronics

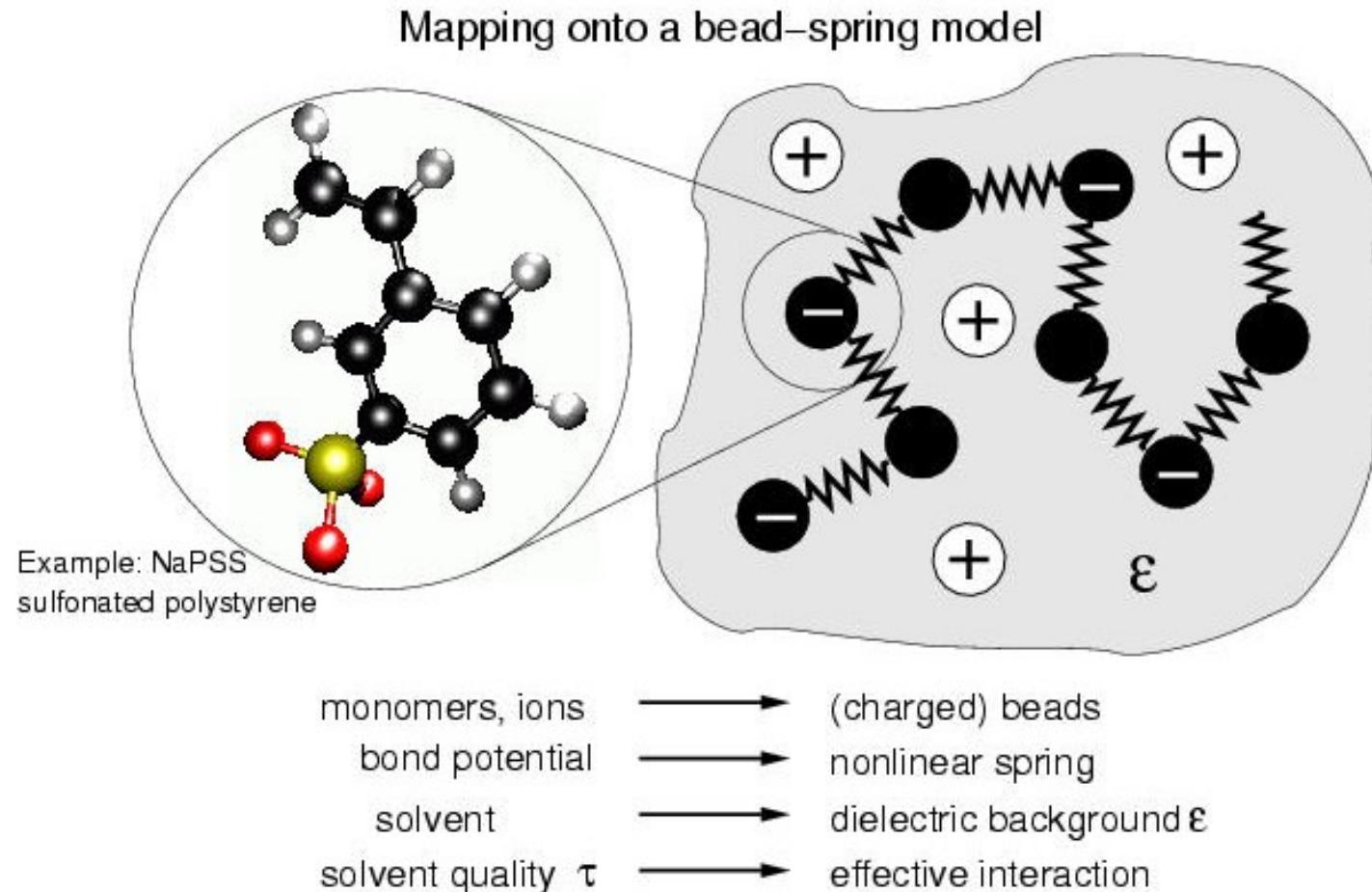
K. Müllen, MPIP



Organic Electronics: Morphology and Processing

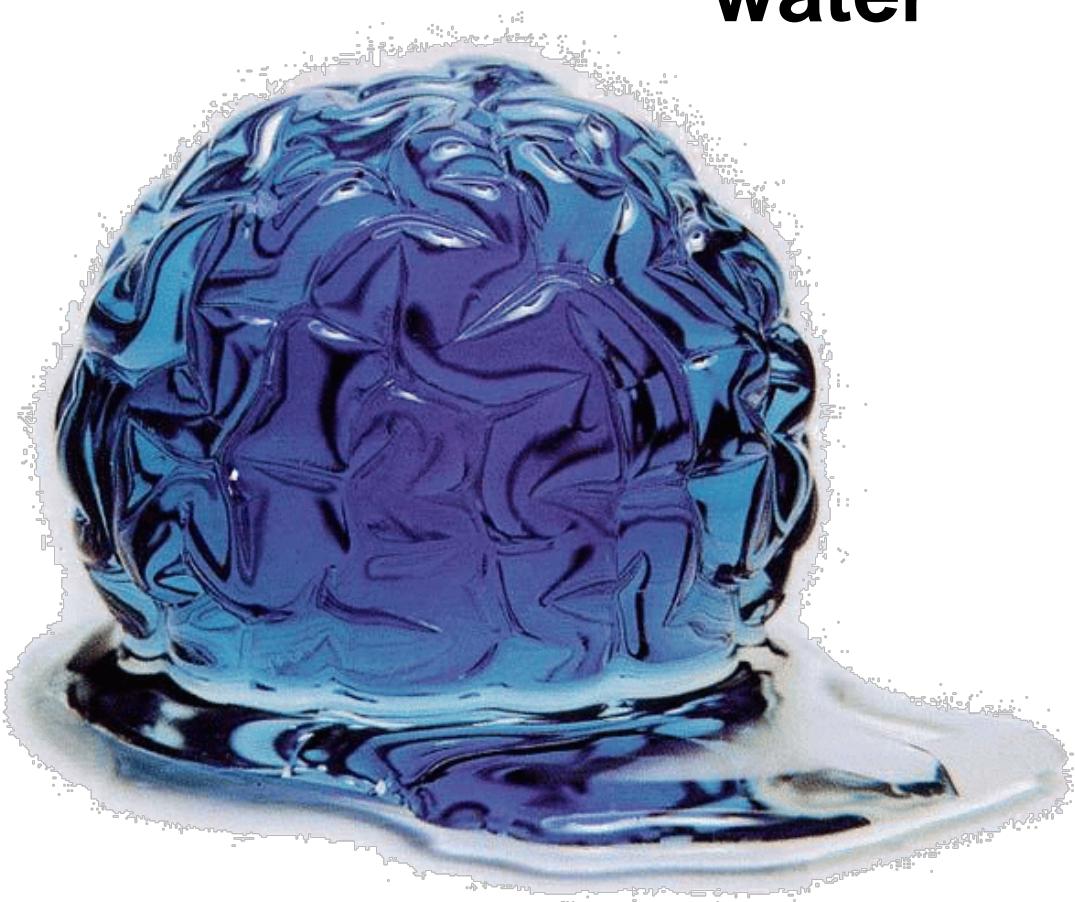


Schematische Darstellung: Polyelektrolyte



Hydrogels

$\approx O(1)$ (!!) Giant Molecule, swollen from water



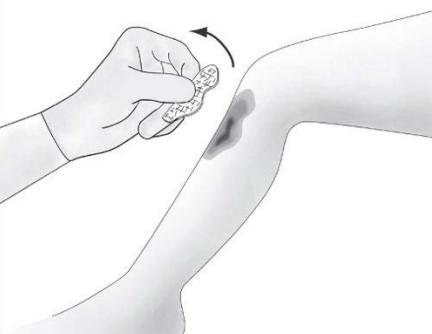
Polymers, usually not water soluble

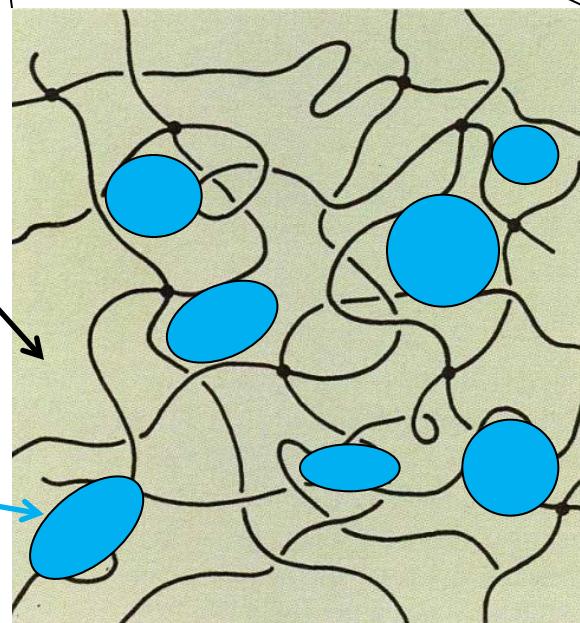
Exceptions e.g.
electrically charged
polymers:

Polyelectrolytes
(holds for most
biopolymers, e.g. DNA)

Hydrogels, Superabsorber

DRACO Wundheilmittel
BASF

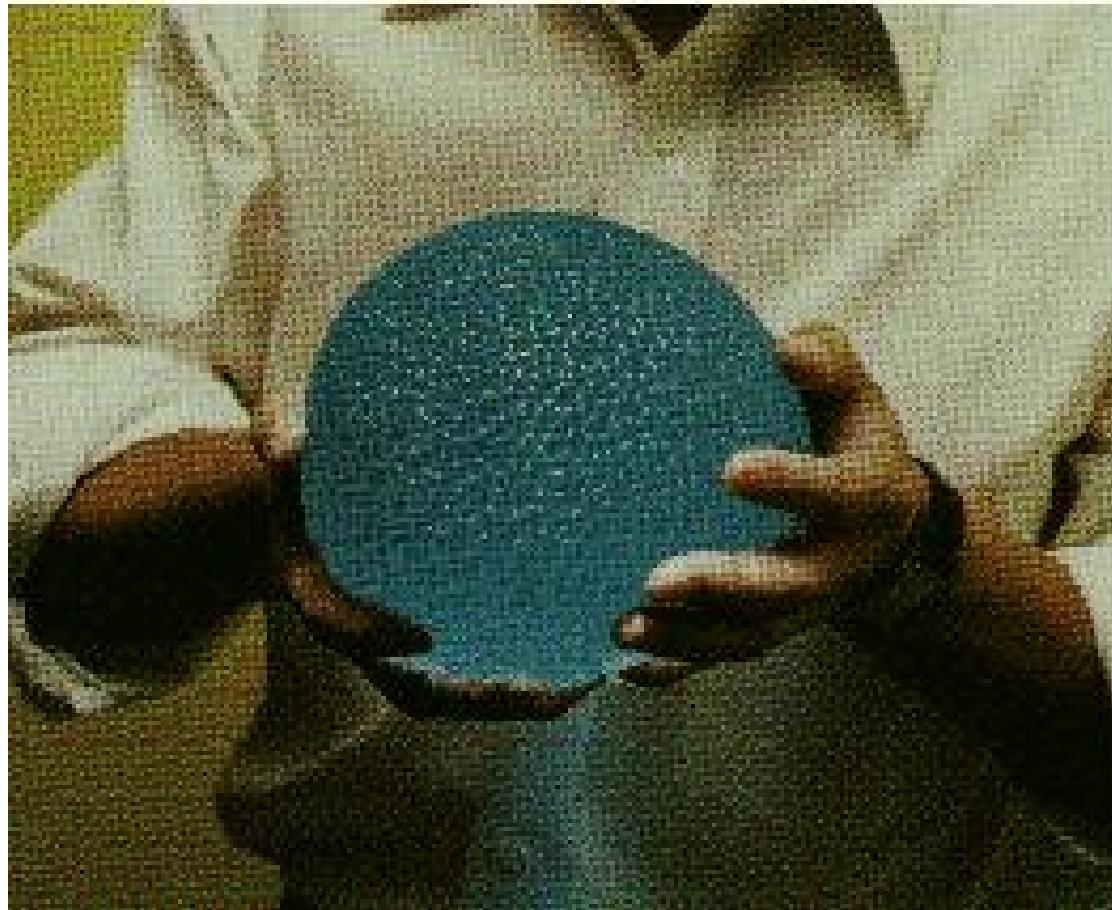




Fat

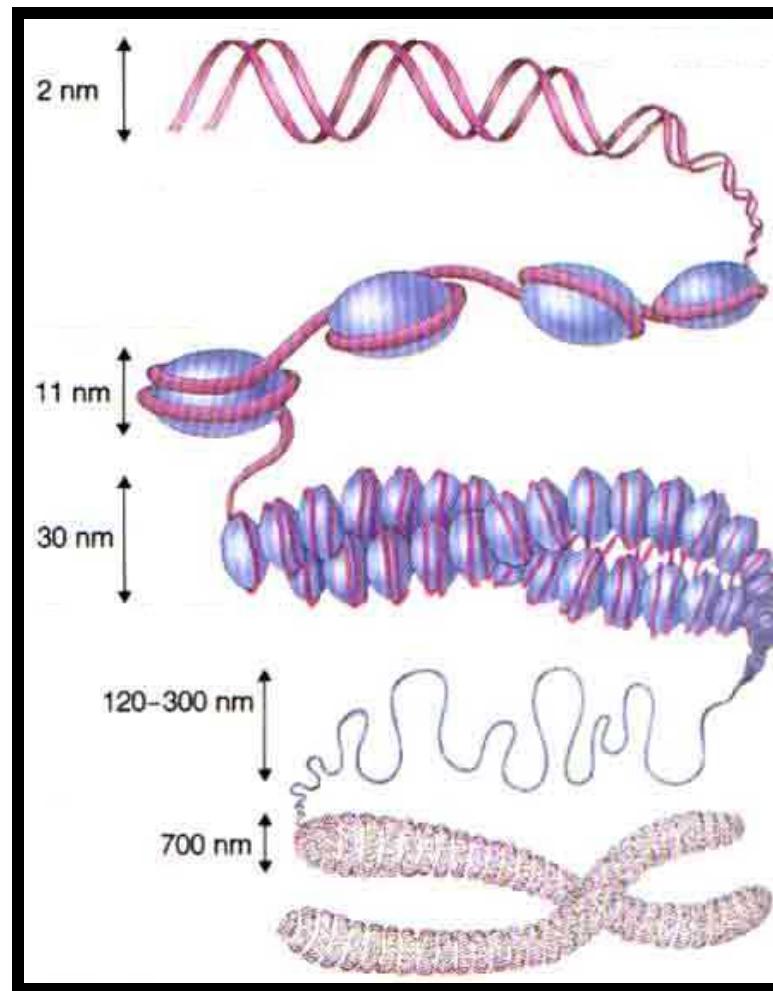
Water

Hydrogel $O(1)$ (!!) Giant Molecules swollen from water

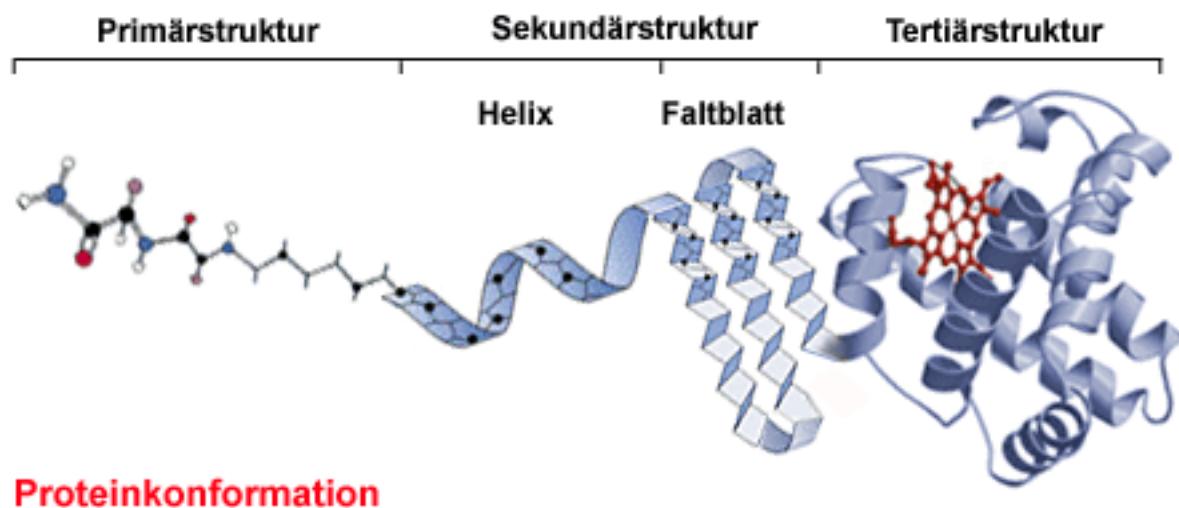


Tanaka, MIT

DNA



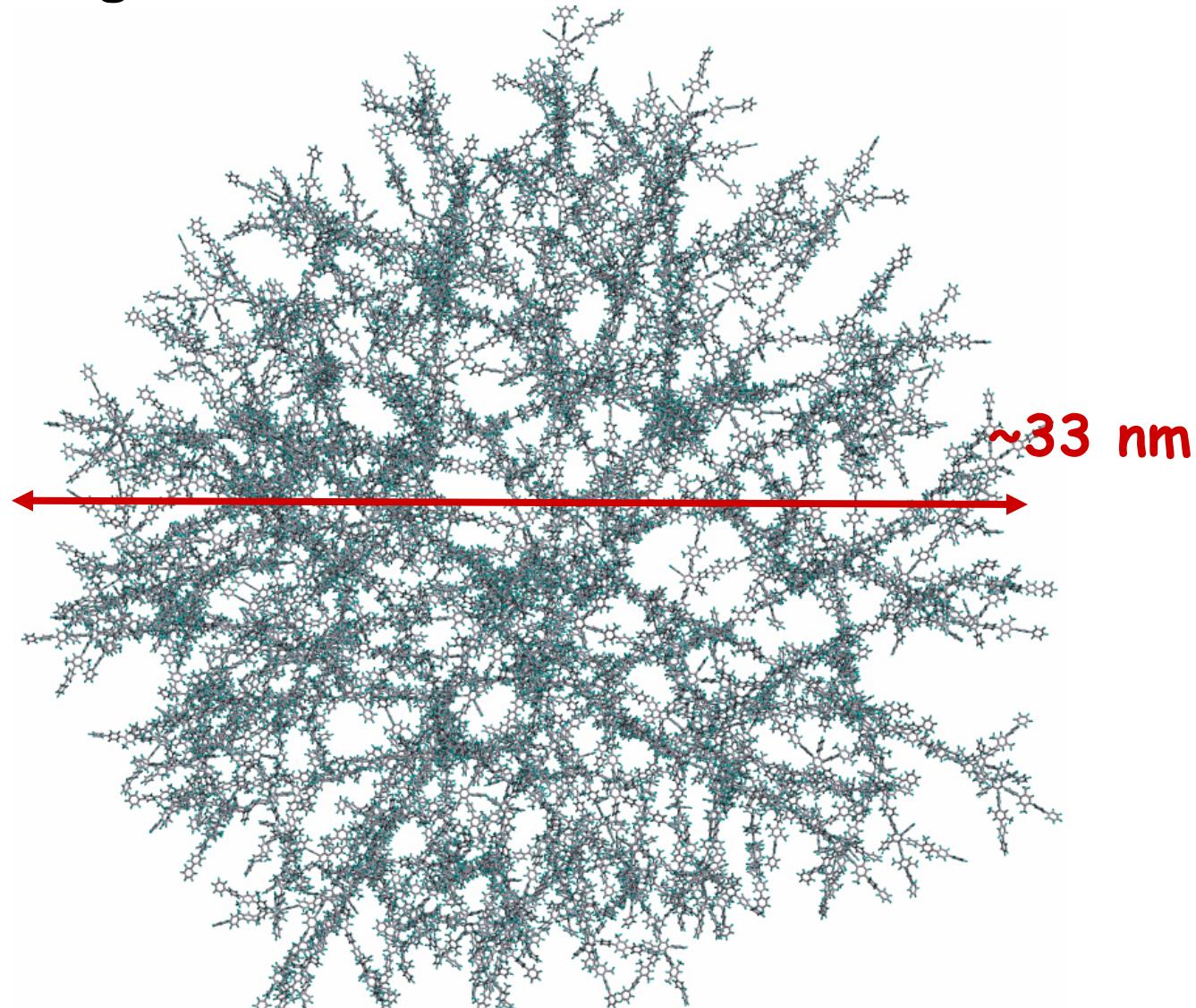
Proteins



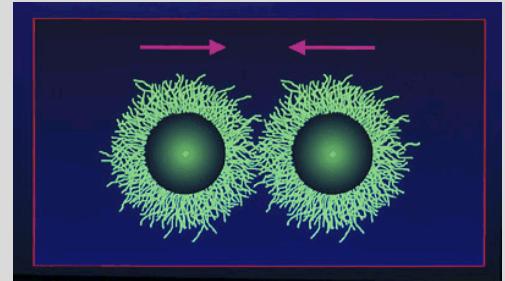
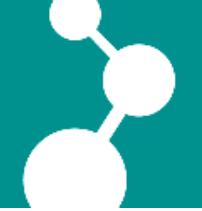
**Monodisperse, soluble,
semi-rigid (3D) Polymer**

MW = 546404 g/mol

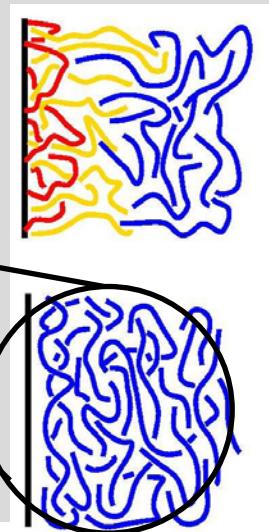
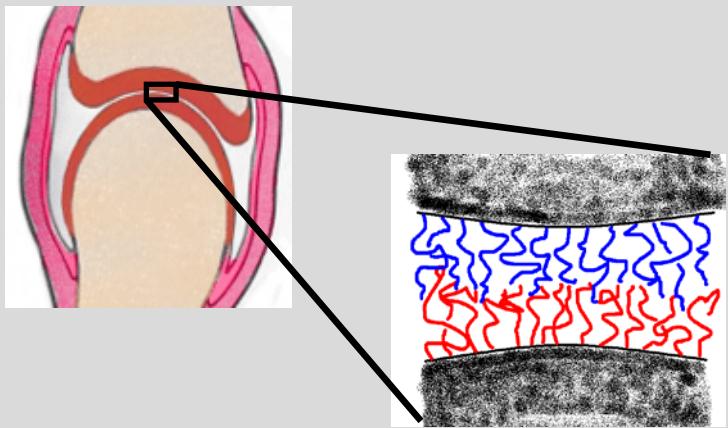
5592 Benzene rings (K. Müllen)



Soft Matter - Polymers



- Adsorption of polymers, wetting of polymers
- Adsorption of a single chain – relation to surface magnetism





Soft Matter

“Soft” means:

- low energy density
- nanoscopic length scales (10\AA ... 1000\AA)
- large fluctuations
- thermal energy $k_B T$
relevant energy scale

100 to 100000 times softer than normal crystals

Energy Scale $k_B T$ for $T=300K$



$$E = 1.38 \cdot 10^{-23} J / K \cdot 300K$$

$$kT \approx 4.1 \cdot 10^{-21} J$$

$$kT \approx 2.5 \cdot 10^{-2} eV$$

$$kT \approx 9.5 \cdot 10^{-4} E_H$$

$$kT \approx 4.1 pNm$$

$$kT \Rightarrow 200 cm^{-1}$$

$$kT \Rightarrow 0.6 kcal / mol$$

$$kT \Rightarrow 2.5 kJ / mol$$

$$E \approx 3 \cdot 10^{-19} J \approx 80kT$$

$$E \approx 4kT - 10kT$$

Electronic structure, CPMD

Quantum Chemistry

Biophysics Membranes, AFM

Spectroscopy

Chemical Bond
Hydrogen Bond

Energy Scale $k_B T$ for $T=300K$



$$E = 1.38 \cdot 10^{-23} J / K \cdot 300K$$

$$kT \approx 4.1 \cdot 10^{-21} J$$

$$kT \approx 2.5 \cdot 10^{-2} eV$$

$$kT \approx 9.5 \cdot 10^{-4} E_H$$

$$kT \approx 4.1 pNm$$

$$kT \Rightarrow 20$$

$$kT \Rightarrow 0.05 kcal / mol$$

$$kT \Rightarrow 2.5 kJ / mol$$

$$E \approx 3 \cdot 10^{-19} J \approx 80kT$$

$$E \approx 4kT - 10kT$$

Electronic structure

Onset of Chemistry

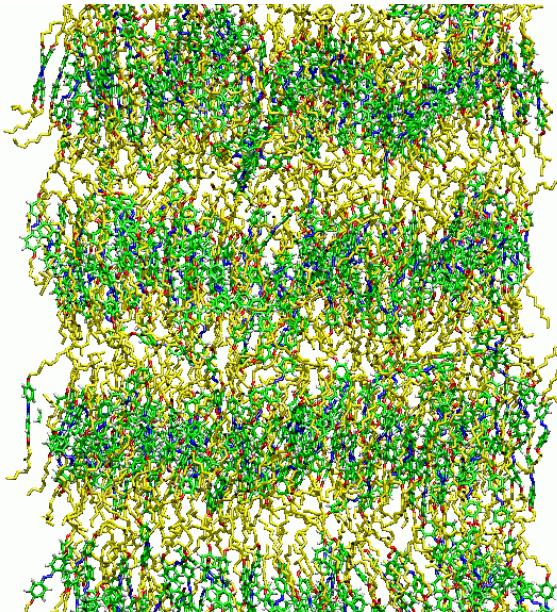
Topophysics Membranes, AFM

Spectroscopy

Chemical Bond
Hydrogen Bond

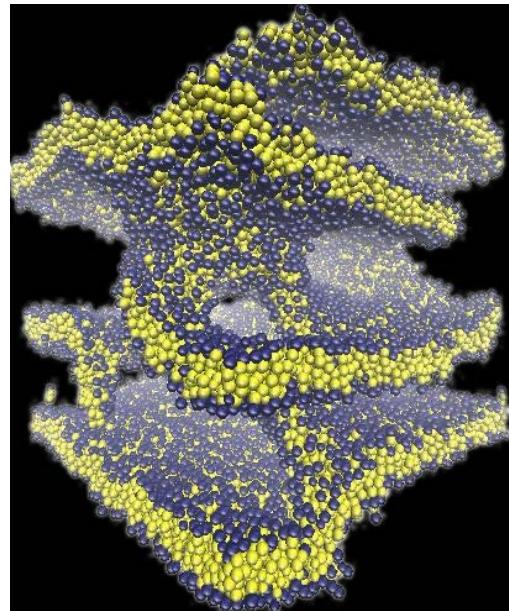
Green light photon $\approx 2-2.5 \text{ eV} \approx 100 \text{ kT}$

Soft Matter – Nanostructured Matter



$$\text{Volume } V = L^3$$

A: Surface
and/or
Interface Area



Nanoscopically Structured Material:

$$V/A \ll 1 \mu\text{m}$$

=> Distinction Bulk vs Surface/Interface not useful

Definition of A depends on the question studied!!

Soft Matter



“Soft” means:

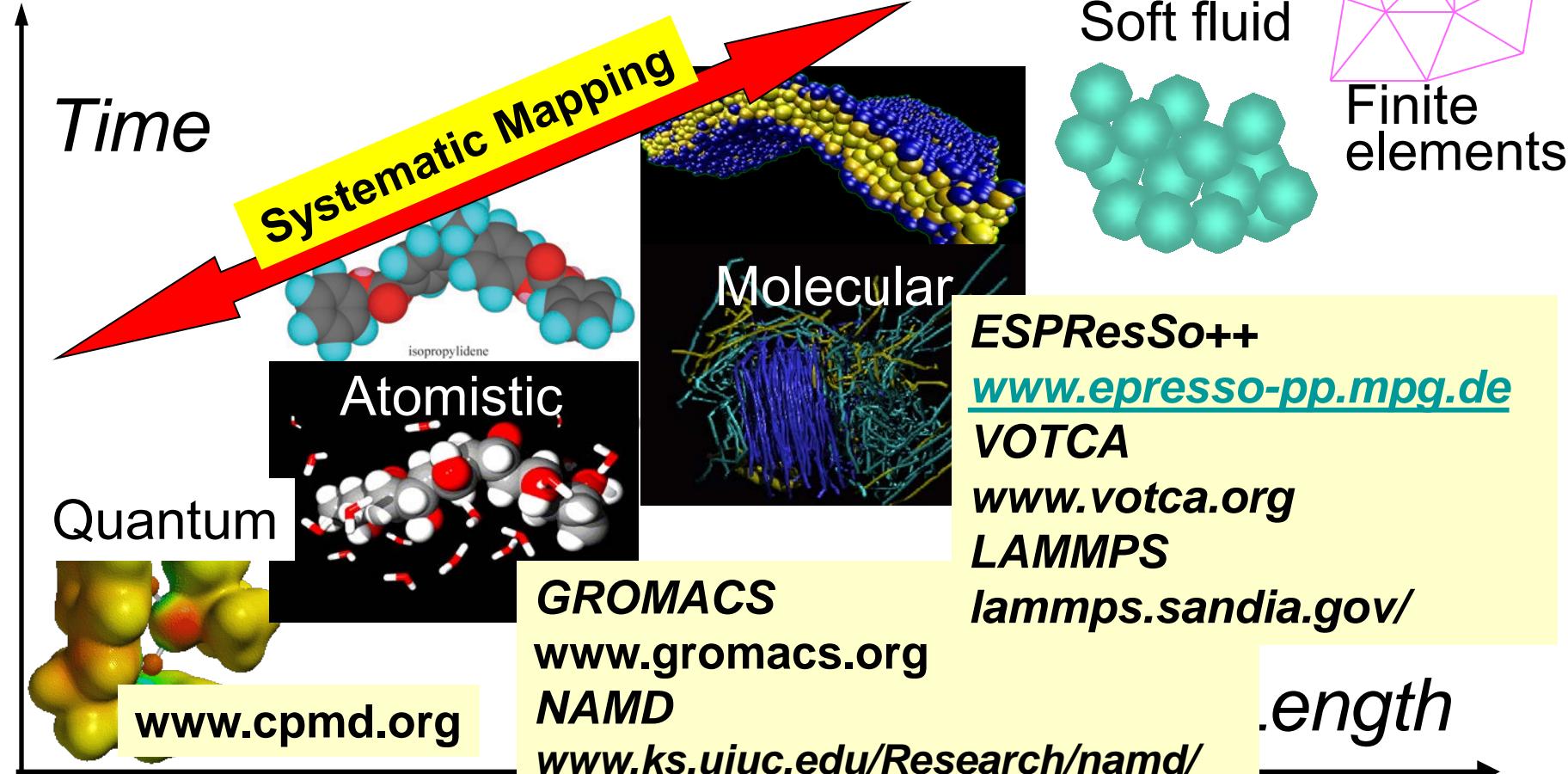
Large fluctuations,
conformational entropy

Or look for energies $\gg k_B T$
to dominate fluctuations

Often dominated by finite size effects!

*Combination of generic and chemistry
specific input needed!*

Characteristic Time and Length Scales



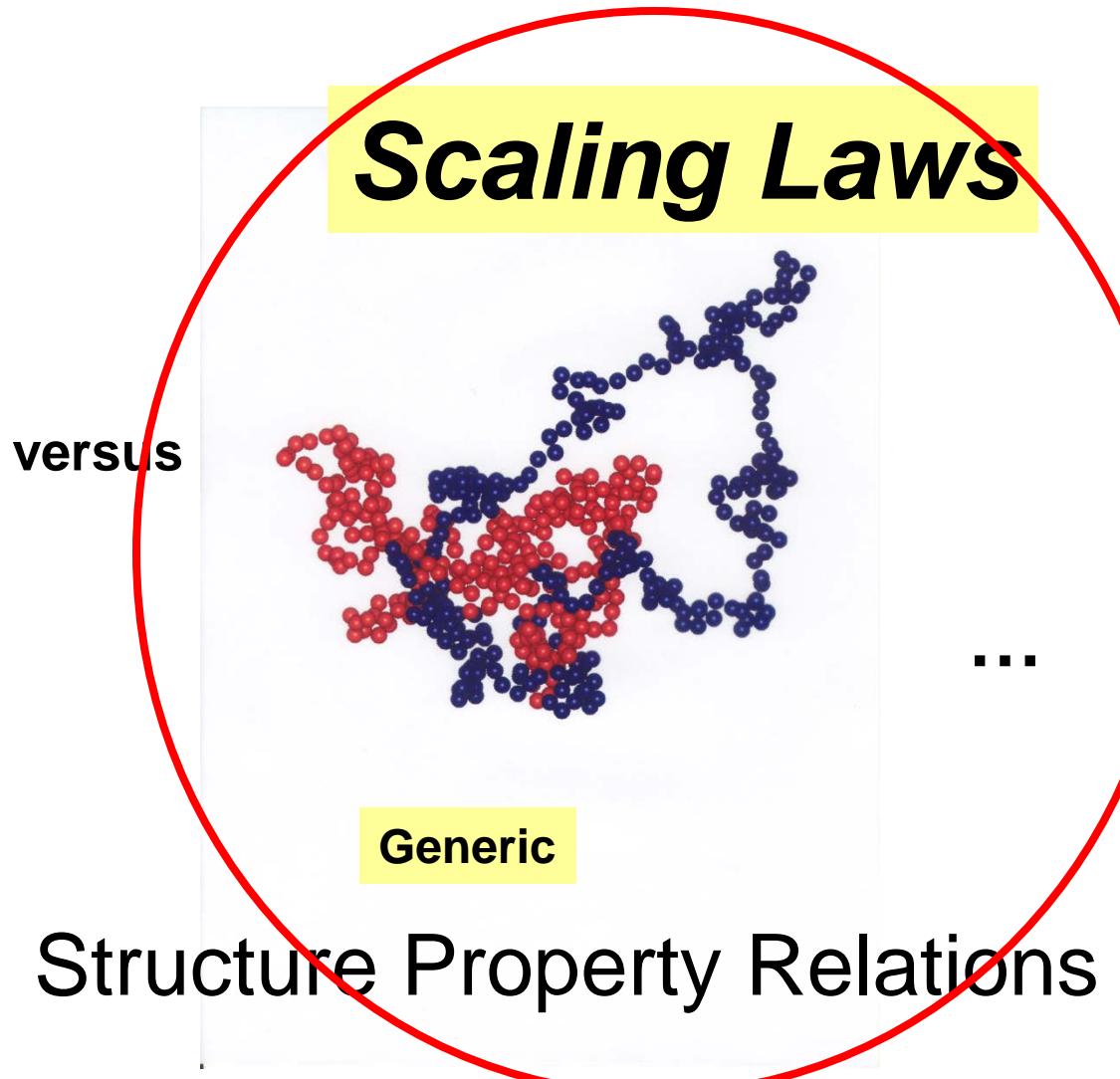
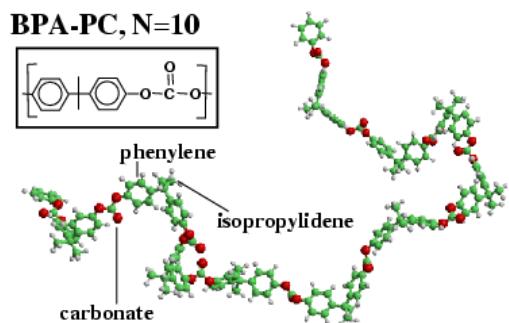
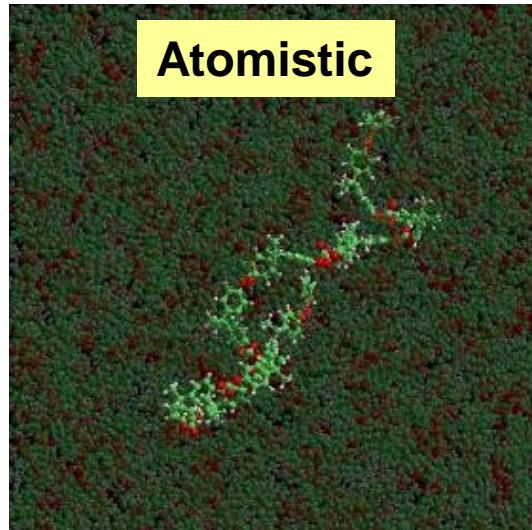
Local Chemical Properties --- Scaling Behavior of Nanostructures
Energy Dominance --- Entropy Dominance of Properties



How can theory, simulations contribute?

Polymers

relevant interactions well defined





Simulation Models for Bilayers

Relevant Variables – More Complex

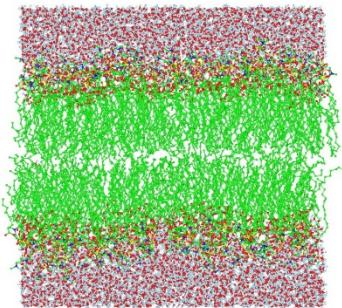


much detail

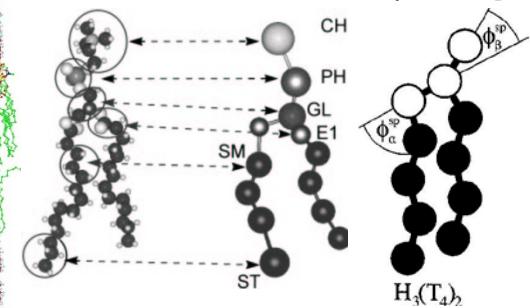
little detail

increasingly coarse grained

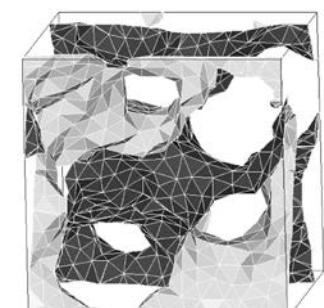
atomistic models



bead-spring models



triangulated surfaces



"standard"
Lennard-Jones

DPD

solvent free

increasing numerical efficiency

But, what do we loose...

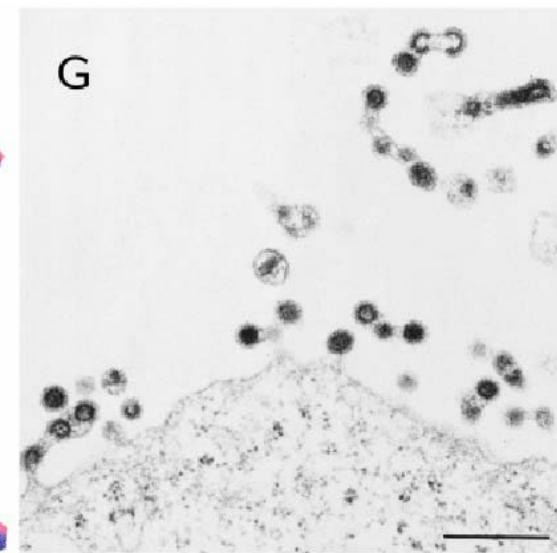
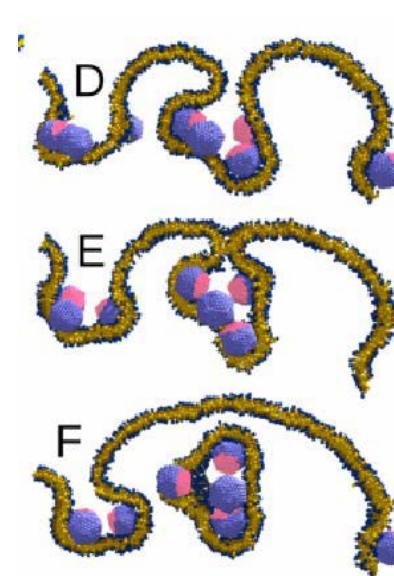
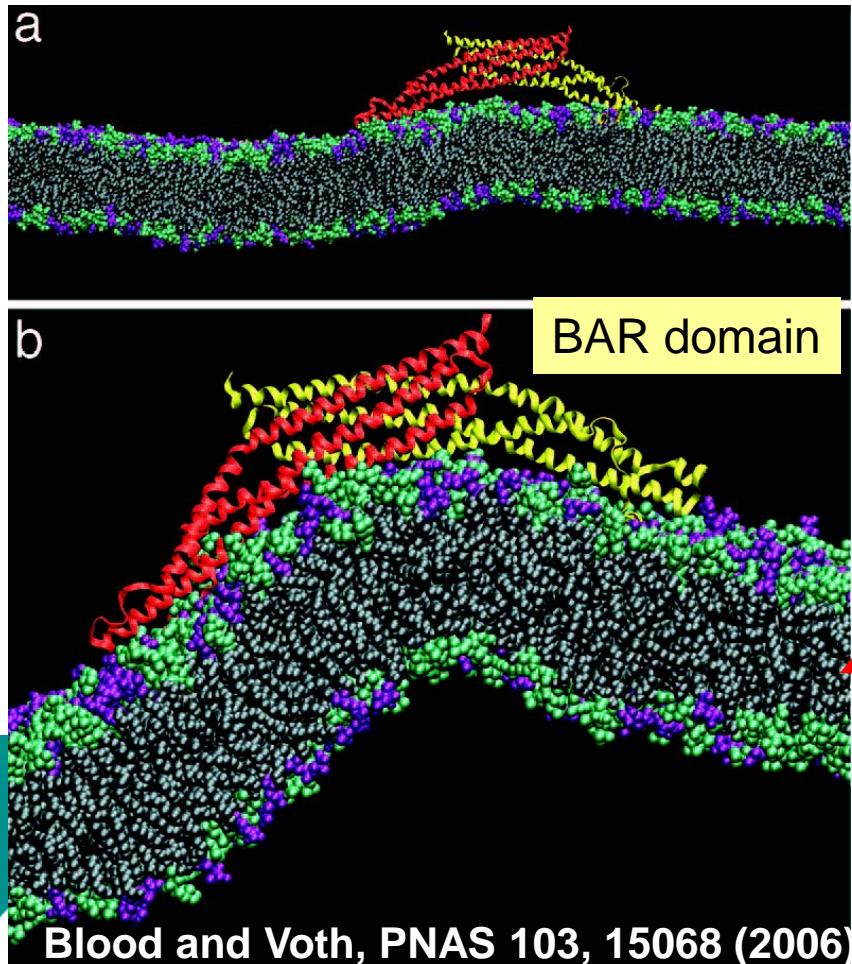


Simulation Models for Bilayers

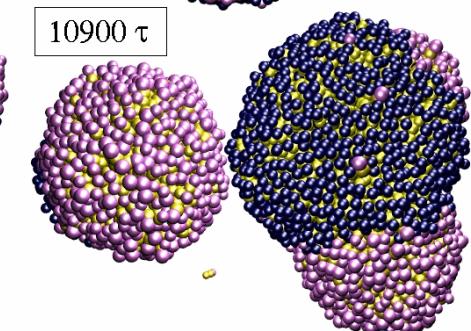
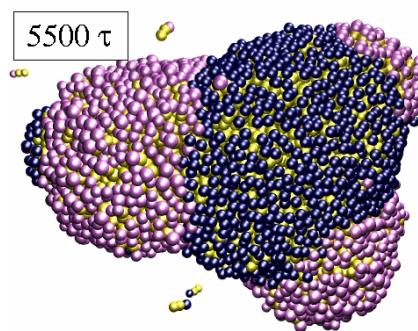
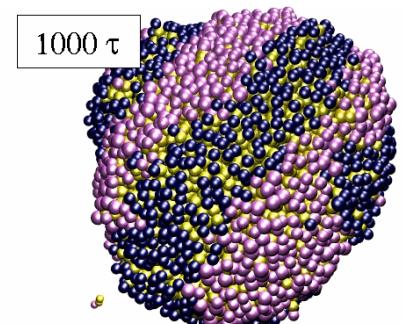
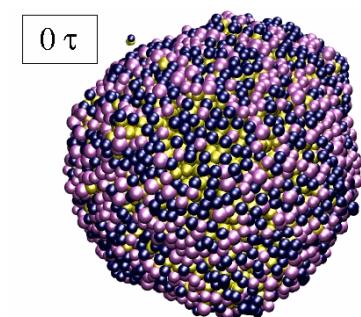
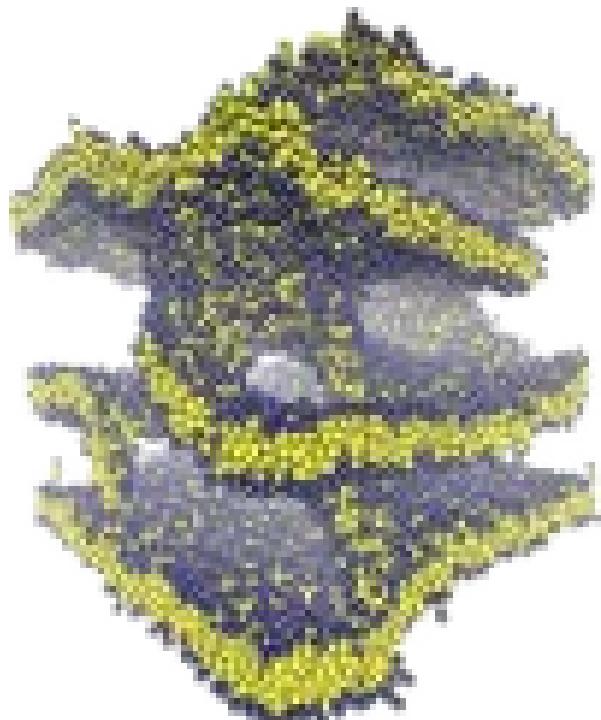
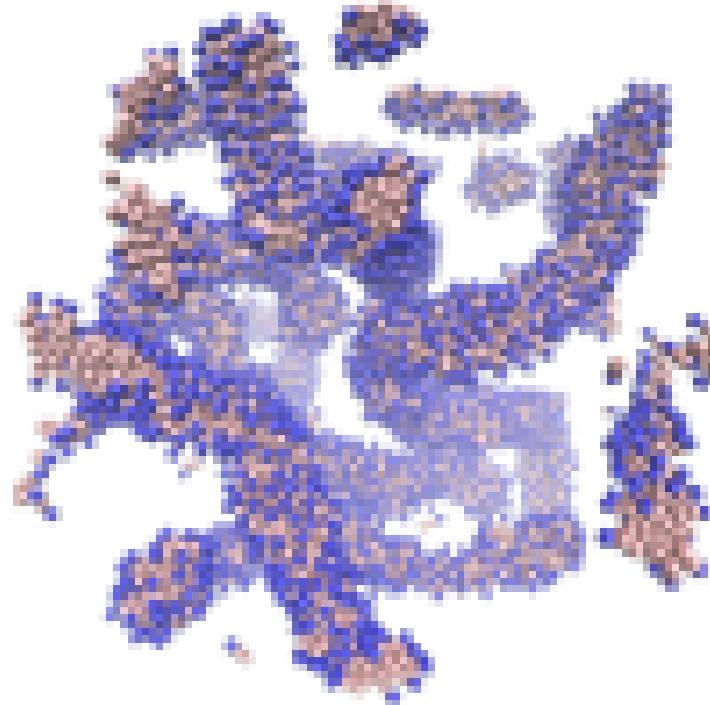
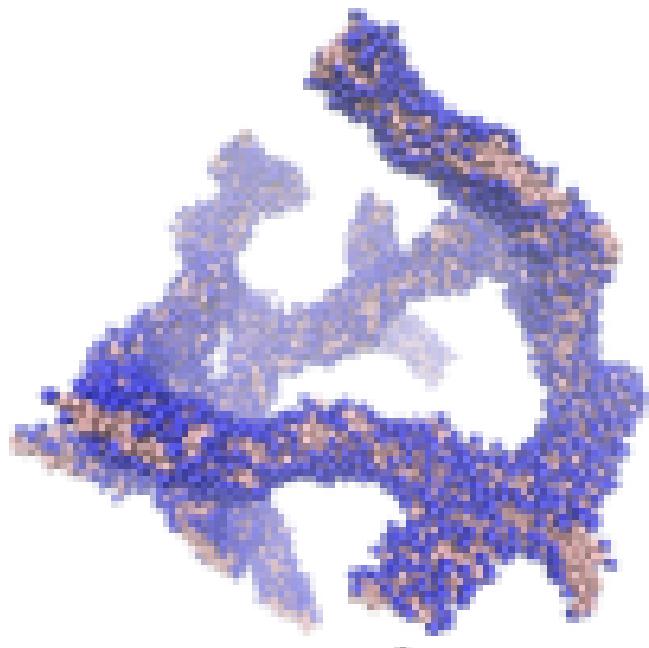
Relevant Variables – More Complex



increasingly coarse grained

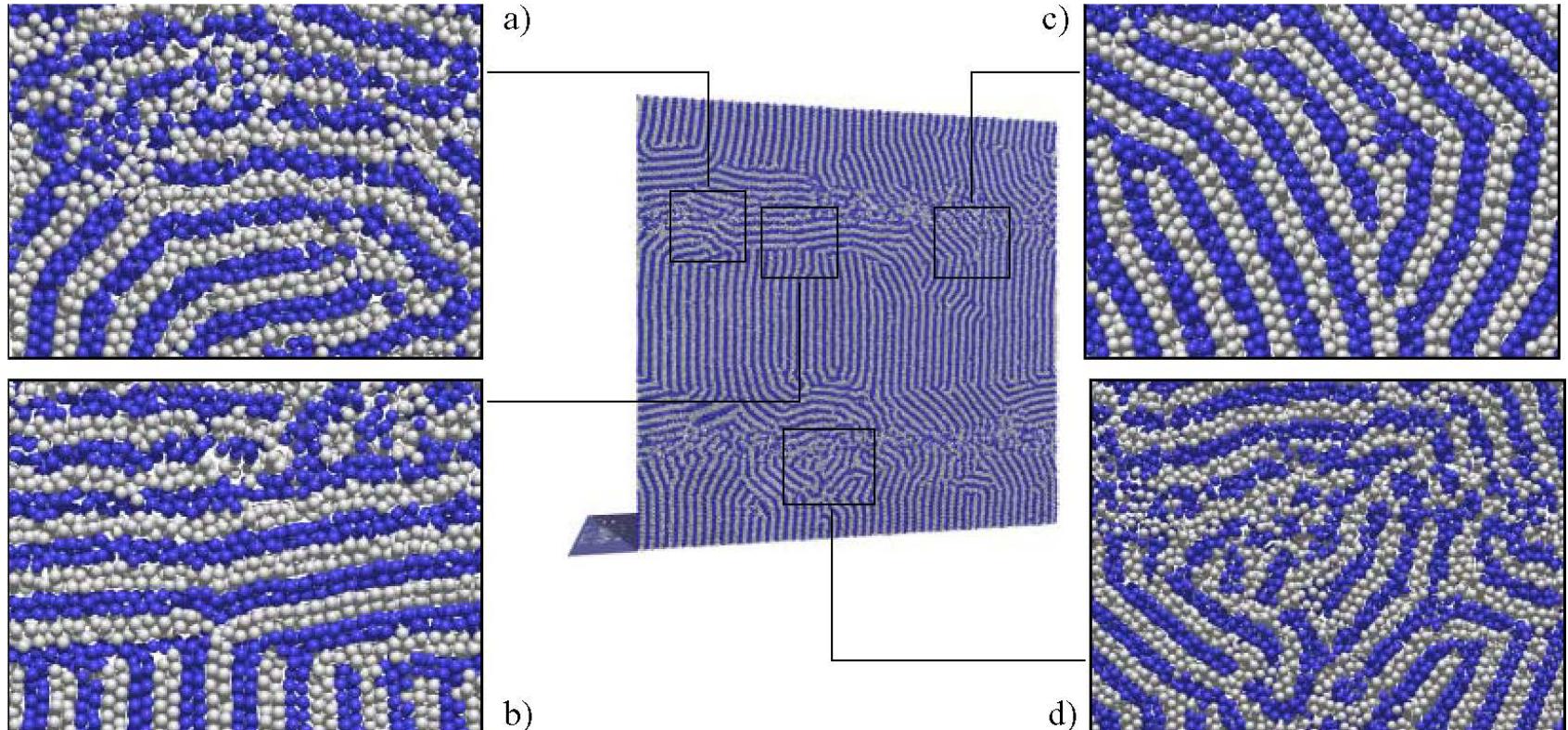


Deserno et al, Nature, 2007



Shearing Lamellar Systems: Dimers

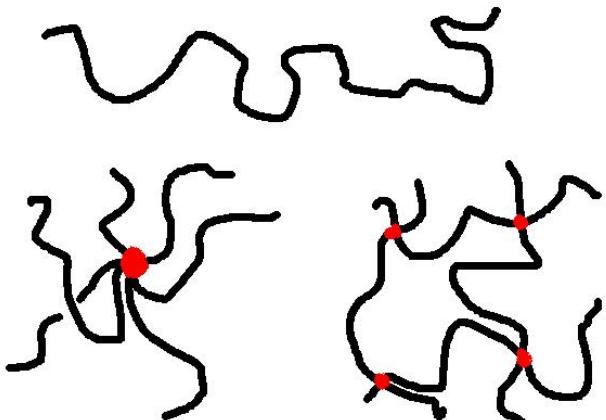
T. Soddeemann, H. X. Guo, B. Dünweg, K. K., PRE 2003



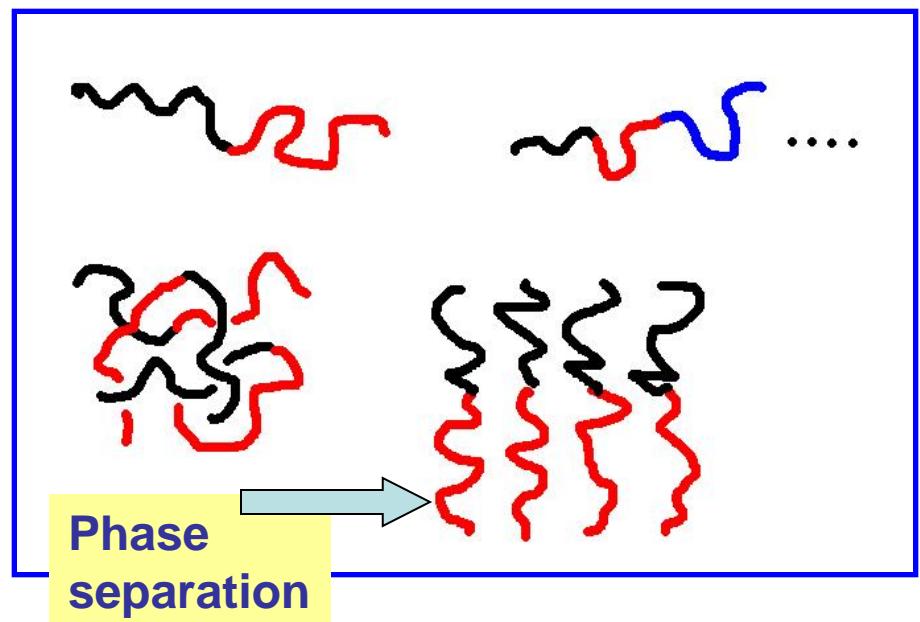


Architectures

Linear Polymers



Block Copolymers



Branched Polymers

Intramolecular entropy :
 $S = kT \Omega(N)$

Soft Matter - Polymers



Polymer mixtures

(A. Sariban, H.P. Deutsch, M. Müller, D. P. Landau, K. Binder,...

later M. Müller, K. Daoulas

(SCFT-particle hybrid methods)

plus a huge body of literature from other groups

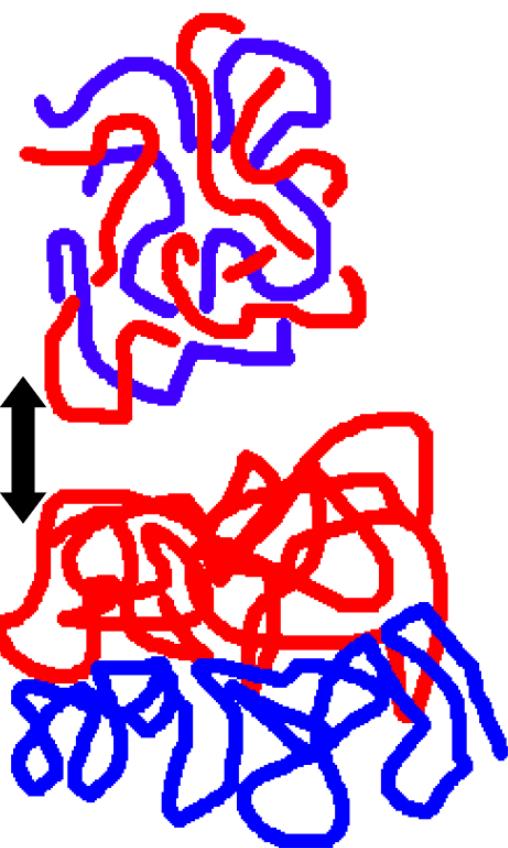
1/N power law

Soft Matter - Polymers



Polymer mixtures

1/N power law for critical interaction x_c



$$\begin{aligned}\langle R^2(N) \rangle &\propto N \\ P(\bar{r}) &\propto \exp(-3\bar{r}^2 / 2\langle R^2(N) \rangle) \\ \rho_{self} &\propto \frac{N}{\langle R^2(N) \rangle^{3/2}} \propto N^{-1/2}\end{aligned}$$

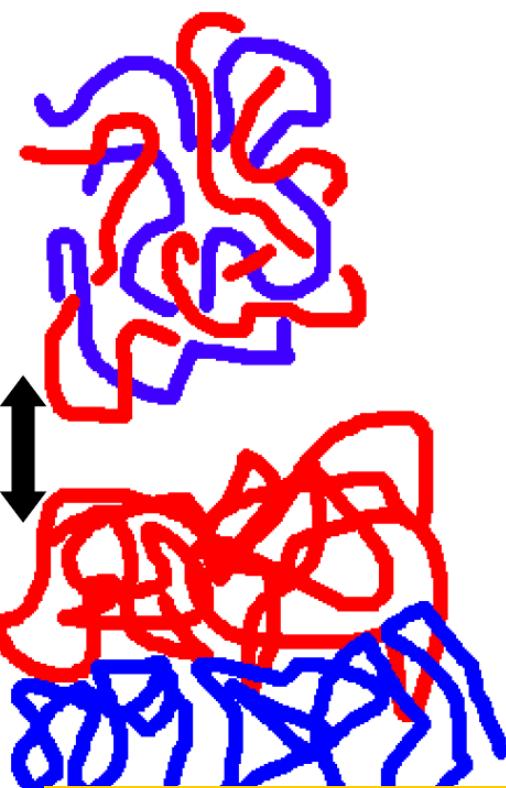
Soft Matter - Polymers



Polymer mixtures

1/N power law for critical interaction χ_c

#AA, #BB, #AB contacts = $O(N)$



$$U_{AB} \propto N \chi_{AB}$$

$$U_{AA} \approx U_{BB} \propto N \chi$$

$$\chi_{eff} = \chi_{AB} - \chi$$

Phase separation, critical interaction



$$\chi_{eff}^c = const \cdot N^{-1}$$

“chemistry” “generic”

Intra-chain entropy invariant => small energy differences => phase separation

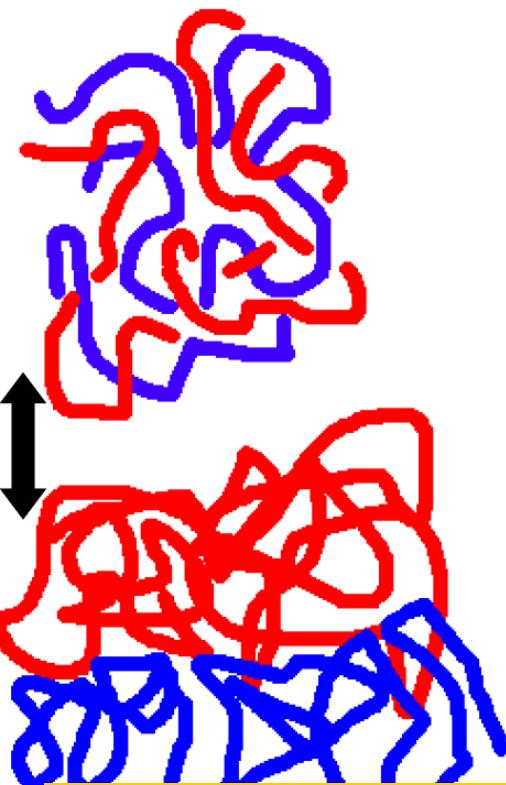
Soft Matter - Polymers



Polymer mixtures

1/N power law for critical interaction x_c

#AA, #BB, #AB contacts = $O(N)$



Free Energy of Mixing:

$$\frac{\Delta F}{kT} = \frac{\phi_A \ln \phi_A}{N_A} + \frac{(1-\phi_A) \ln(1-\phi_A)}{N_B} + \phi_A(1-\phi_A)\chi_{eff}$$

A-B interaction energy

Translational entropy per chain

$$\phi_A + \phi_B = 1$$

Phase segregation:

$$\frac{\partial^2 \Delta F}{\partial \phi_A^2} = \frac{\partial^3 \Delta F}{\partial \phi_A^3} = 0 \quad \Rightarrow \quad (\chi_{eff} N_A)_{crit} = 2$$

Intra-chain entropy invariant => small energy differences => phase separation

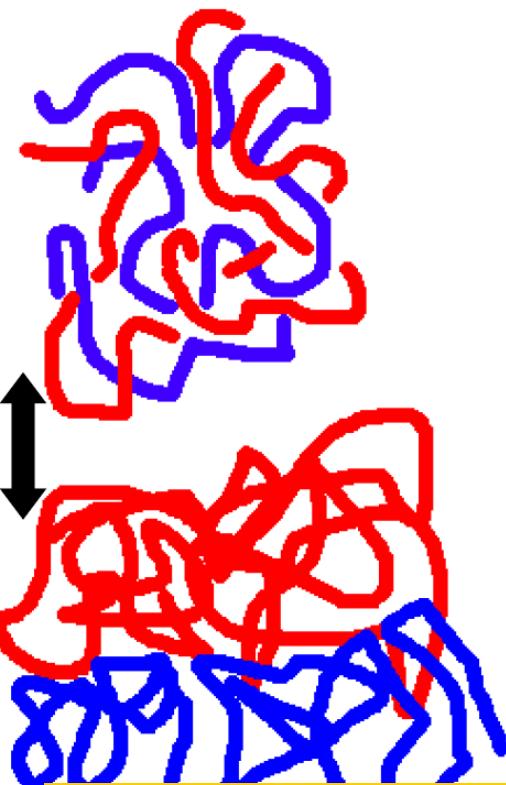
Soft Matter - Polymers



Polymer mixtures

1/N power law for critical interaction χ_c

#AA, #BB, #AB contacts = $O(N)$



Free Energy of Mixing:

A-B interaction energy



Highly disputed proposal by
K. S. Schweizer

y per chain

$$\chi_c \propto N^{1/2}$$

segregation:

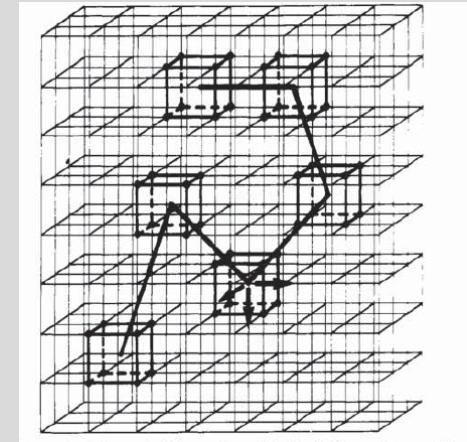
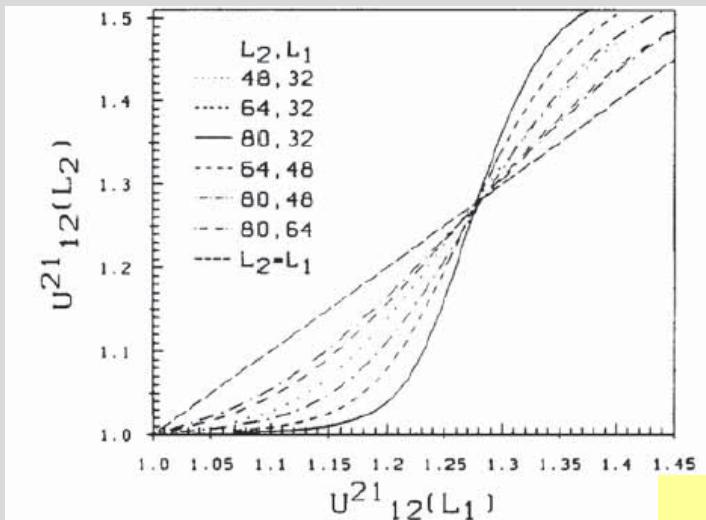
Intra-chain entropy invariant => small energy differences => phase separation

Soft Matter - Polymers

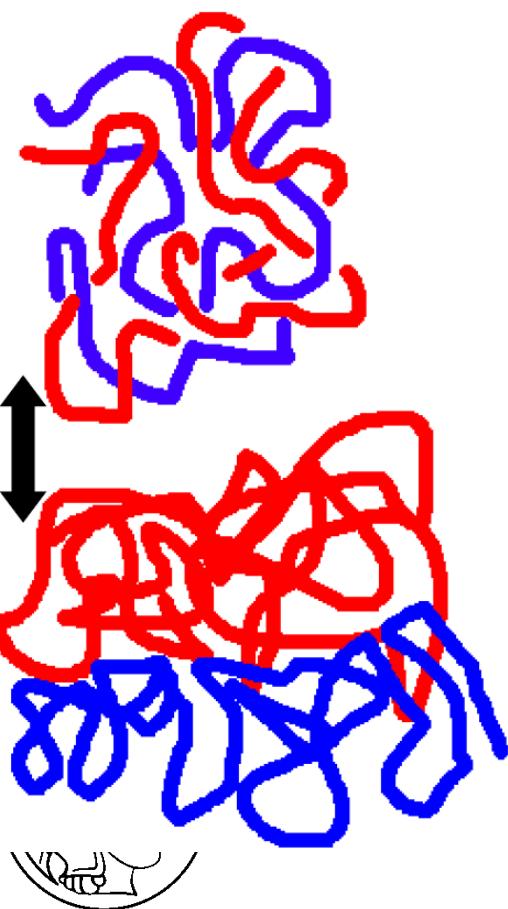


Check by Simulation: H.P. Deutsch, K. Binder 1992

Monte Carlo Simulation
Extensive finite size
scaling analysis employing
Cumulants of
A-B density differences



$$\frac{1}{\chi_{eff}^{crit}} = \frac{kT}{\varepsilon} = 2.15N + 1.35$$



Soft Matter - Polymers

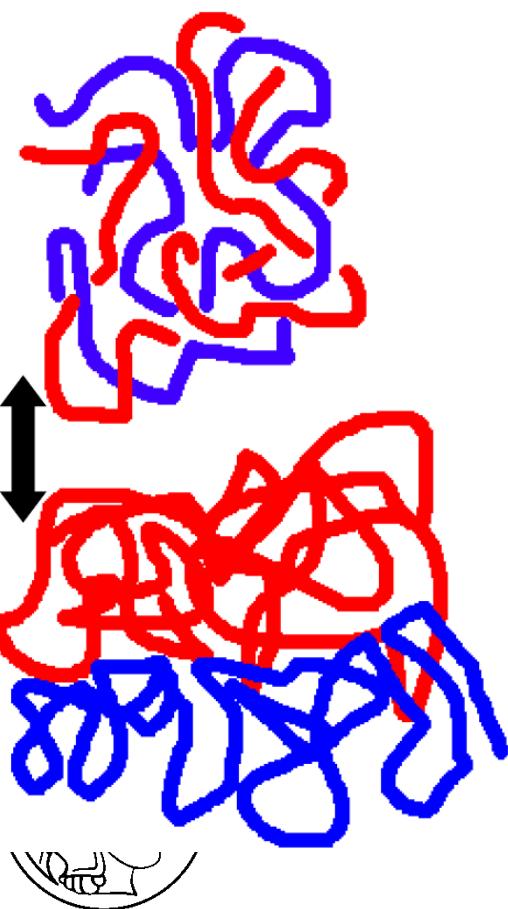


Check by Experiment (**shortly after simulations**):
F. S. Bates et al. 1992

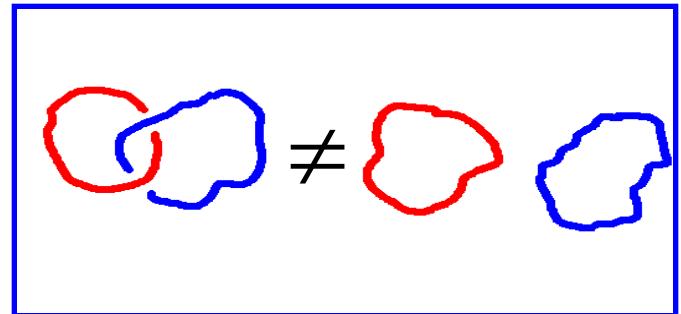
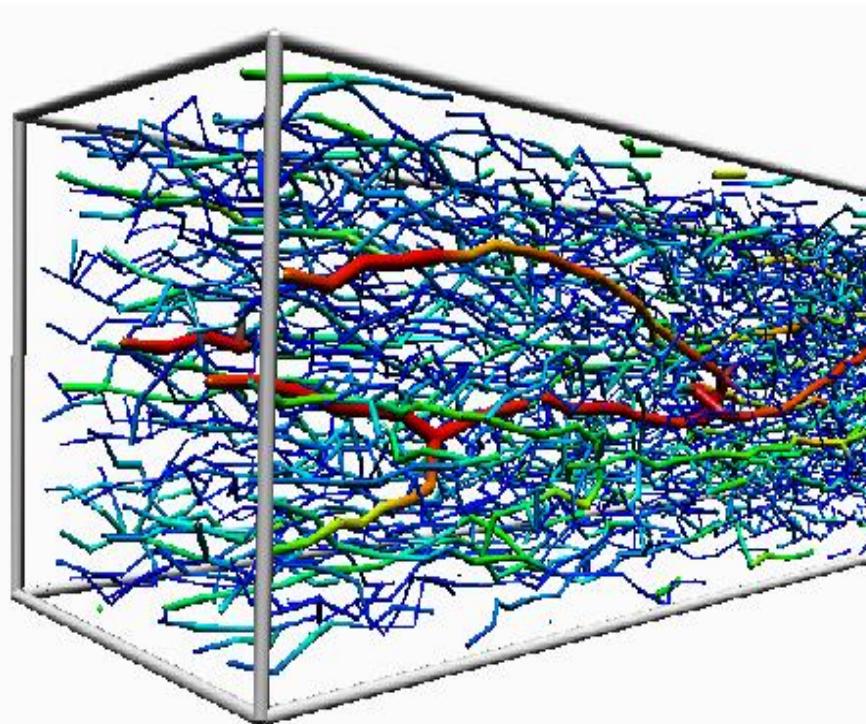
N-scattering experiments on partially deuterated
Poly(ethylene-propylene) copolymers

$(C_5D_xH_{(10-x)})_N$ mixed with $(C_5H_{10})_N$

$$\chi_{eff}^{crit} \propto N^{-\delta} \quad \delta = 1.01 \pm 0.05$$



Universal Confinement: Links & Stress



Local stress distribution

- **SMALL STRESS**
- **LARGE STRESS**

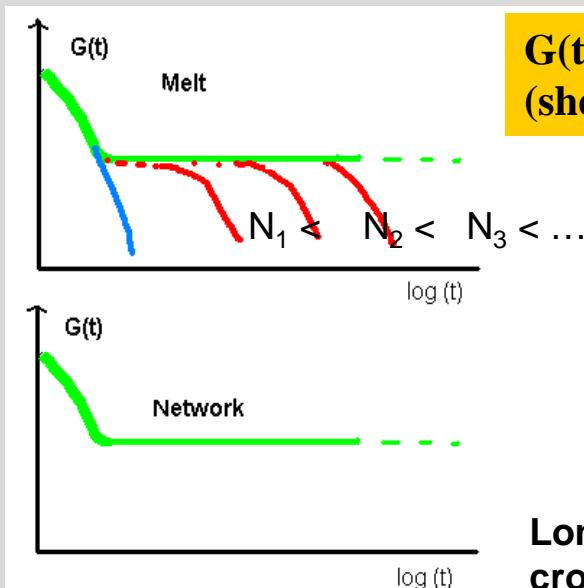
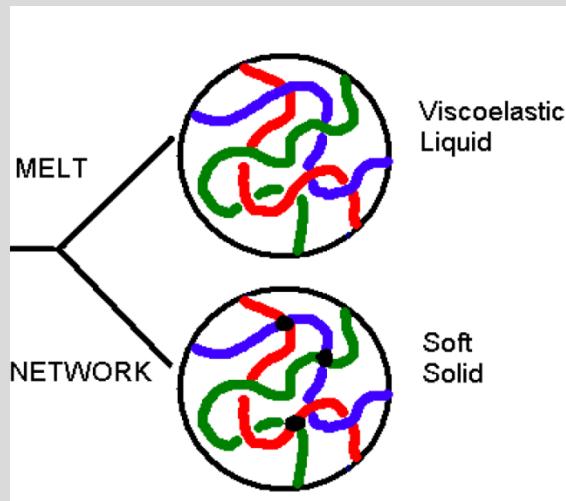
Soft Matter - Polymers



Edwards, de Gennes

Tube model, reptation concept

Classical Problem: MELT \leftrightarrow NETWORK



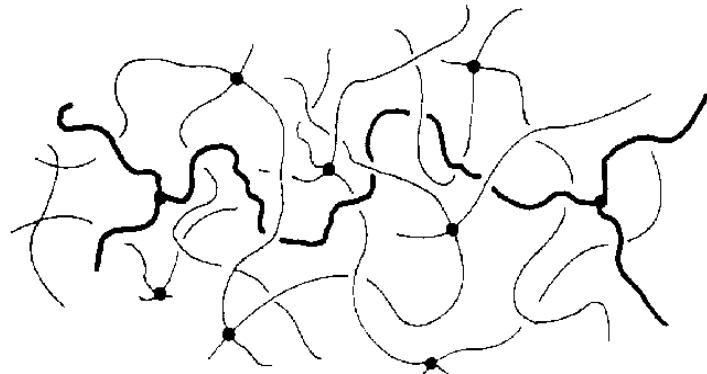
G(t): Time dependent (shear) modulus:

Long chains between crosslinks, $G(t)$ independent of N

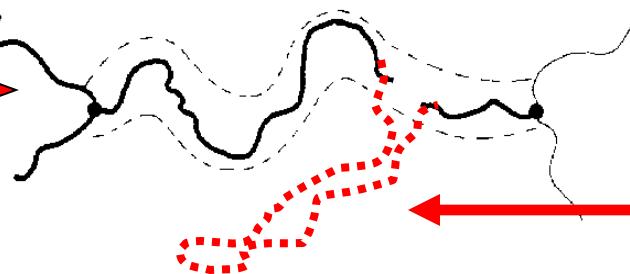
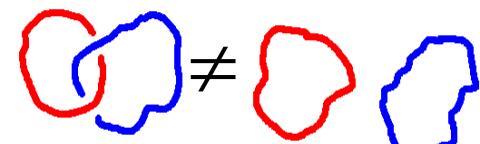




Tube Model

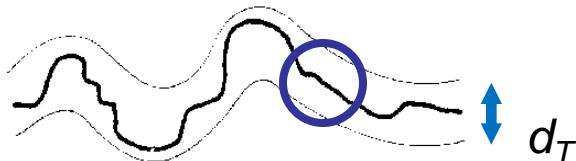


↓ ↓ Edwards 1967



“Tube leakage”:
Entropy penalty ($O(L/2)$)
for defects too large

↓ ↓ De Gennes 1971

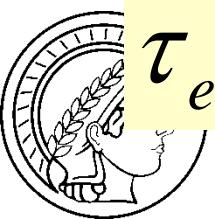


$$d_T \propto N_e^{1/2}$$

$$G^0 \propto N_e^{-1}$$

$$\tau_e \propto N_e^2$$

Unique length d_T





Tube Model

Concept, based entirely on conformational statistics of networks (soft solid) and melts (liquid).

$$\rightarrow d_T \propto N_e^{1/2}$$

N_e should be determined from conformational statistics of the polymers.

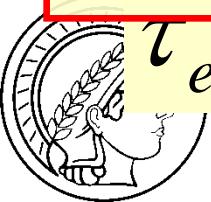
Unique length d_T

Edwards 1967

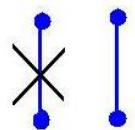
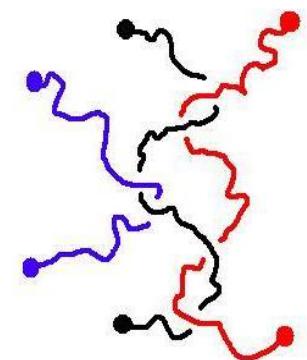
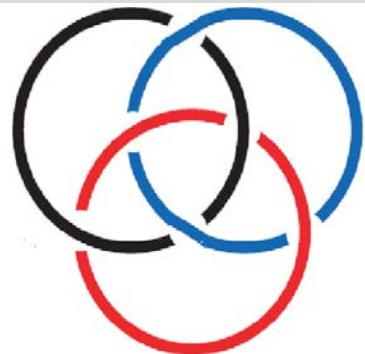
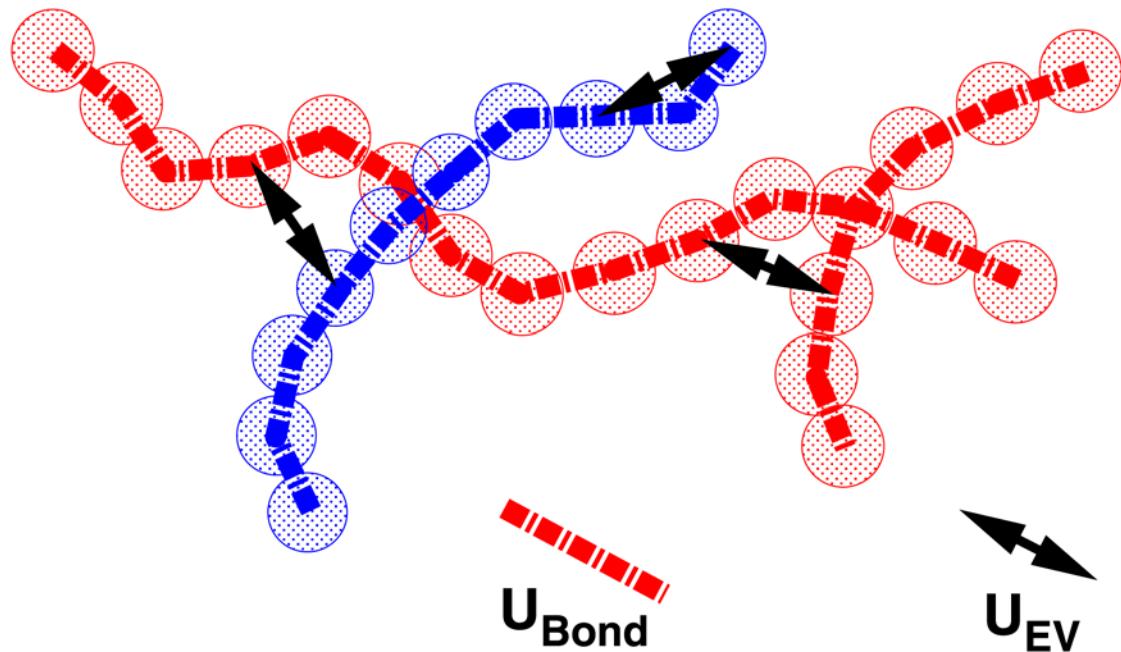
“Tube leakage”: Entropy penalty ($O(L/2)$) for defects too large

De Gennes 1971

$$d_T$$



(Semi) generic methods: MC, MD



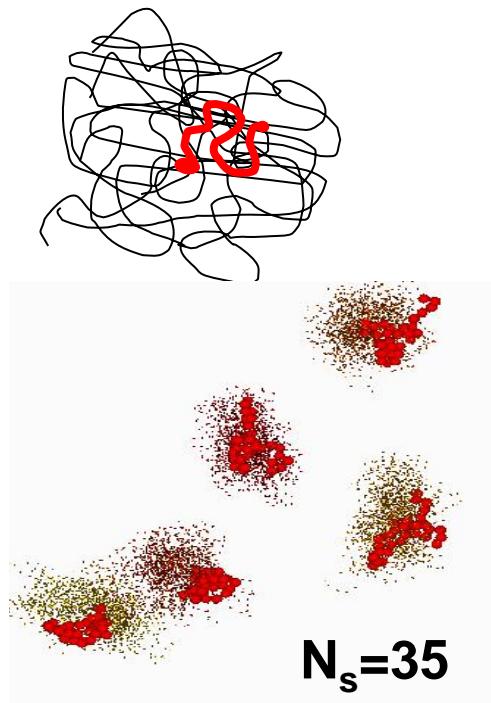
$$U_{EV}(r) = \epsilon \left(\left(\frac{\sigma}{r} \right)^{12} - \left(\frac{\sigma}{r} \right)^6 + \frac{1}{4} \right)$$

$$U_{\text{Bond}} = -30 k_B T \ln \left\{ 1 - \left(\frac{r}{R_c} \right)^2 \right\}$$

•Endlinked melts



Edgardo R. Duering, KK, Gary S. Grest
JCP, 1994



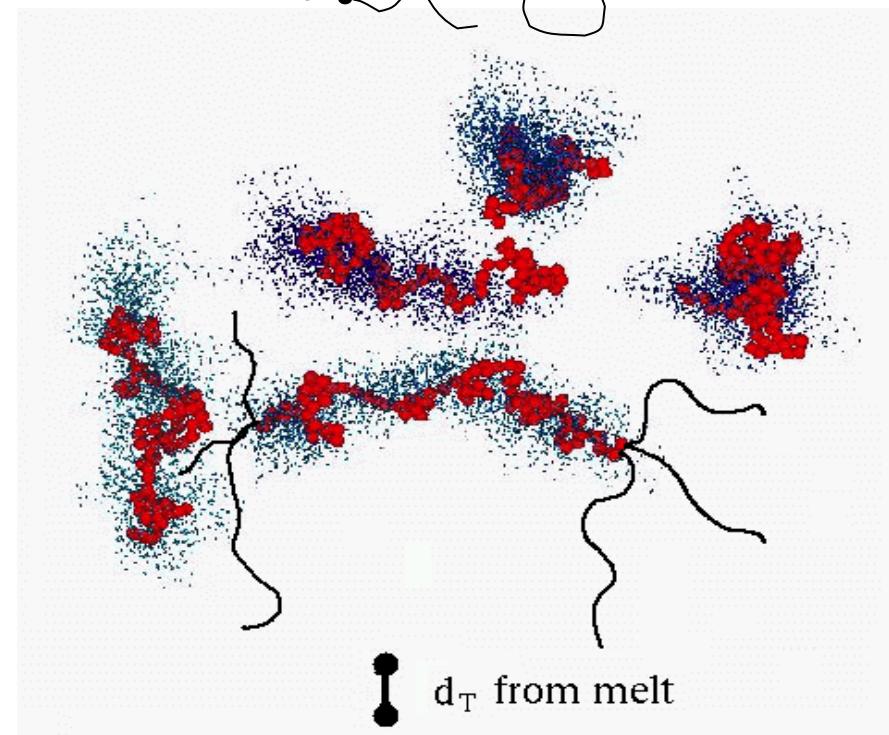
$$R^2 \odot d_T^2$$

$$N_e \odot 70$$



Initial conformation

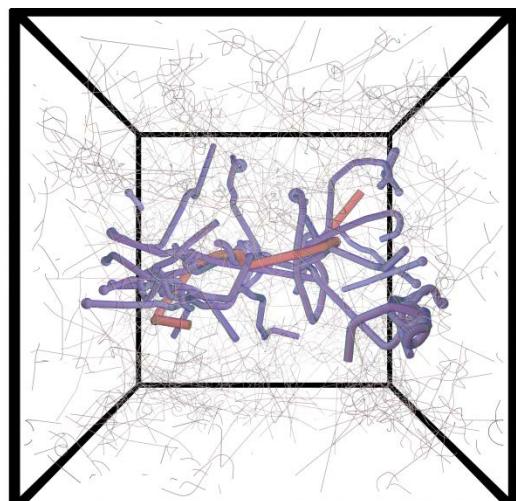
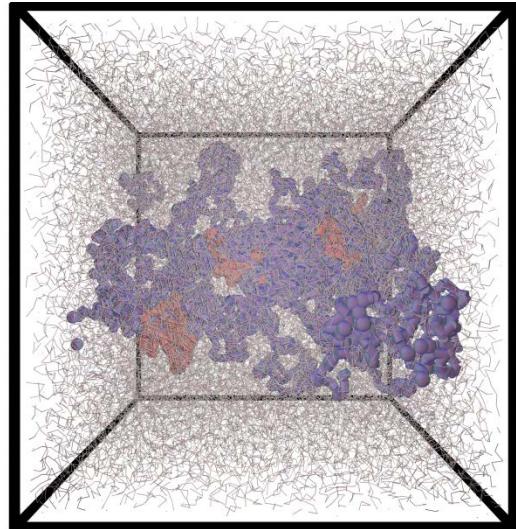
..... All conformations



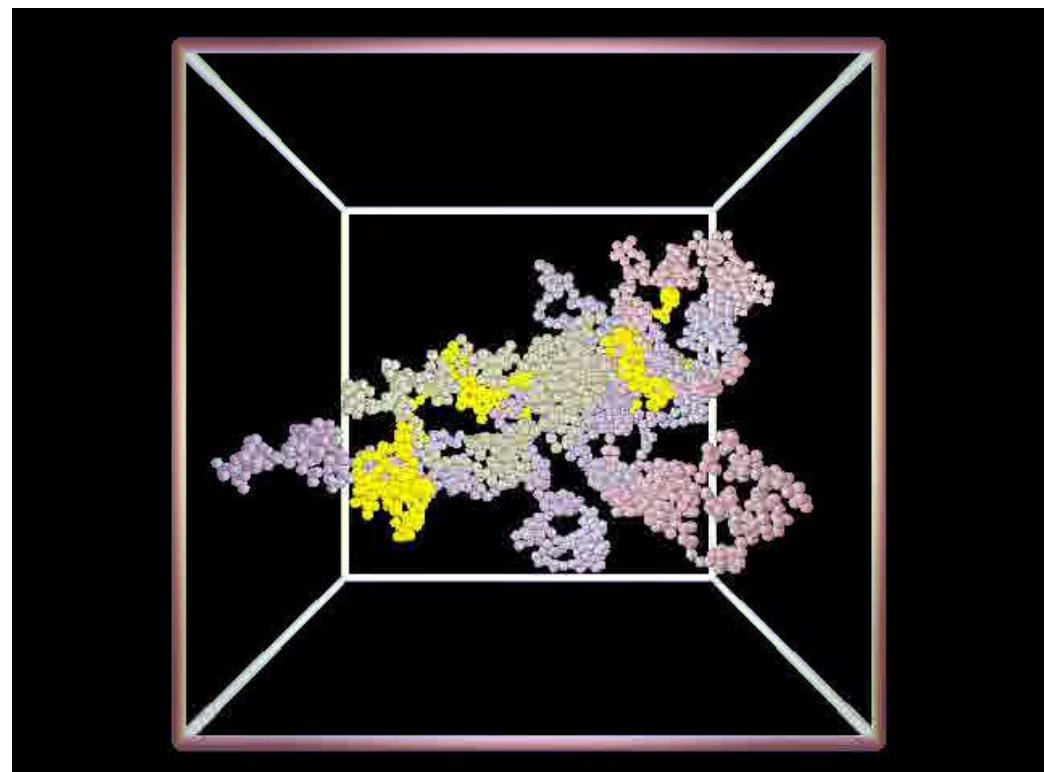
$$N = 100$$



Entanglements: Primitive Path Analysis



- Evolution of entangled chain cluster
(Everaers et al, Science 2004)

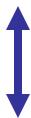


Related work: Theodorou, Kroeger, Larson, Briels...



Characteristic Lengths: l_K and p

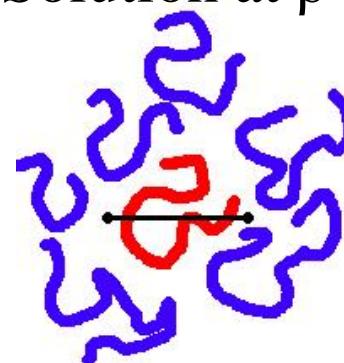
$$p = \frac{N}{\rho} \frac{1}{\langle R^2(N) \rangle}$$



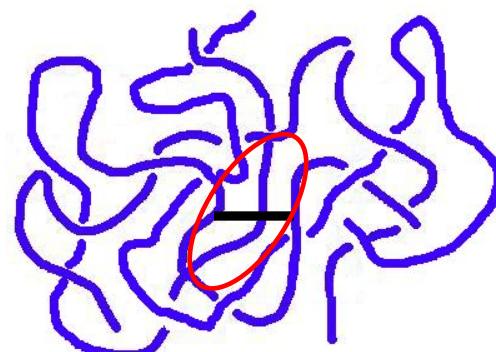
ρ/N = chain density

Packing
Length p

Solution at ρ^* : $p \approx R$



Melt: $p \approx$ strand-strand distance



Plateau Modulus G_N^0 Different Polymers

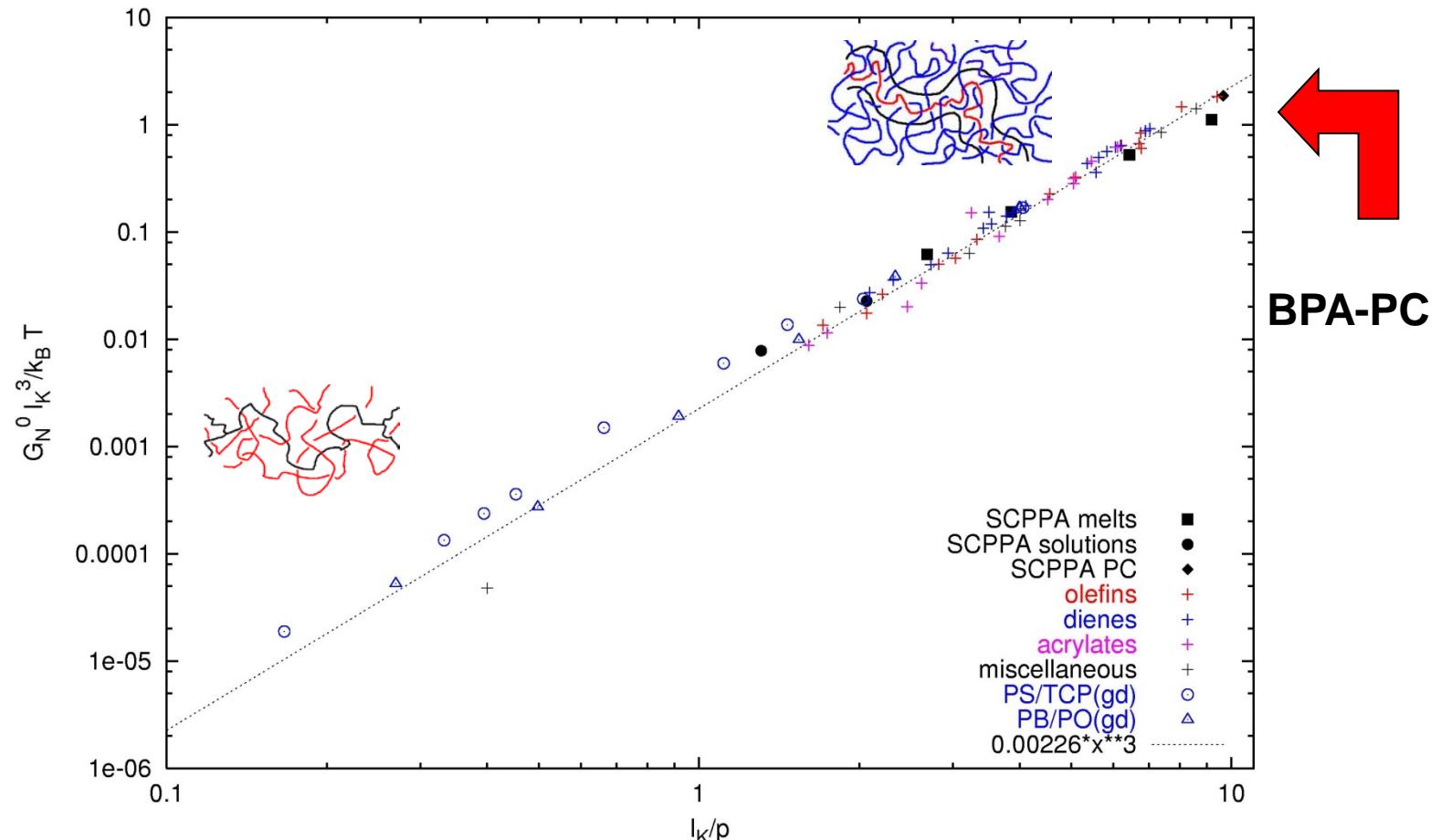
Melts and Solutions: Topological analysis and experiment



$G_N^0 l_k^3 / kT$ vs l_k / p

NO adjustable Parameters

GE plot (smdil lk= $\langle r^2 \rangle / L$)



Soft Matter

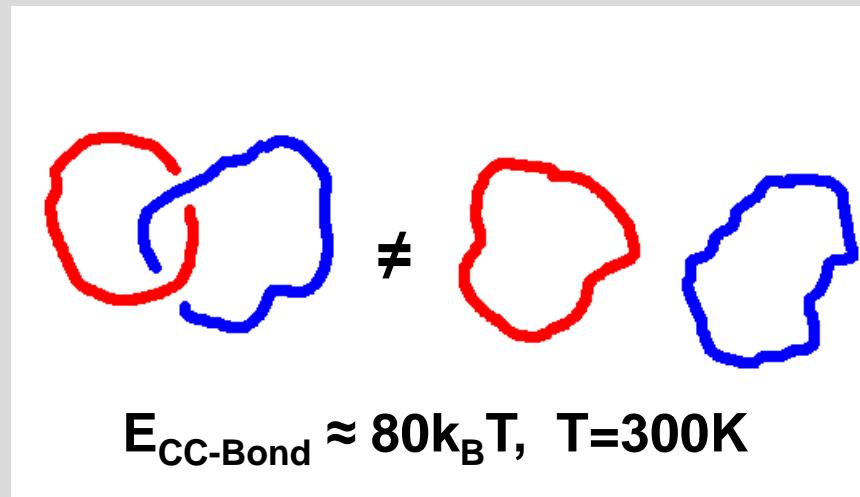


“Soft” means:

Large fluctuations,
conformational entropy, $O(L)$
(L contour length of the molecule)

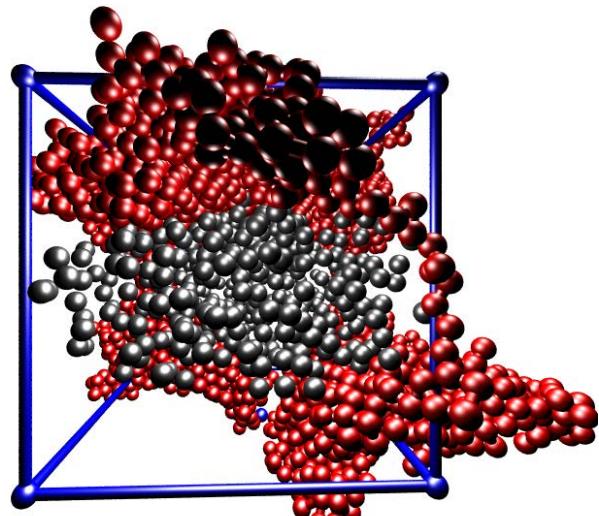
and/or

if energies $\gg k_B T$
fluctuations are constrained



Shrinking Gels: Poor Solvent Hydrogels

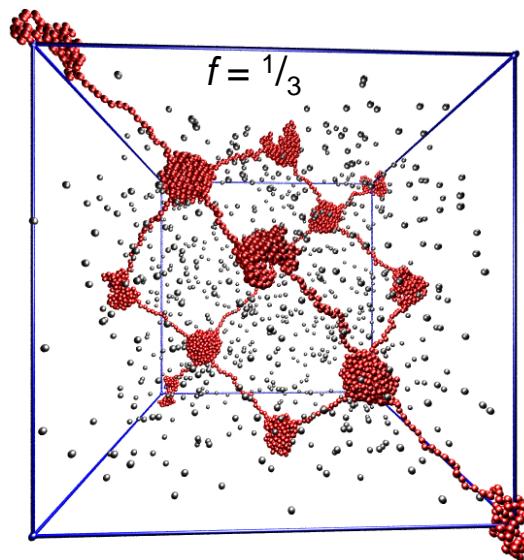
high salt
Low charge fraction



$$f = 1/8$$

...

low salt
high charge fraction

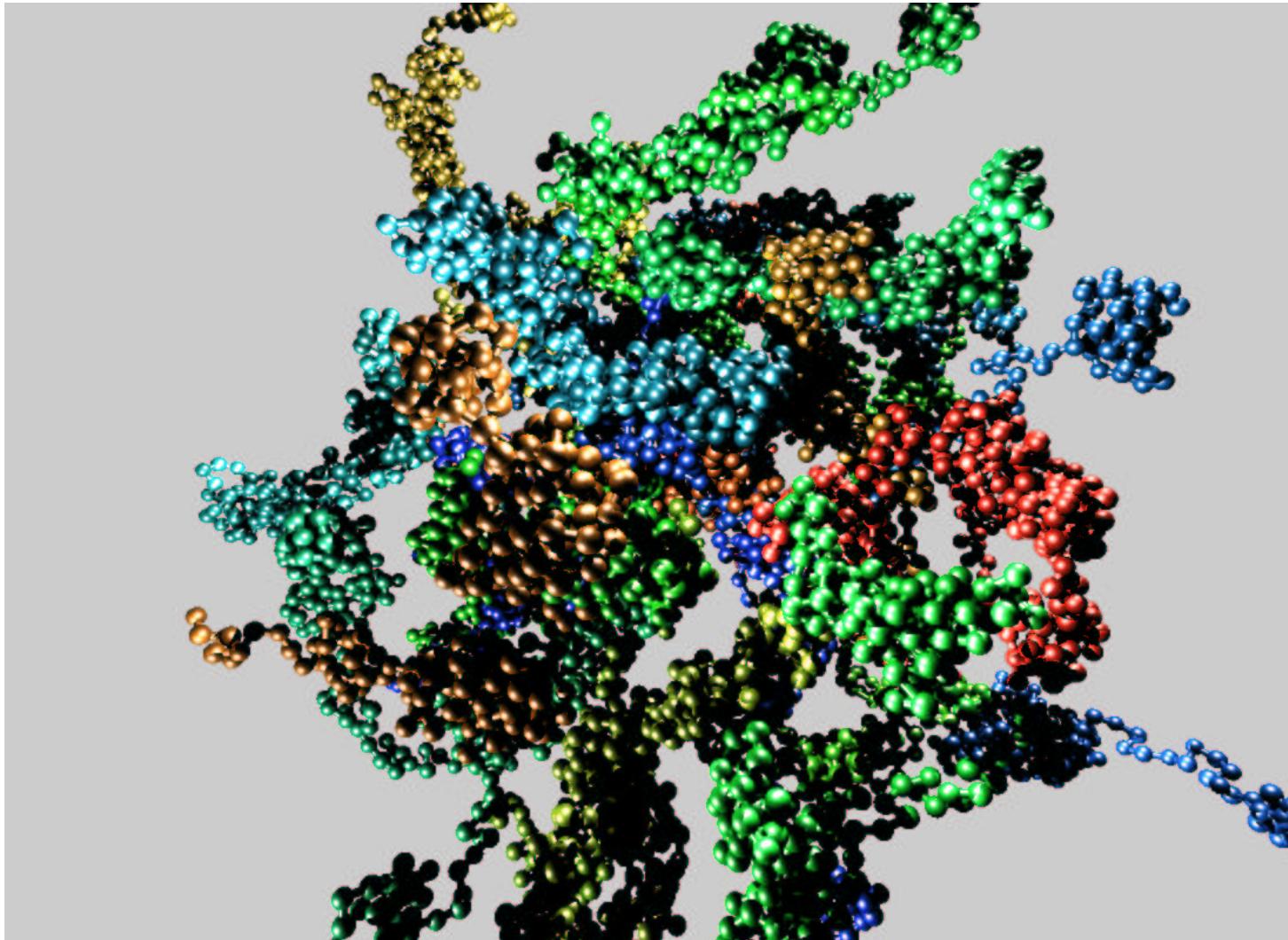


$$f = 1/1$$

Segregated strands?

- osmotic pressure of “counterions” enforces swelling (\leftrightarrow good solvent picture)
- nodes act as condensation nuclei
- “pearls” form due to balance with surface tension (\leftrightarrow Rayleigh instability)

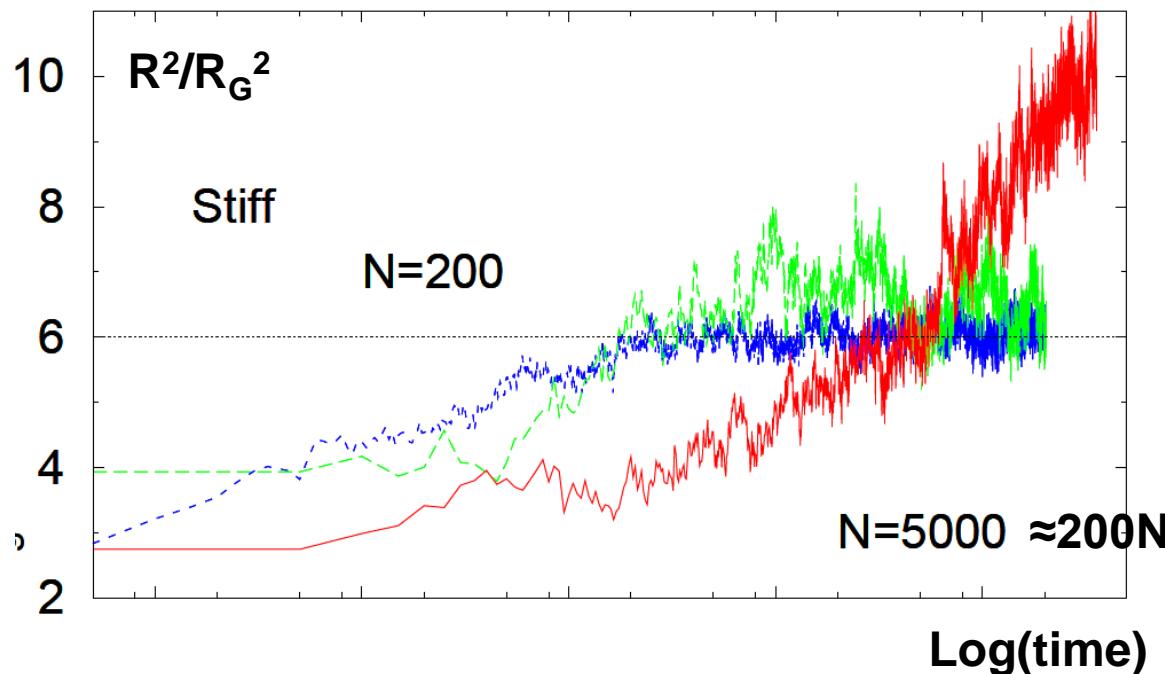
Collapsing polyelectrolyte chains





Melt of collapsed linear polymers: Make “use” topological constraints

T. Vettorel, KK, MTS 2010

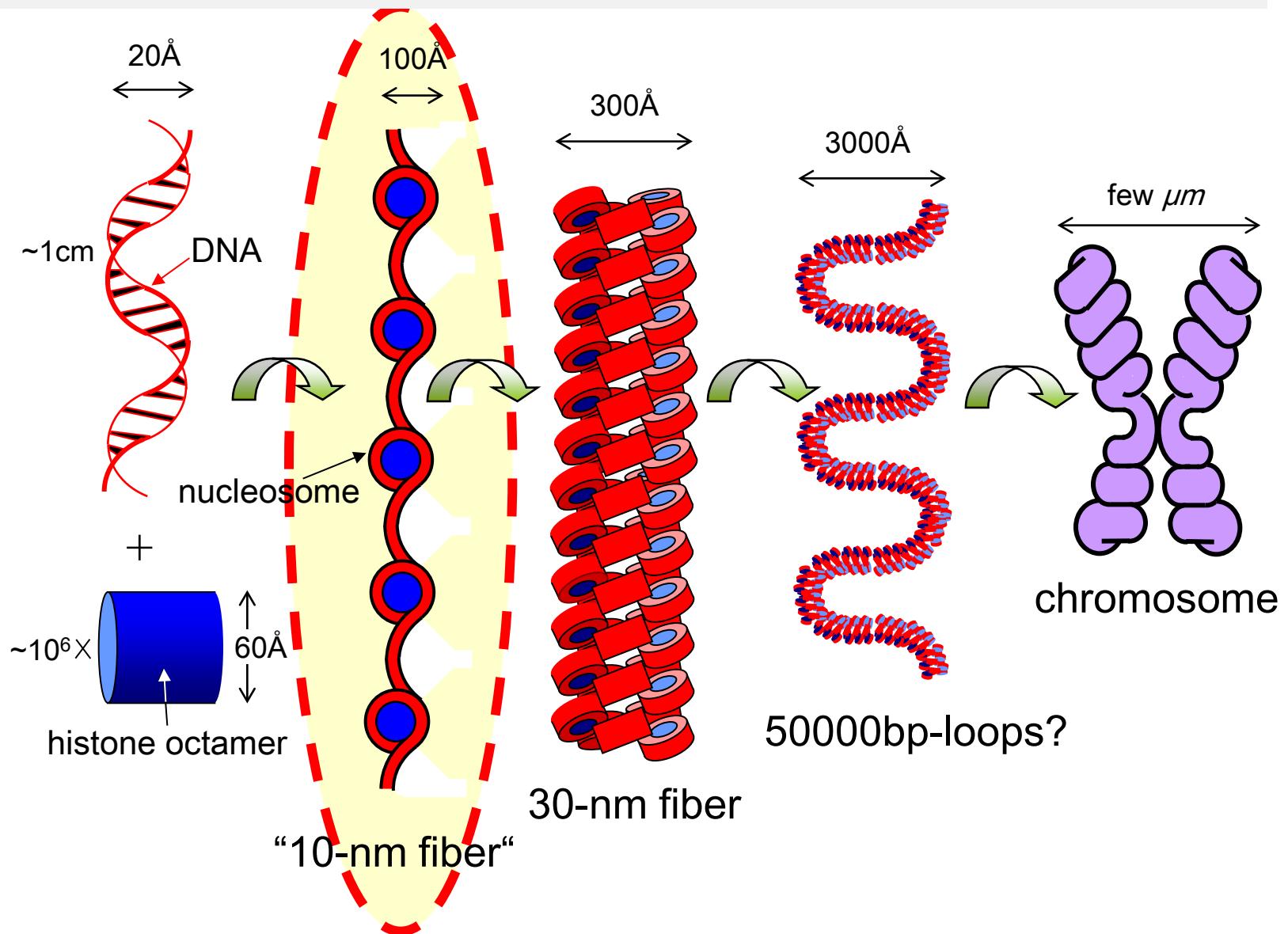


See also work on crystals: Rastogi, Spiess, McLeish



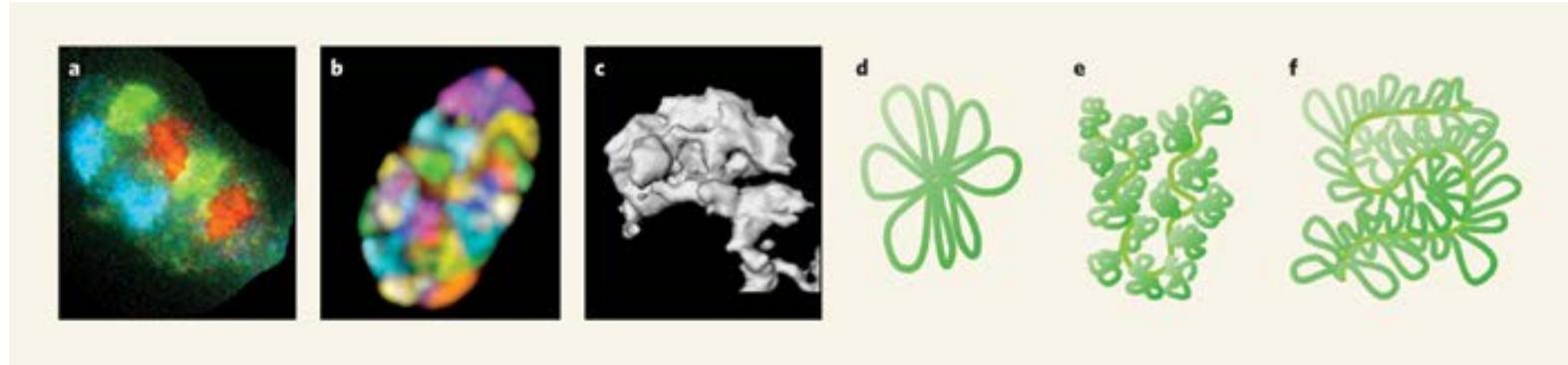
Structure of Chromatin

from review H. Schiessel, „The physics of chromatin“, J. Phys.: Condens. Matter 15 (2003) R699

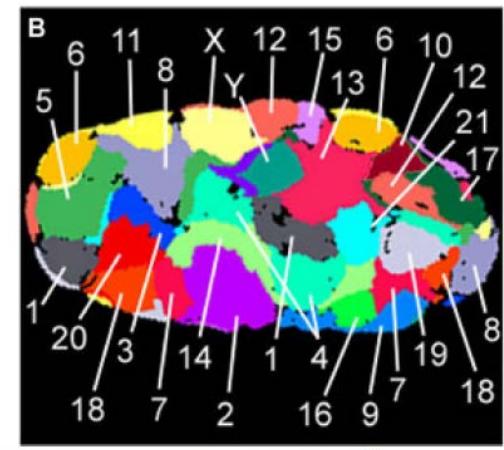
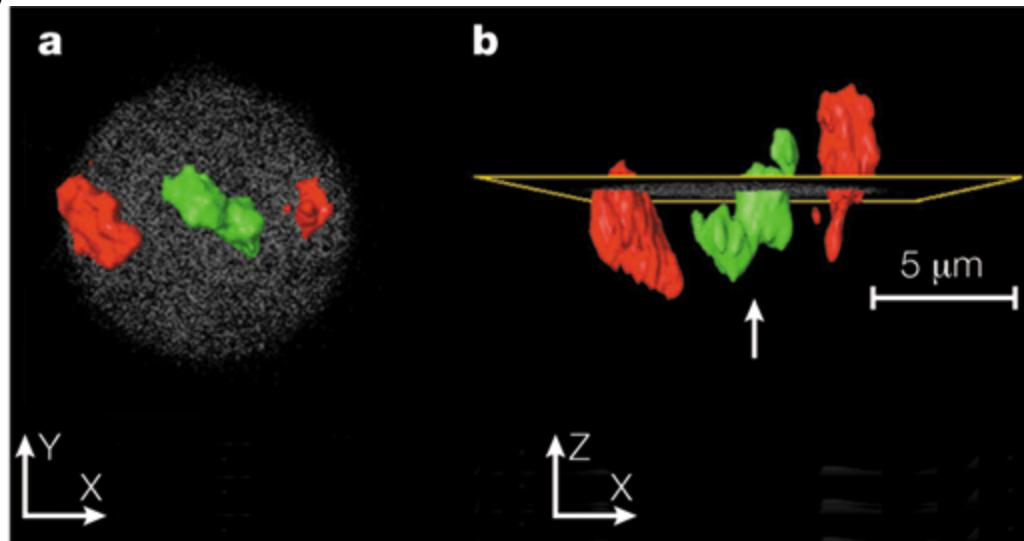


Chromosome Territories

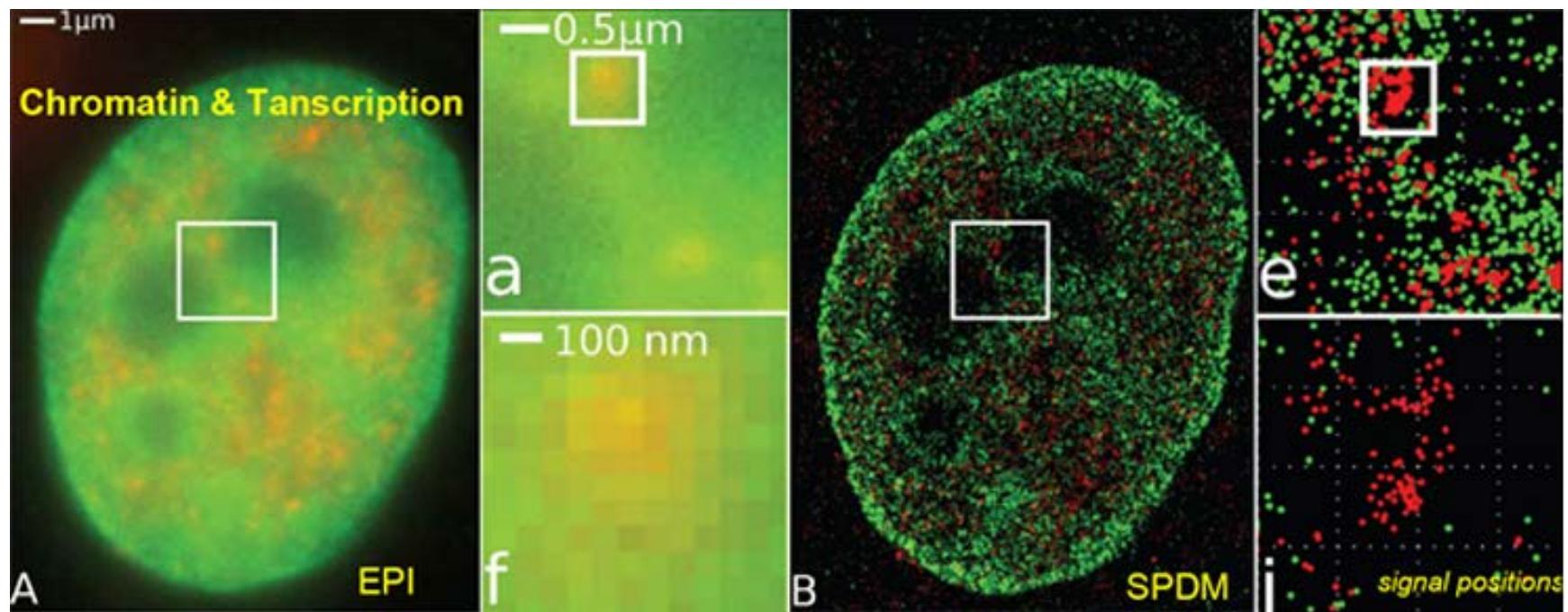
K.J. Meaburn, T. Mistelli, Nature 445, 379-381(2007)



T. Cremer, C. Cremer, Nature Reviews Genetics vol. 2, pp. 292-301
(2001)

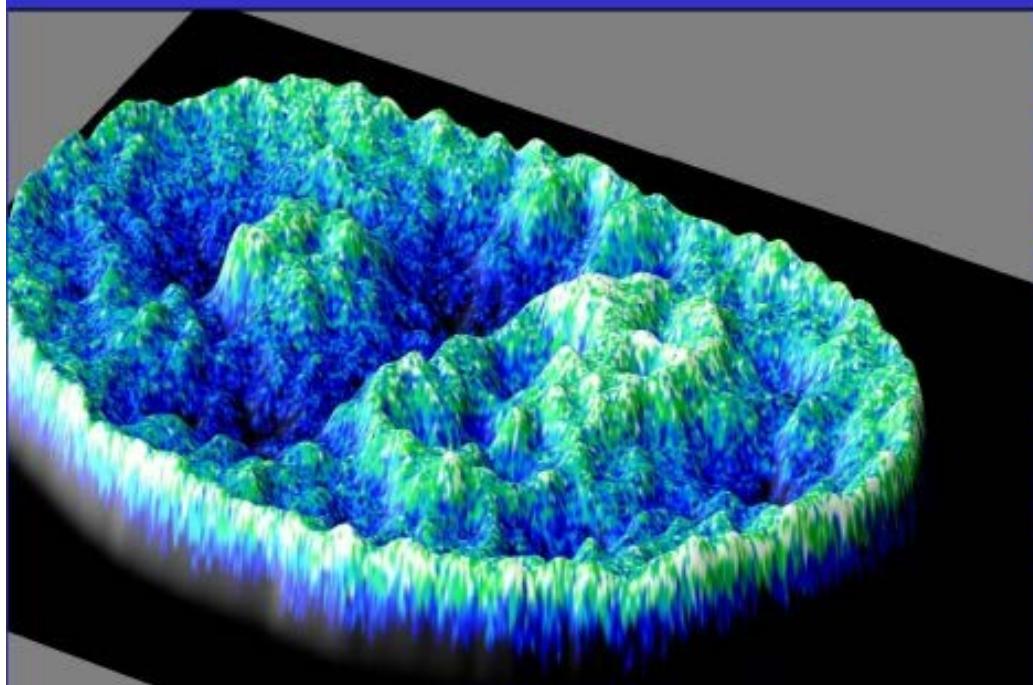


Cremer Lab, IMB Mainz



Cremer Lab, IMB Mainz

Superresolution Microscopy* of DNA Distribution in a human Cell Nucleus.



Intensity Image
after Labelling
with a DNA
specific
Fluorochrome:

~ 5,000 DNA
sites / μm^2

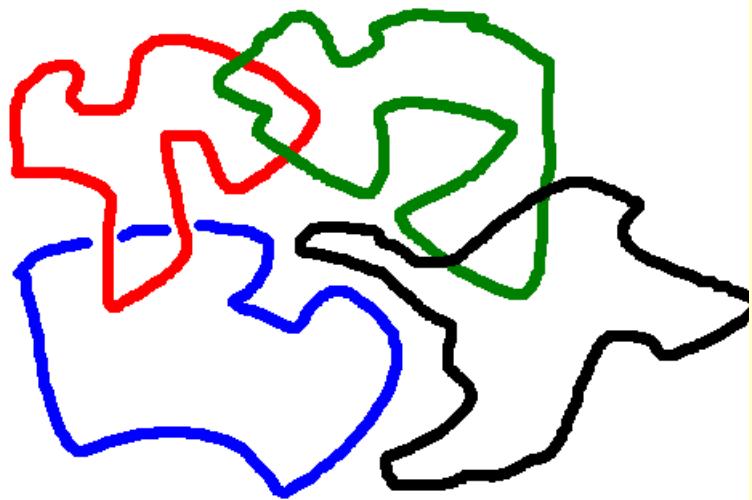
Maximum ~ 27,000
DNA sites/ μm^2

Szczurek et al. (2014)

*Localization Microscopy (SPDM)



Melt of non concatenated ring polymers



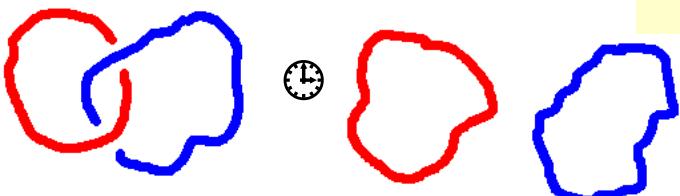
Conformation of the rings?

Random walk: $R^2(N) \propto N^{2v}$, $v=1/2$

=> Rings interpenetrate each other

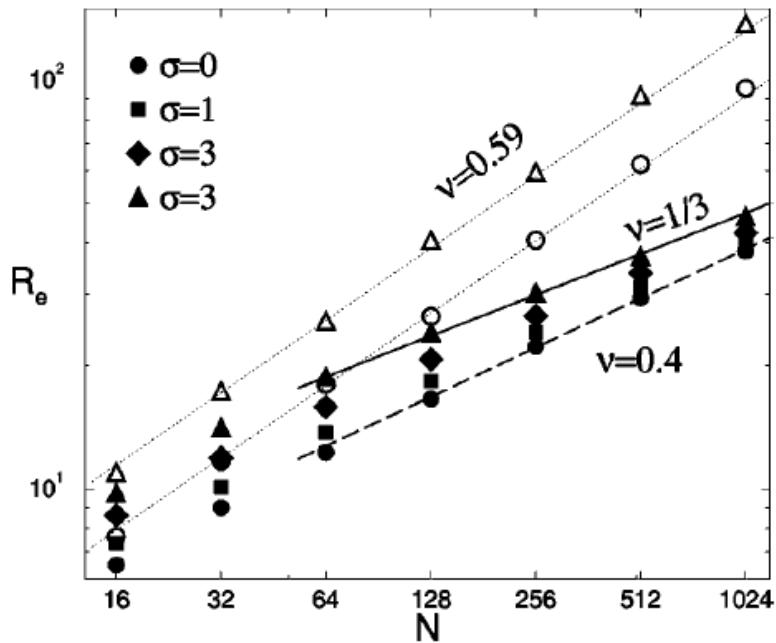
Collapsed ring: $R^2(N) \propto N^{2v}$, $v = 1/3$

⇒ Rings segregate,
finite number of neighbors





Topological effects in ring polymers. II. Influence of persistence length

M. Müller,^{1,2} J. P. Wittmer,^{2,3,*} and M. E. Cates²

Speculation of higher order contributions

$$F \propto kT(R^d / N)^\alpha$$

Osmotic pressure

$$+ kTN / R^2$$

Surface term of EV interaction



$$\nu = (\alpha + 1)/(3\alpha + 2)$$

$$\alpha = 1 \rightarrow \nu = 0.4$$

$$\alpha = \infty \rightarrow \nu = 1/3$$

New Flory theory, based on lattice animal concept
 $\nu=1/3$, $d=3$
 A. Y. Grosberg,
 Soft Matter Nov. 2013

Too(?) simple Physical picture

Grosberg, Vettorel, KK

- If rings collapse, (naively) expect core plus corona

(first simple picture)

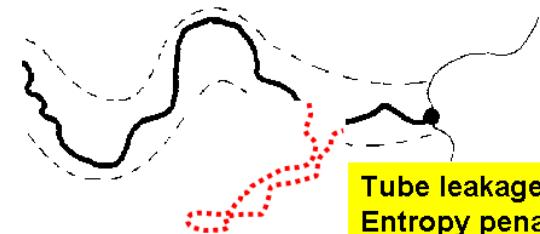


$$d_T^2 \approx N_e$$

$$N \leq N_e$$

$$N_e \leq N \leq O(\approx 10N_e)$$

$$N \geq O(\approx 10N_e)$$



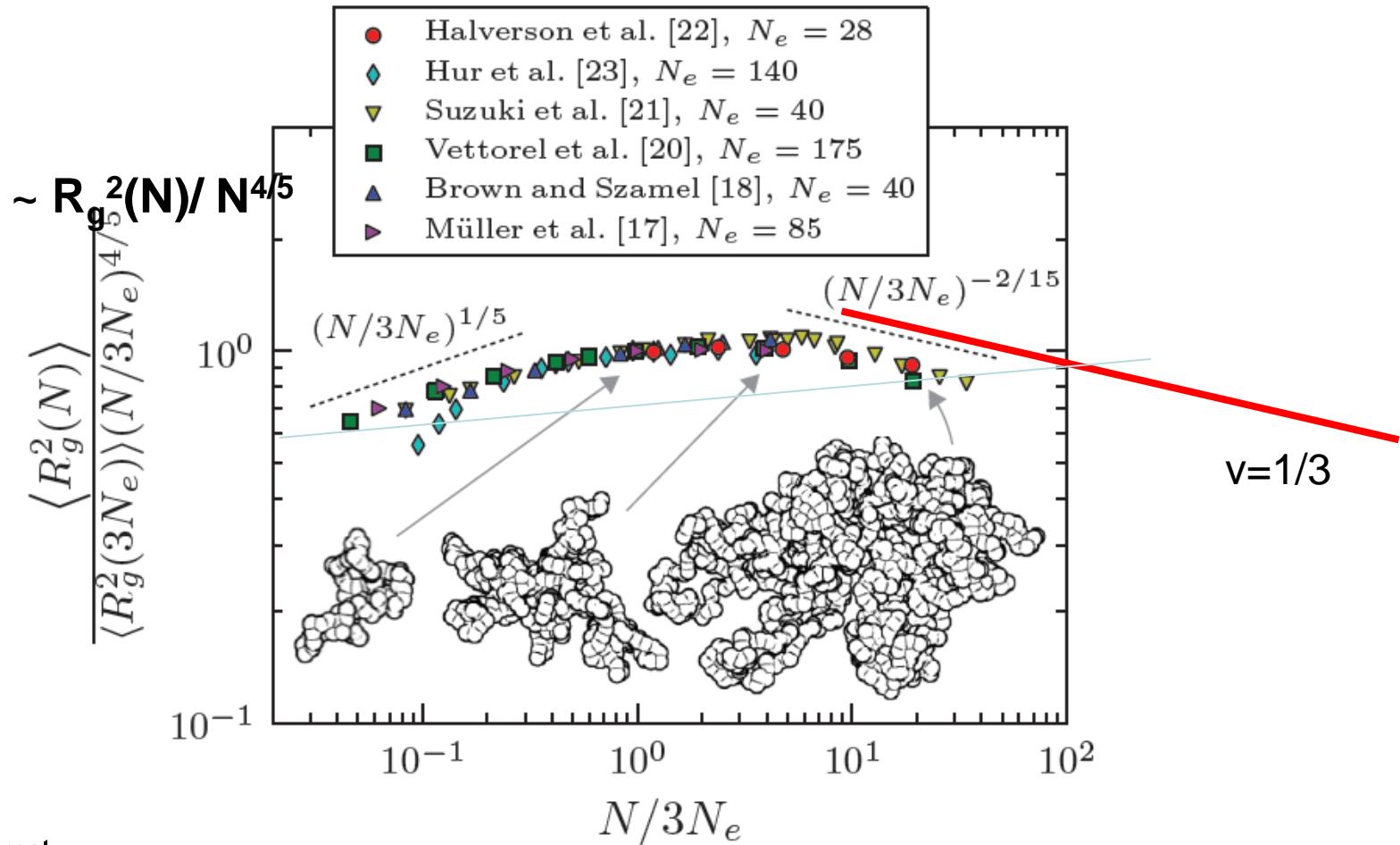
Tube leakage:
Entropy penalty for
defects too large

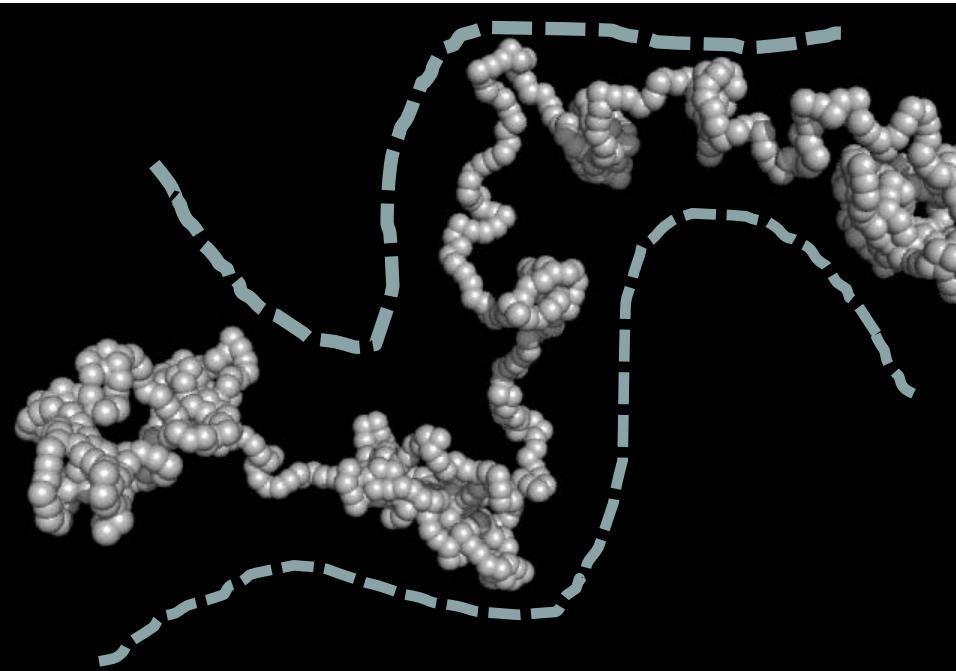
No effect, random walk

Begin “shrinking”,

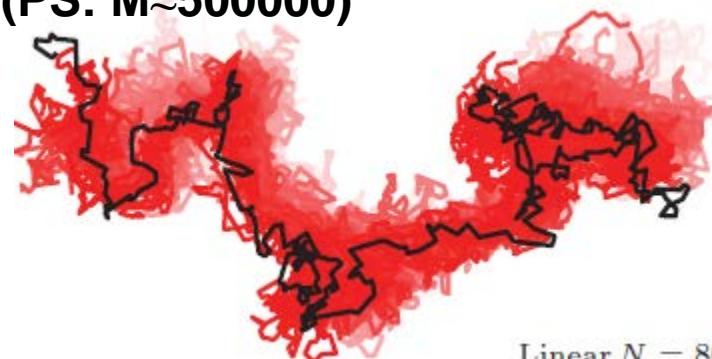
Segregation regime

Ring extensions: Collection of data



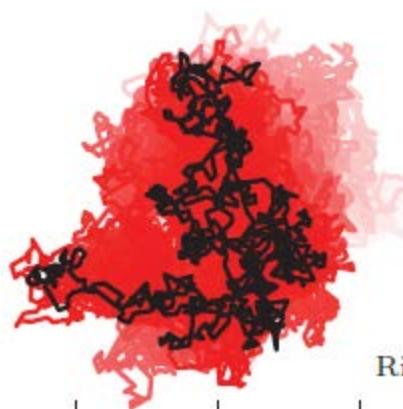


Linear, $N=800 \approx 29 N_e$
(PS: $M \approx 500000$)

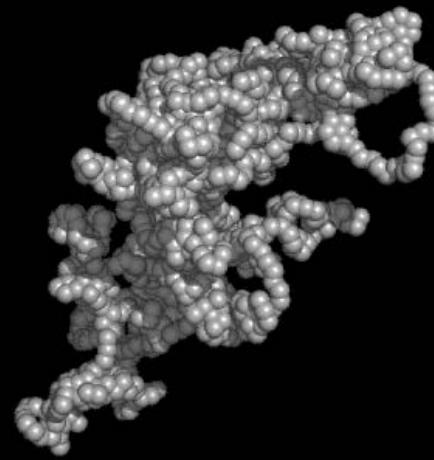


Linear $N = 800$

chain id = 78, monomers per chain = 800
time = $5.664e+06$ tau $Rg(t)/\langle Rg \rangle = 1.22$ $\sqrt{g_3(t)} / \langle Rg \rangle = 0.00$



Ring $N = 1600$
(PS: $M \approx 1000000$)



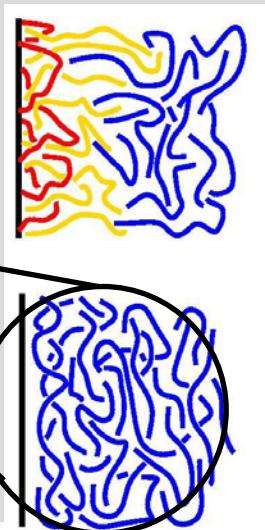
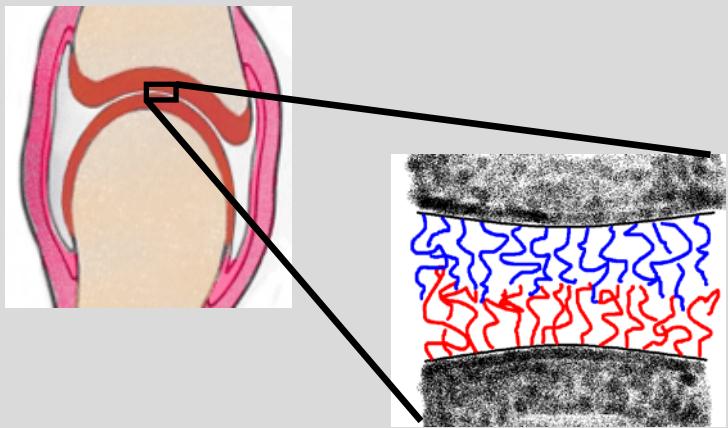
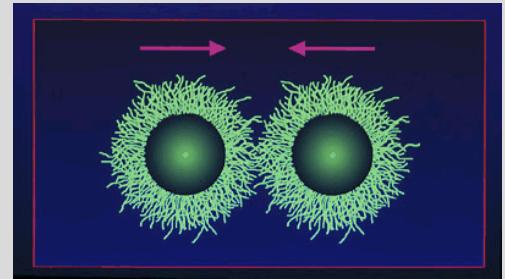
chain id = 29, monomers per ring = 1600
time = $1.00e+04$ tau $Rg(t)/\langle Rg \rangle = 1.13$ $\sqrt{g_3(t)} / \langle Rg \rangle = 0.00$

Surfaces



Back to Spin analogy:

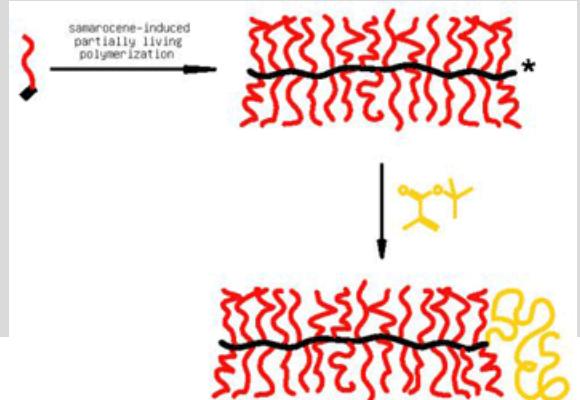
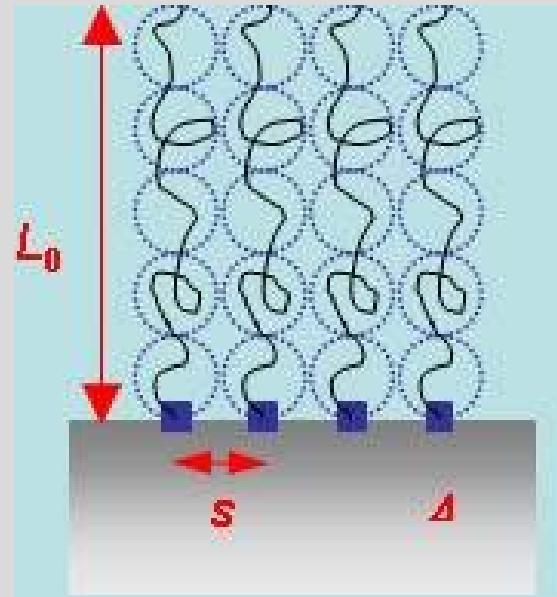
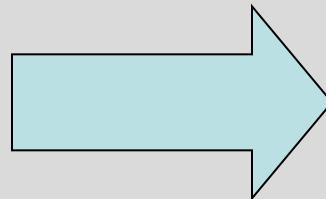
- Adsorption of polymers, wetting of polymers
- Adsorption of a single chain – relation to surface magnetism



Soft Matter - Polymers

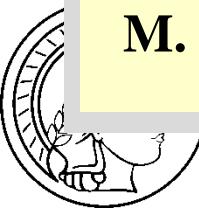


Many chains close to surfaces: Brushes



Bottle – Brushes

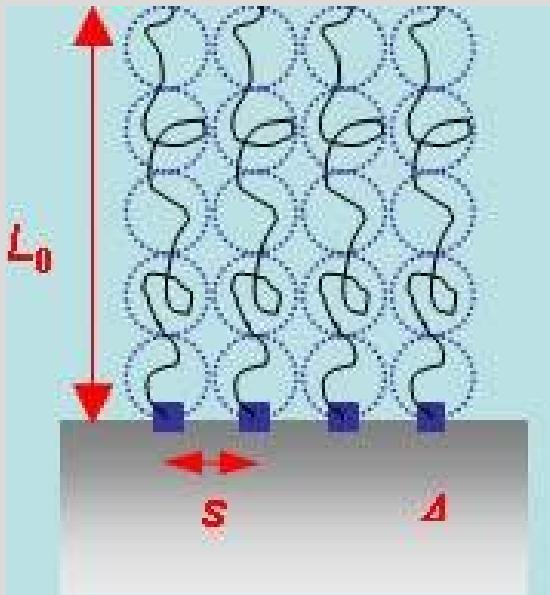
M. Schmidt et al, K. Binder, H.P. Hsu, W.
Paul et al, SFB625



Soft Matter - Polymers



Many chains close to surfaces: Brushes

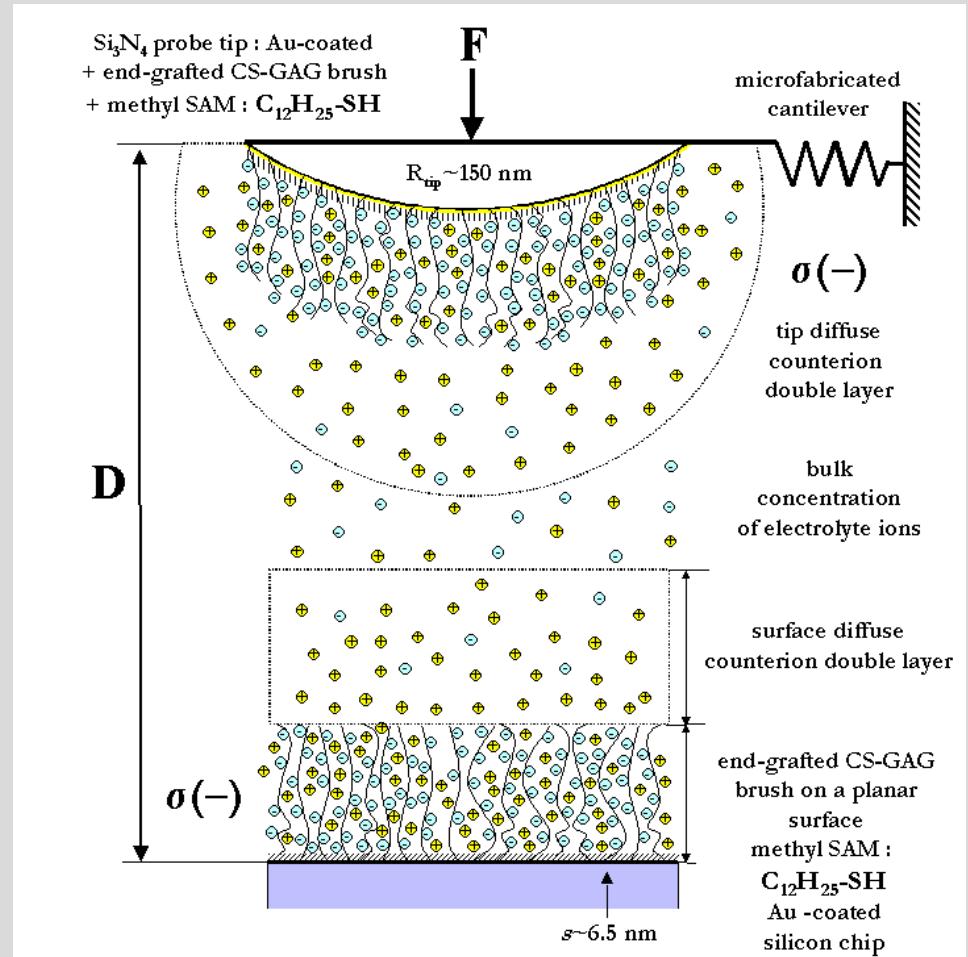


e.g. J. Klein et al

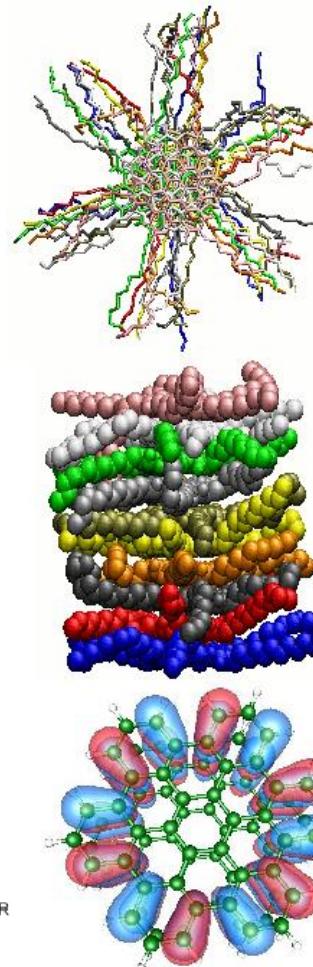
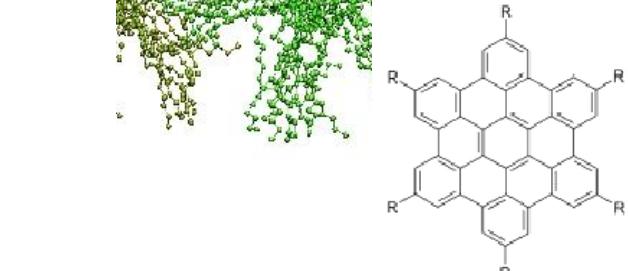
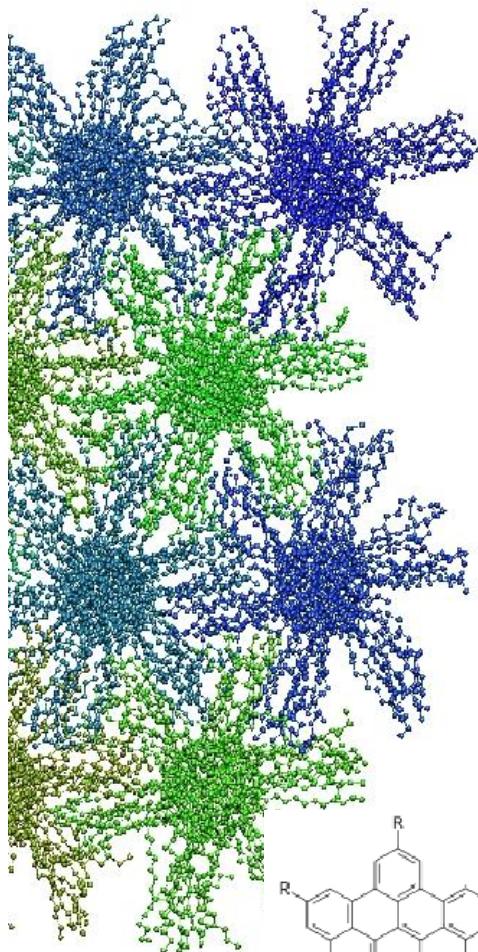
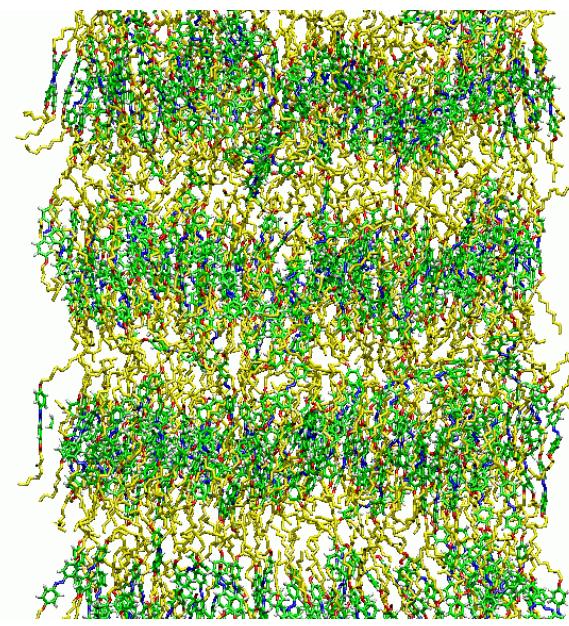
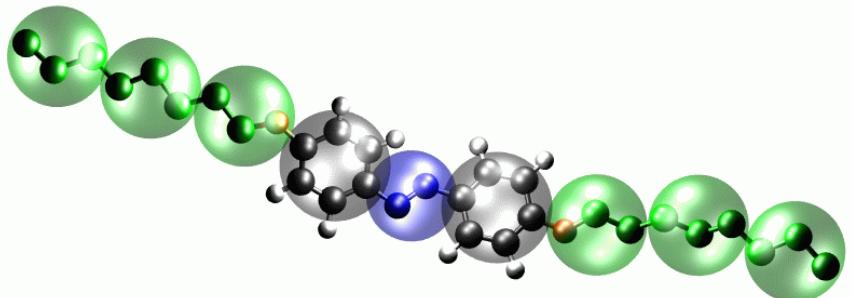
Forces between surfaces etc.

Wetting phenomena ...

SCF Theory, simulations...



Goal: Application to complex systems



C. Peter et al, Soft Mat. 2008

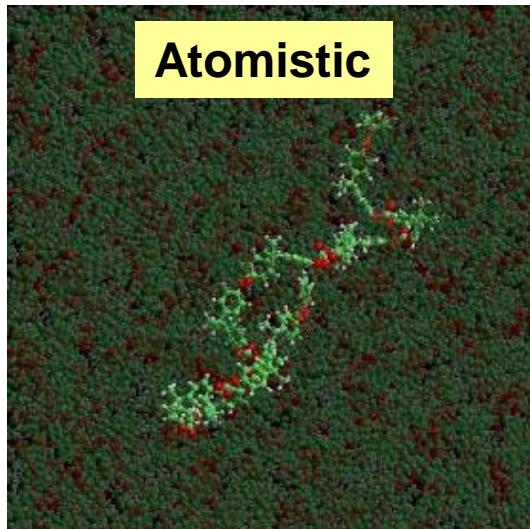
Andrienko et al, PRL 2007, Nature Mat. 2009

Polymers

relevant interactions well defined

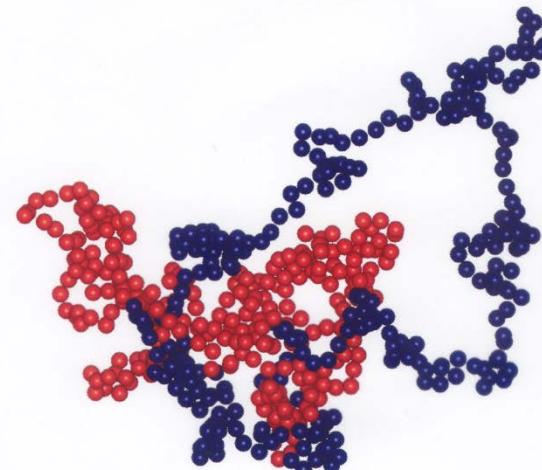


Atomistic

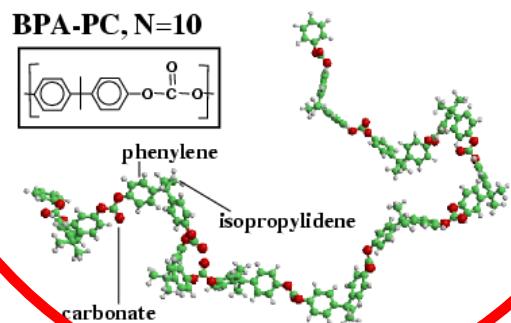


versus

Scaling Laws



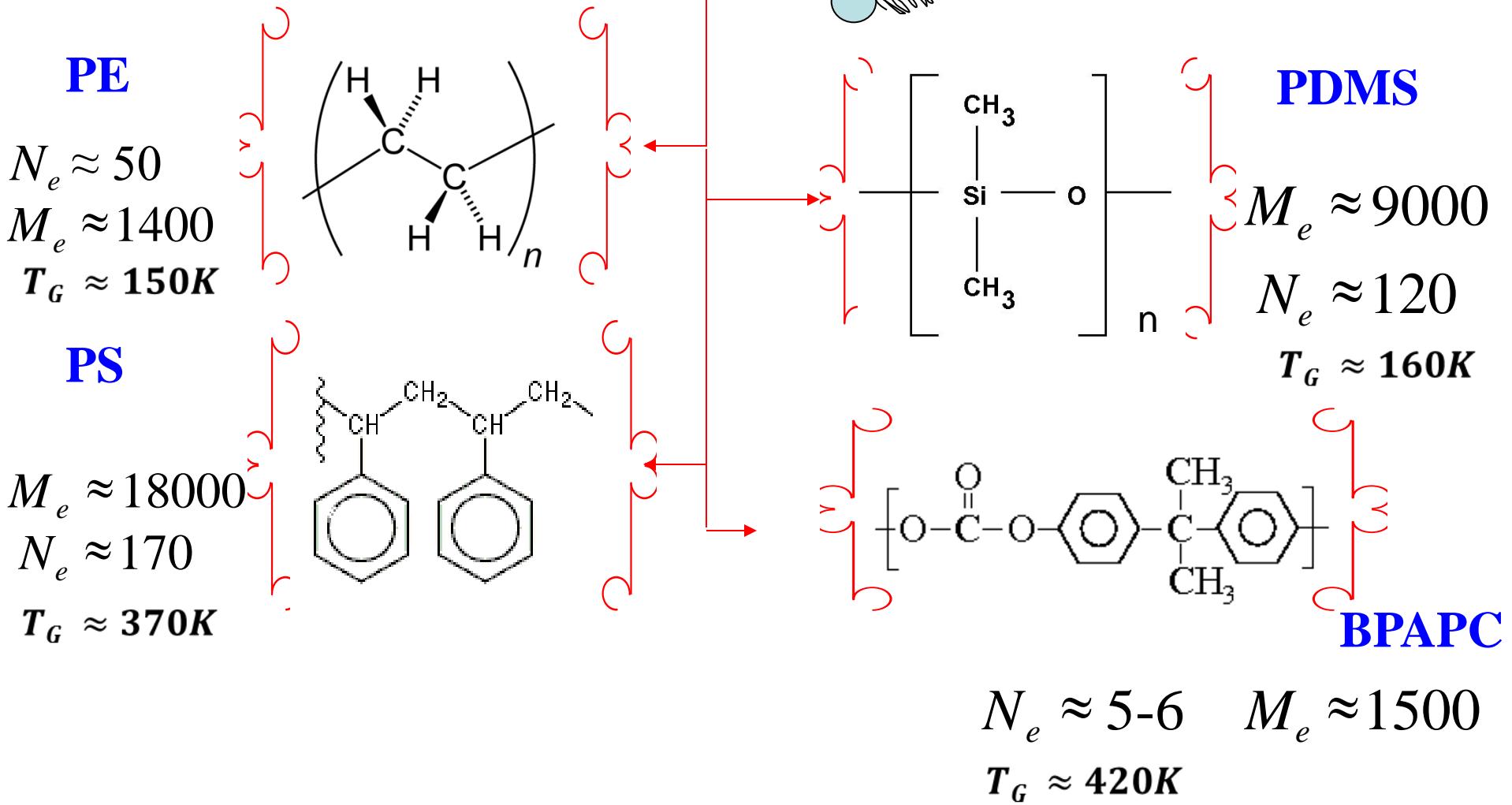
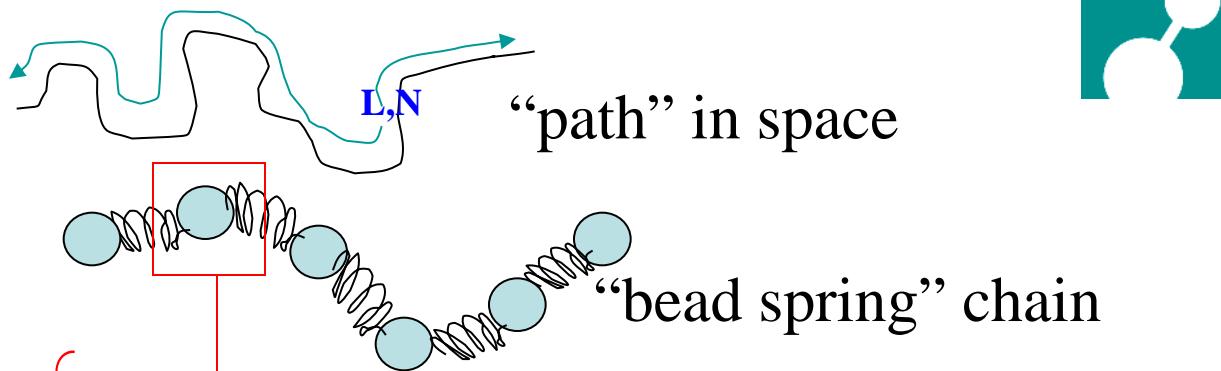
Generic



Structure Property Relations

Polymers

"oil soluble"



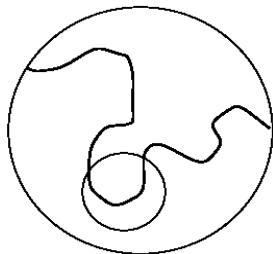
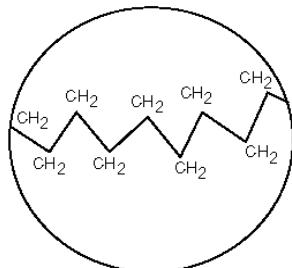
Glass transition temperature T_G of polymers

Polymer	T_G (K)
•	
• Polyethylene (LDPE)	150K
• <u>Polypropylene</u> (PP, atactic)	250K
• <u>Poly(vinyl acetate)</u> (PVAc)	300K
• <u>Polyethylene terephthalate</u> (PET)	350K
• <u>Poly(vinyl alcohol)</u> (PVA)	360K
• <u>Poly(vinyl chloride)</u> (PVC)	355K
• <u>Polystyrene</u> (PS)	370K
• <u>Polypropylene</u> (PP, isotactic)	273K
• <u>Poly(3-hydroxybutyrate)</u> (PHB)	273K
• <u>Poly(methylmethacrylate)</u> (PMMA, atactic)	380K
• <u>Polycarbonate</u> (BPA –PC)	420K
• HIP(TMC)- Polycarbonate	$\approx 520K$

Most polymers are glassy at room temperature



Example Viscosity $\eta=AN^x$ of a polymer melt (extrusion processes)



Microscopic

$$L \approx 1\text{\AA} - 3\text{\AA}$$

$$T \approx 10^{-13} \text{ sec}$$

materials/ chemistry specific Prefactor

(e.g. function of T_G , glass transition)

$$\eta = A N^x$$

“Energy dominated”

Mesoscopic

$$L \approx 10\text{\AA} - 50\text{\AA}$$

$$T \approx 10^{-8} - 10^{-4} \text{ sec}$$

generic/universal Properties

$$\eta = A N^x \quad X = 3.4$$

M molecular weight

“Entropy dominated”

$$\boxed{\eta = A N^x}$$

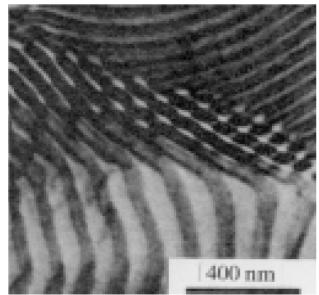
varies many decades

e.g.: increase $\eta \Rightarrow 10\eta$

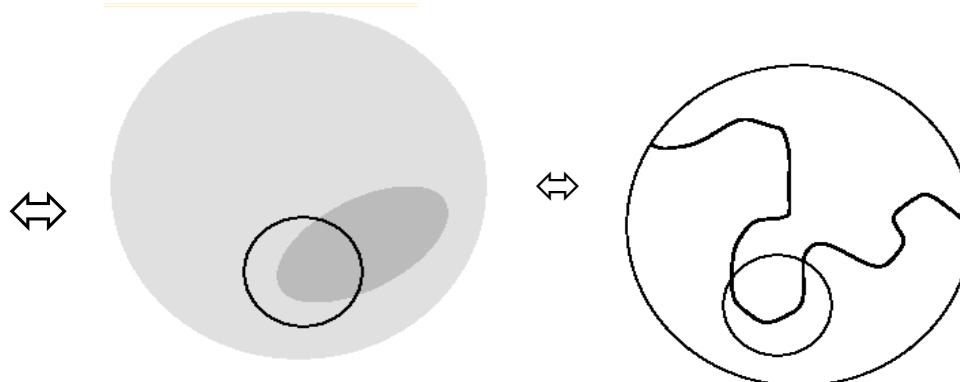
specific: $T=500\text{K} \Rightarrow 470\text{K}$ (BPA-PC)

generic: $N \Rightarrow 2N$

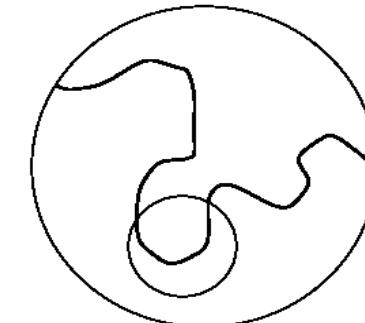
Time and length scales



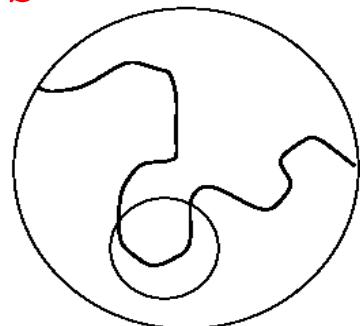
Macroscopic
domains etc.



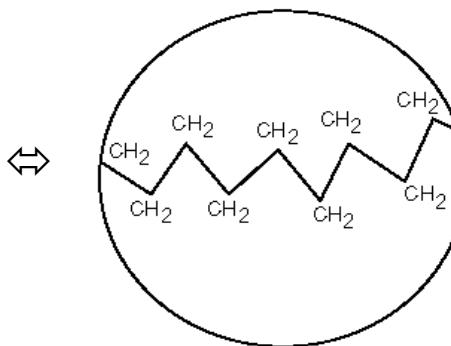
Semi macroscopic
 $L \approx 100\text{\AA} - 1000\text{\AA}$
 $T \approx 0$ (1 sec)



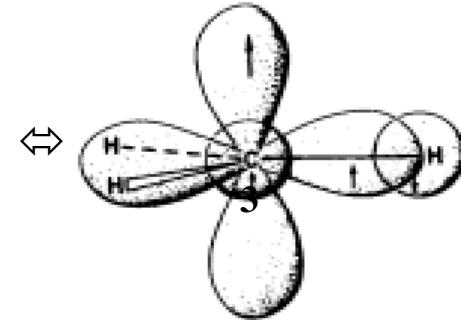
Mesoscopic
 $L \approx 10\text{\AA} - 50\text{\AA}$
 $T \approx 10^{-8} - 10^{-4}$ sec
Entropy dominates



Mesoscopic
 $L \approx 10\text{\AA} - 50\text{\AA}$
 $T \approx 10^{-8} - 10^{-4}$ sec
Entropy dominates

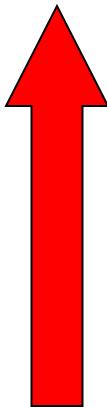


Microscopic
 $L \approx 1\text{\AA} - 3\text{\AA}$
 $T \approx 10^{-13}$ sec
Energy dominates



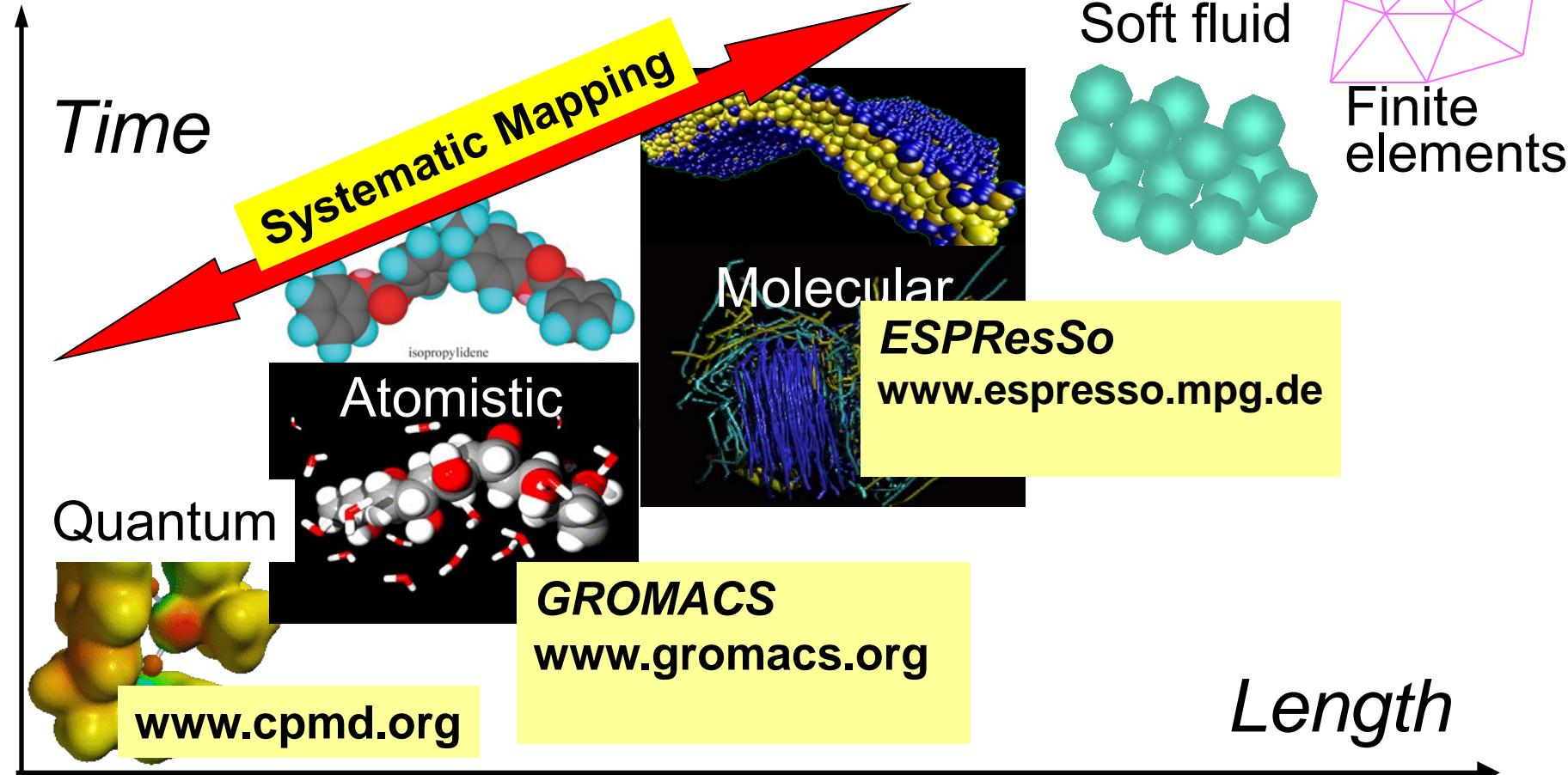
(Sub)atomic
electronic structure
chemical reactions
excited states

Material Properties



generic/universal AND chemistry specific

Characteristic Time and Length Scales



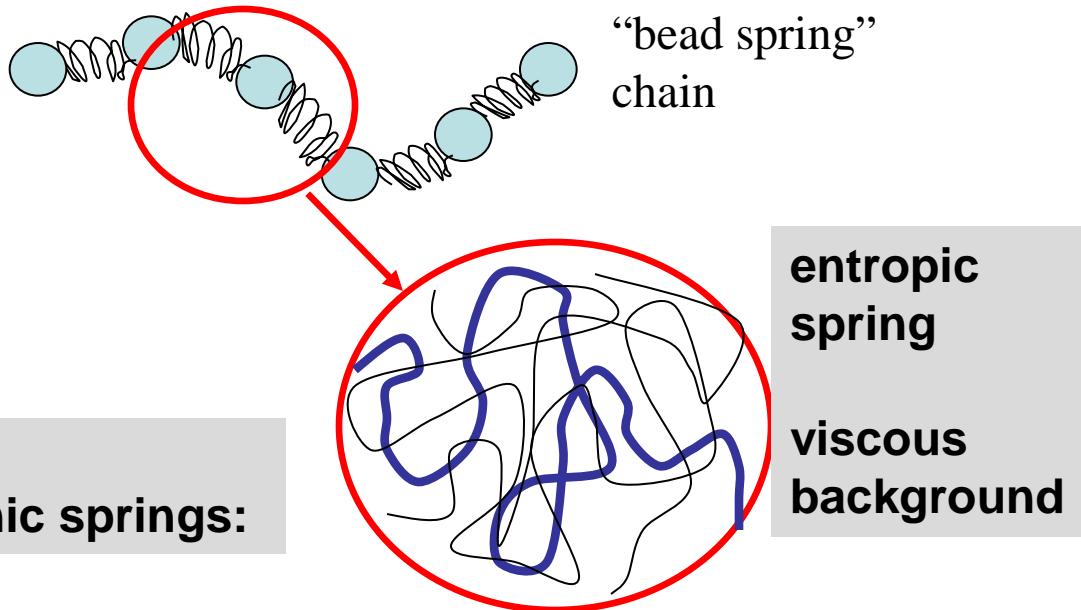
Local Chemical Properties --- Scaling Behavior of Nanostructures
Energy Dominance --- Entropy Dominance of Properties

Detour: Rouse Reptation – Most Basic Aspects



“Rouse Model”₁₉₅₃

Brownian dynamics of
chain (random walk) of harmonic springs:



$$\zeta \dot{r}_i = \frac{3kT}{l^2} (2r_i - r_{i-1} - r_{i+1}) + f_i$$

$$\langle f_i(t) f_j(t') \rangle = 2dkT \delta_{ij} \delta(t - t')$$



Rouse model

Rouse
Modes

$$\mathbf{X}_p(t) = \frac{1}{N-1} \sum_{i=1}^{N-1} r_i(t) \cos\left(\frac{p\pi(i+1/2)}{N-1}\right)$$

$$p = 0, 1, 2, 3, \dots$$

Equation of
motion

$$\zeta_p \dot{\mathbf{X}}_p = k_p \mathbf{X}_p + f_p$$

Relaxation time

$$\boxed{\tau_p} = \frac{\zeta \langle R^2(N) \rangle^2}{3\pi^2 l l_K k_B T p^2} \propto \boxed{\frac{N^2}{p^2}}$$

$$\tau_{Rouse} = \tau_{p=1}$$



Rouse model

Modulus

$$G(t) = \frac{\rho k_B T}{N} \sum_p \exp(-2tp^2 / \tau_R)$$

Viscosity

$$\eta = \int_0^\infty G(t) dt = \rho \zeta \frac{\langle R^2 \rangle}{36} \propto N$$

Chain diffusion constant

$$D = \frac{k_B T}{N \zeta}$$

Agrees to experiments of short chains



Rouse model

Modulus

$$G(t) = \frac{\rho k_B T}{N} \sum_p \exp(-2tp^2 / \tau_R)$$

Networks!!
(classical)

viscosity

$$\eta = \int_0^\infty G(t) dt = \rho \zeta \frac{\langle R^2 \rangle}{36} \propto N$$

Chain diffusion constant

$$D = \frac{k_B T}{N \zeta}$$

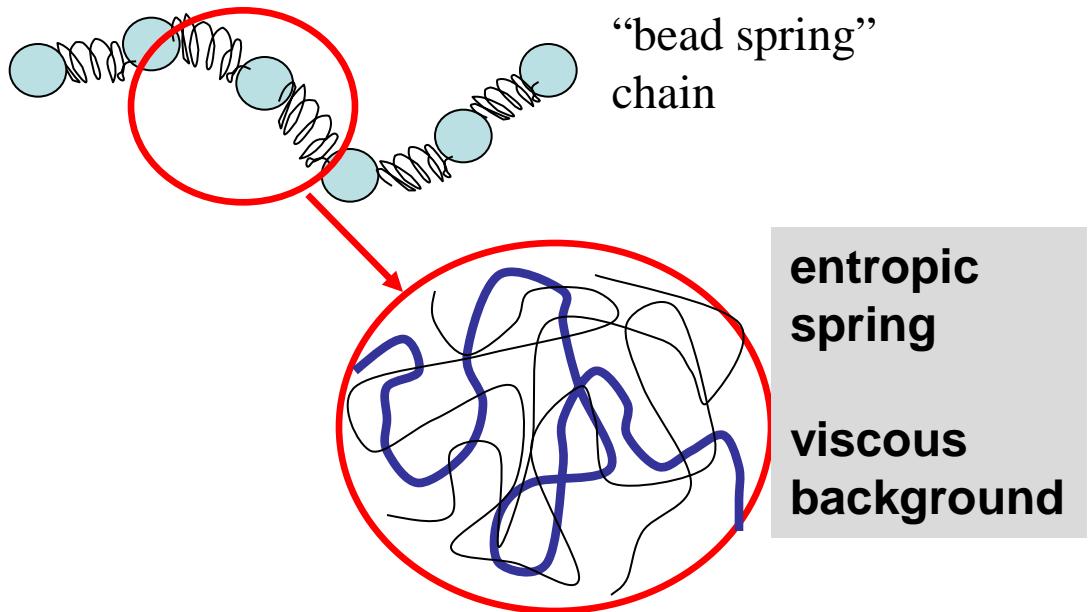
Agrees to experiments of short chains



Mean square displacements

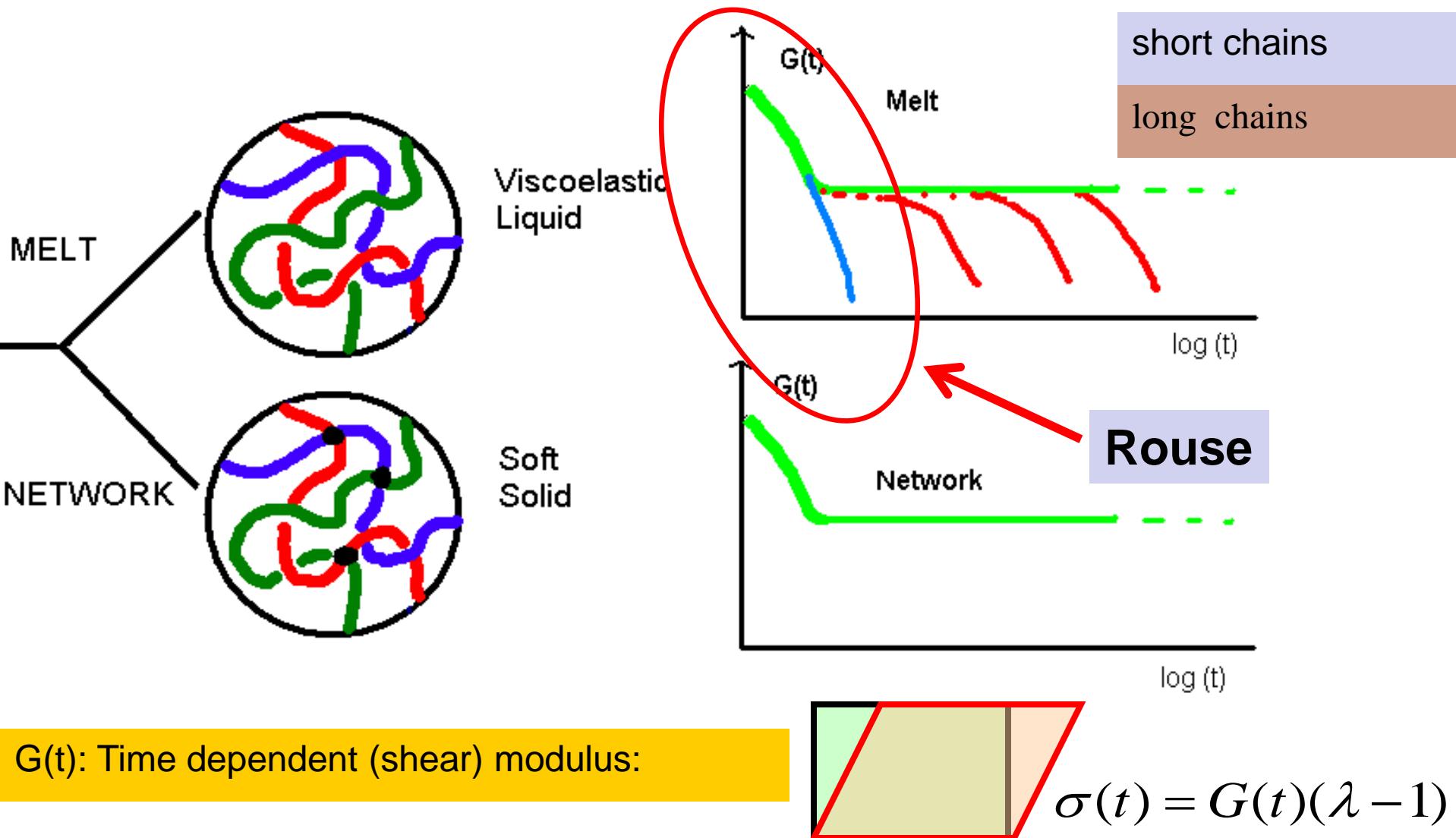
“Rouse Model”₁₉₅₃

Brownian dynamics of
chain of harmonic springs:

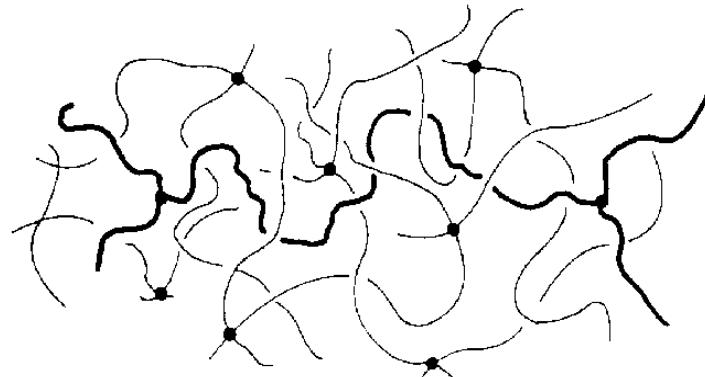


$$g_1(t) = \langle (r_i(t) - r_i(0))^2 \rangle \propto \begin{cases} t, & t < \tau_0 \\ t^{1/2}, & \tau_0 < t < \tau_R \\ t, & t > \tau_R \end{cases}$$

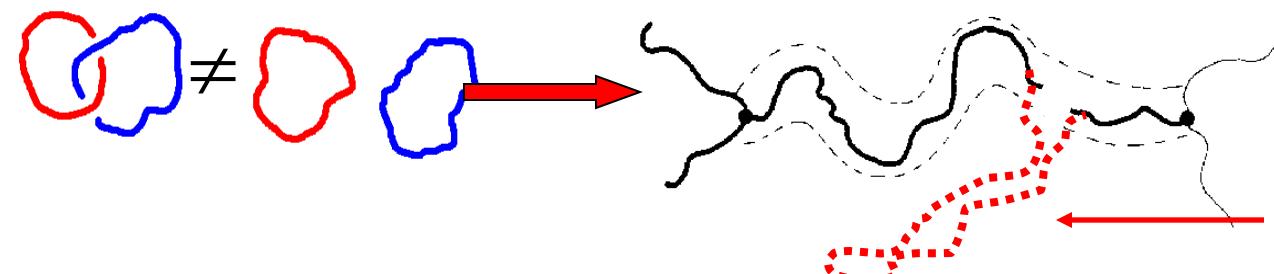
MELT \leftrightarrow NETWORK



Tube Model



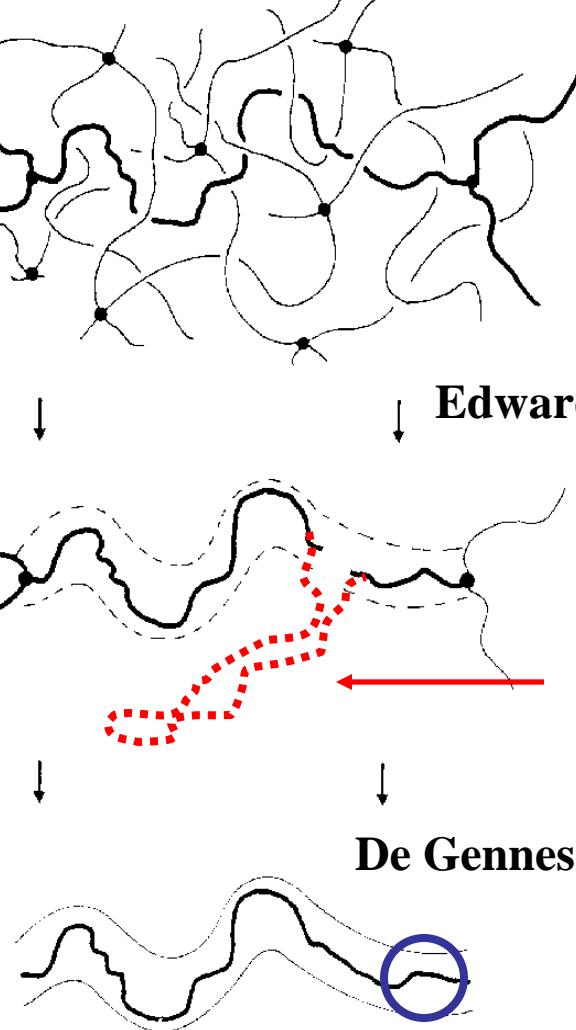
↓ Edwards '67



“Tube leakage”:
Entropy penalty for
defects too large

$$d_T \propto N_e^{1/2}$$
$$G^0 \propto N_e^{-1}$$

Unique length d_T



↓ De Gennes '71



Rouse model subject to tube constraints

Polymer Melts

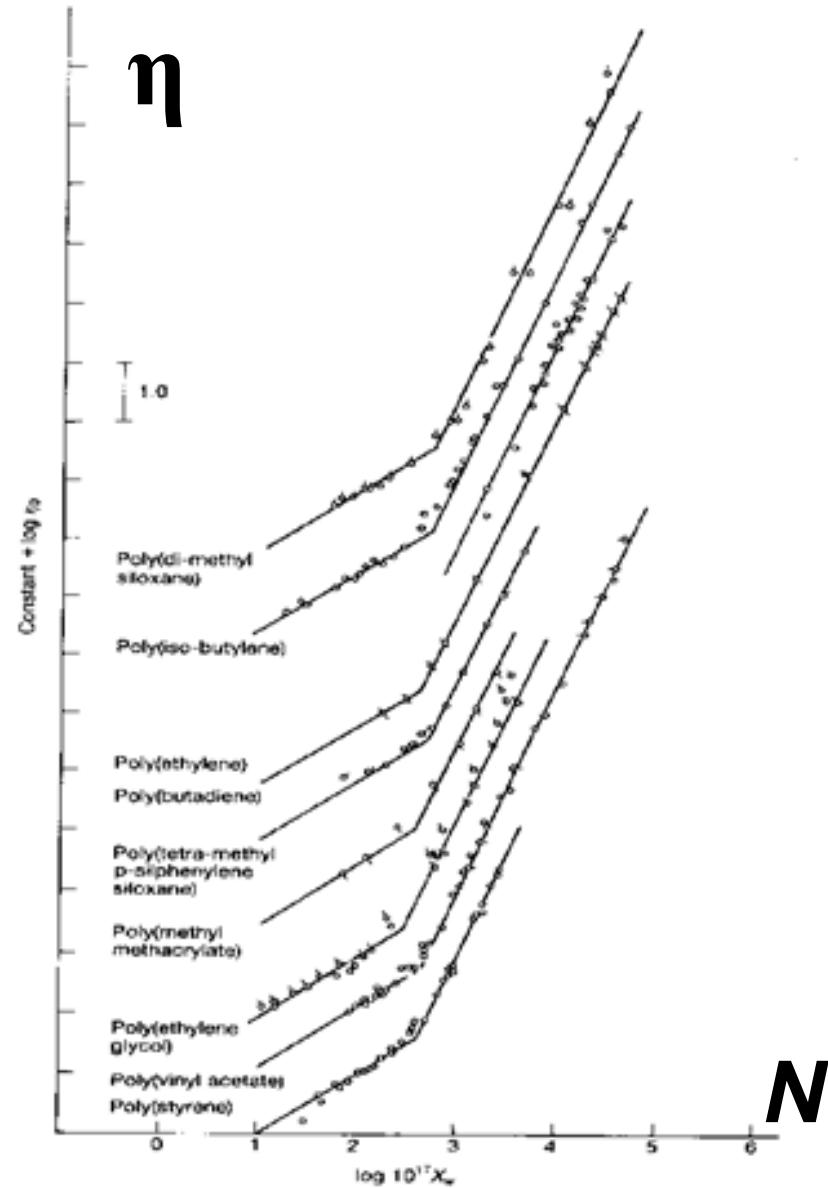
Short Chains $N < N_e$

$N_e < N$ Long Chains

• $D \propto N^{-1}$	Diffusion	$D \propto N^{-2}$
• $\tau_R \propto N^2$	Relaxation Time	$\tau_d \propto N^{3\dots}$
• $\eta \propto N^1$	Viscosity	$\eta \propto N^{3.4}$
• $G^0 = 0$	Modulus	$G^0 \sim \frac{kT}{N_e} \propto N^0$
•		$(t \ll \tau_d)$
• Rouse Model		Reptation Model
• Chain motion subject		Chain motion subject to
• background and heat bath		TUBE CONCEPT

UNIVERSAL PROPERTIES

Melt viscosity vs chain length



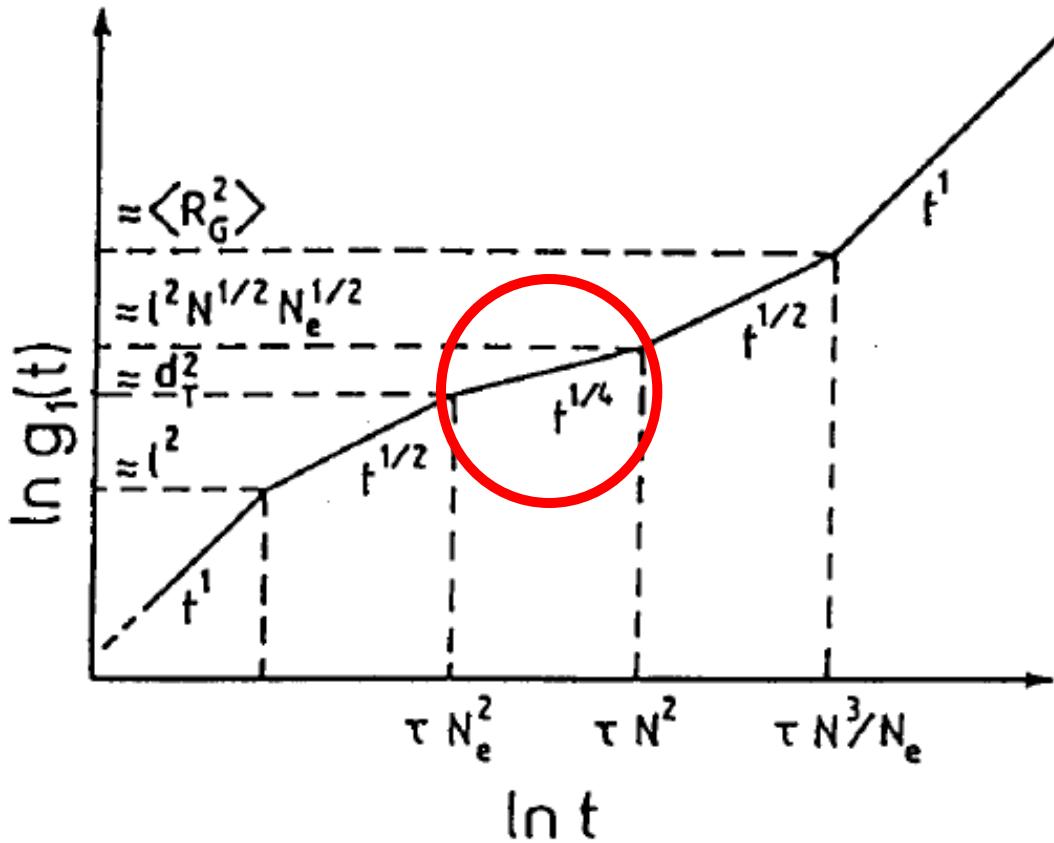
From Doi and Edwards



Mean square displacements: reptation

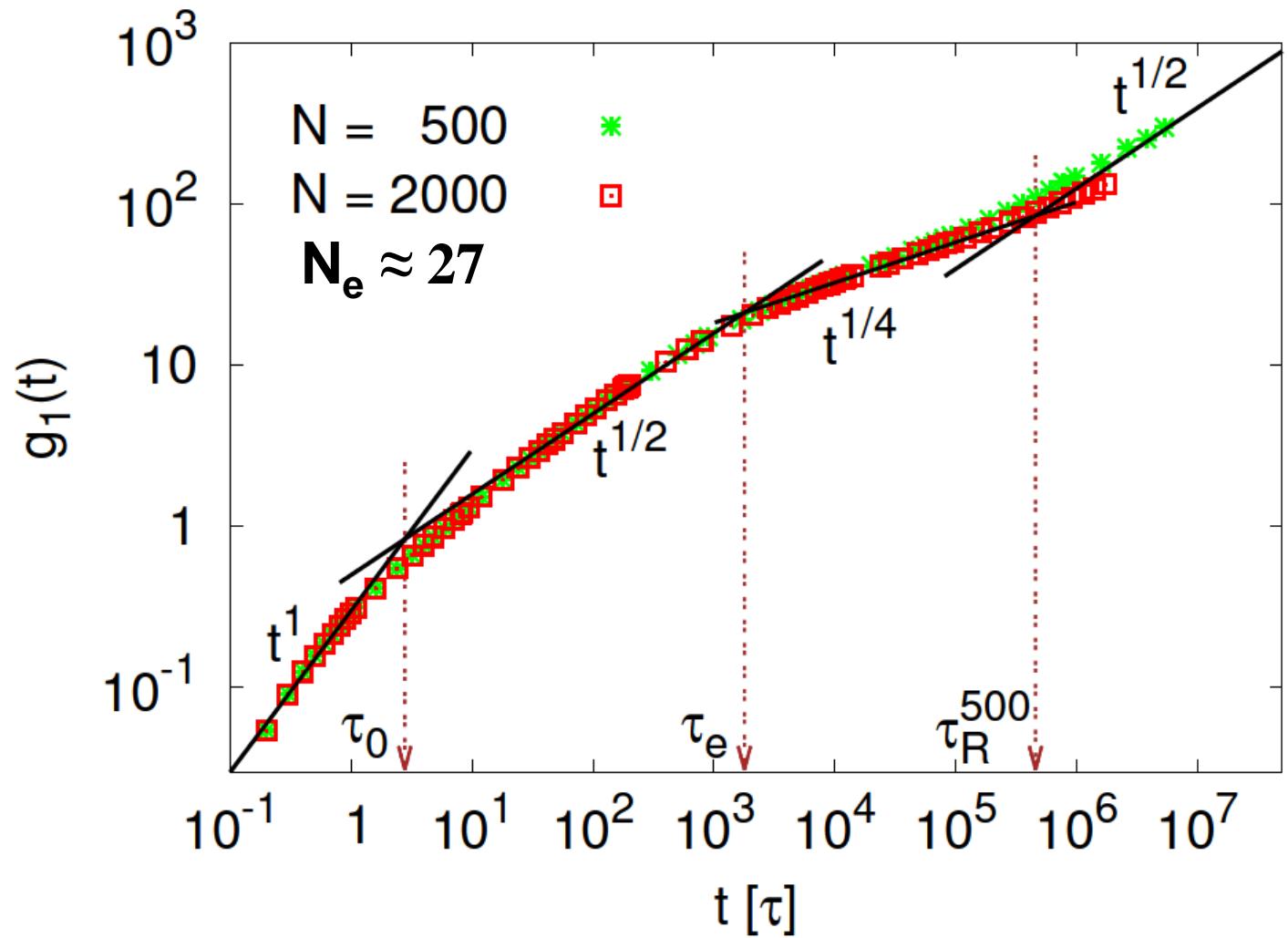
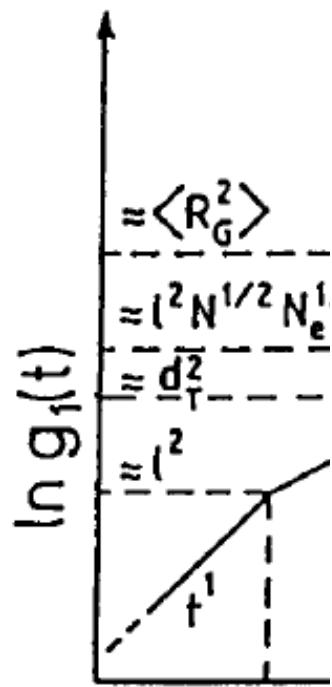
$$g_1(t) = \langle (r_i(0) - r_i(t))^2 \rangle \propto \begin{cases} t, & t < \tau_0 \\ t^{1/2}, & \tau_0 < t < \tau_e \approx N_e^2 \\ t^{1/4}, & \tau_e < t < \tau_R \approx N^2 \\ t^{1/2}, & \tau_R < t < \tau_d \approx N^3 / N_e \\ t, & t > \tau_d \end{cases}$$

$$g_3(t) = \langle (r_{cm}(t) - r_{cm}(0))^2 \rangle \propto \begin{cases} t, & t < \tau_e \\ t^{1/2}, & \tau_e < t < \tau_R \\ t, & t > \tau_R \end{cases}$$



Famous $t^{1/4}$ power law for mean square displacement





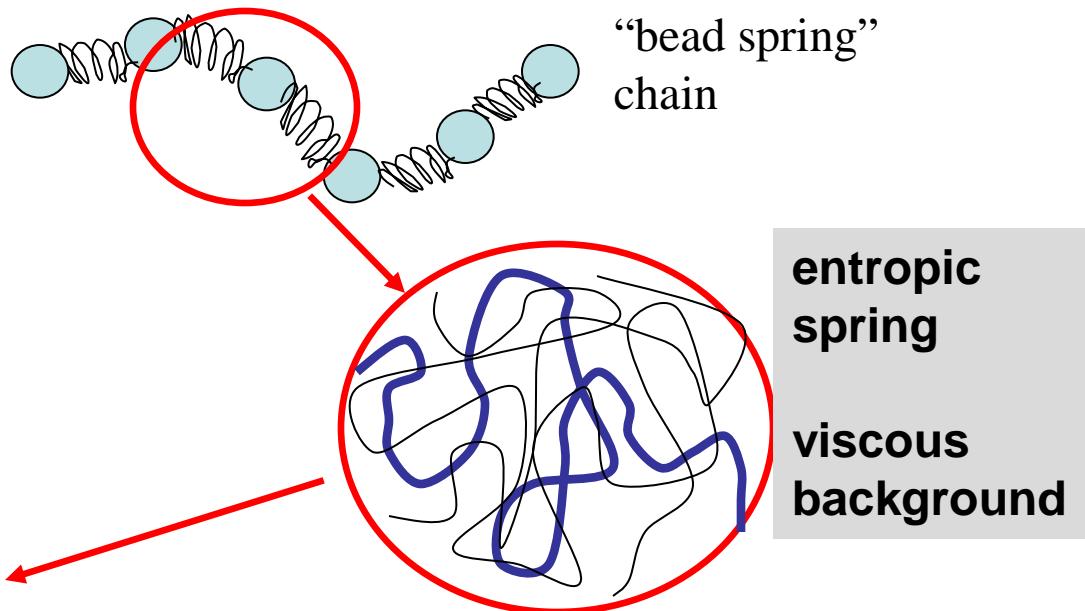
H. P. Hsu



End Detour: Rouse Reptation – Most Basic Aspects



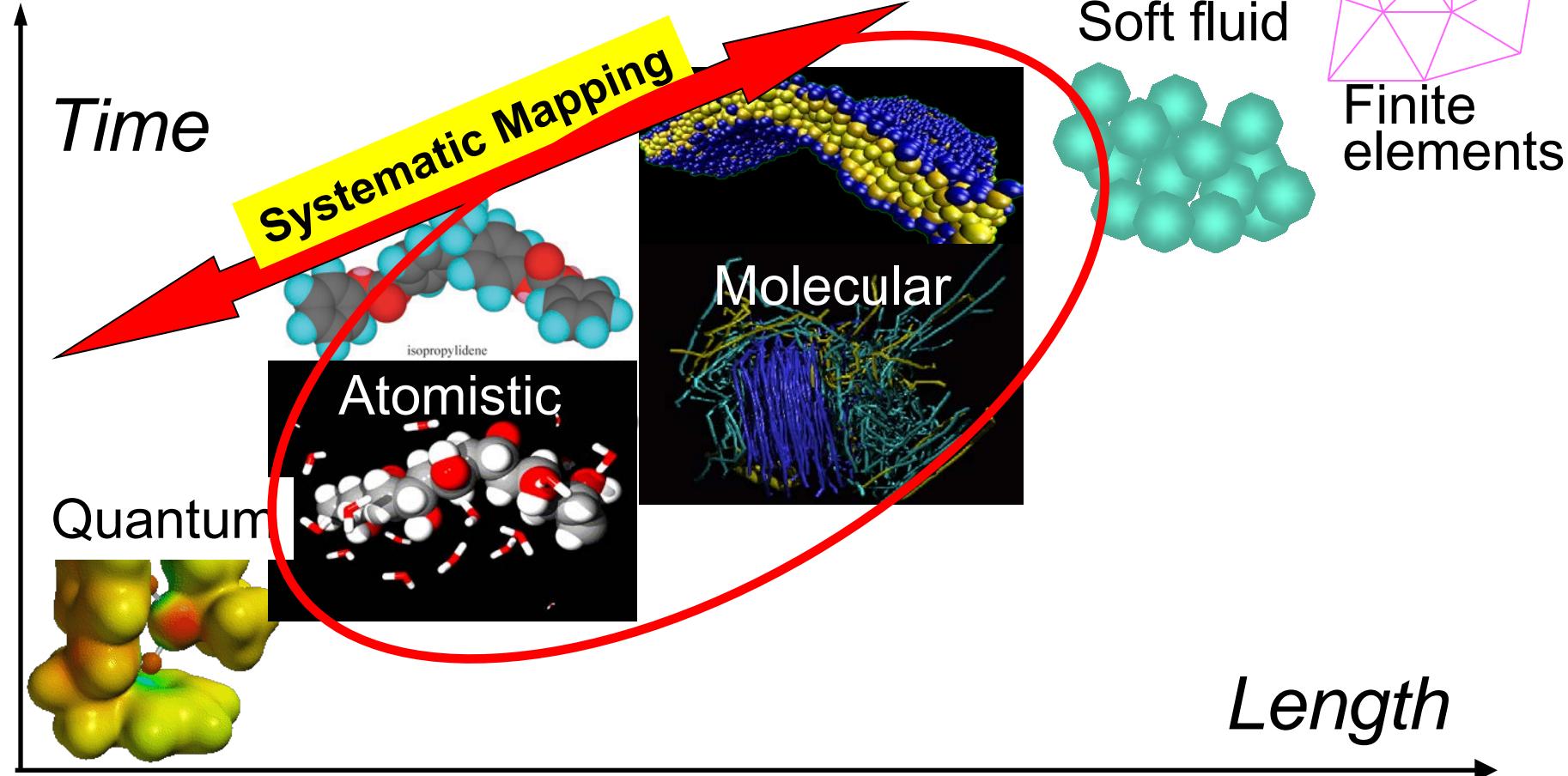
“Rouse Model”¹⁹⁵³



“Reptation/Tube Model”^{1967,71}



Coarse Graining



Local Chemical Properties --- Scaling Behavior of Nanostructures
Energy Dominance --- Entropy Dominance of Properties



Coarse Graining (CG)

Micro-Meso Simulation

Interplay Energy \leftrightarrow Entropy
Free Energy Scale: $k_B T$

atomistic model \Rightarrow *parameterize cg model*

Different general strategies:

Structure based

Force based

Potential based

Methods to calculate CG interactions:

(Iterative) Boltzmann Inversion

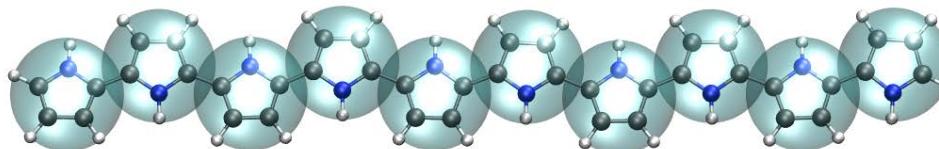
Inverse MC,

Force Matching

Relative Entropy

Coarse-graining of atomistic systems

Atomistic system
(polypyrrole chain)



Coarse-grained system
(10 PPY repeat units)

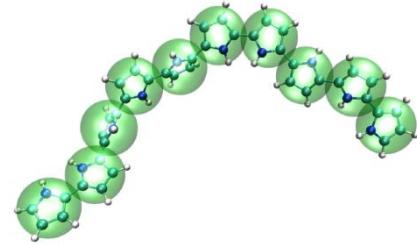
$$\mathbf{r}^n = \{\mathbf{r}_1, \dots, \mathbf{r}_n\}$$

$$\mathbf{p}^n = \{\mathbf{p}_1, \dots, \mathbf{p}_n\}$$

$$h(\mathbf{r}^n, \mathbf{p}^n) = \sum_{i=1}^n \frac{1}{2m_i} \mathbf{p}_i^2 + u(\mathbf{r}^n)$$

$$\mathbf{R}^N = \{\mathbf{R}_1, \dots, \mathbf{R}_N\}$$

$$\mathbf{P}^N = \{\mathbf{P}_1, \dots, \mathbf{P}_N\}$$



$$H(\mathbf{R}^N, \mathbf{P}^N) = \sum_{I=1}^N \frac{1}{2M_I} \mathbf{P}_I^2 + U(\mathbf{R}^N)$$

canonical ensemble (NVT)

$$p_{rp}(\mathbf{r}^n, \mathbf{p}^n) = p_r(\mathbf{r}^n)p_p(\mathbf{p}^n)$$

$$p(\mathbf{r}^n) \sim \exp[-\beta u(\mathbf{r}^n)]$$

$$p(\mathbf{p}^n) \sim \exp\left[-\beta \sum_{i=1}^n \frac{\mathbf{p}_i^2}{2m_i}\right]$$

$$P_{RP}(\mathbf{R}_N, \mathbf{P}_N) = P_R(\mathbf{R}_N)P_P(\mathbf{P}_N)$$

$$P(\mathbf{R}_N) \sim \exp[-\beta U(\mathbf{R}^N)]$$

$$P(\mathbf{P}_N) \sim \exp\left[-\beta \sum_{I=1}^N \frac{\mathbf{P}_I^2}{2M_I}\right]$$

Basic standard methods

- **Boltzmann inversion (structure based cg)**

The idea of Boltzmann inversion stems from the fact that in a canonical ensemble *independent* degrees of freedom q obey the Boltzmann distribution, i. e.

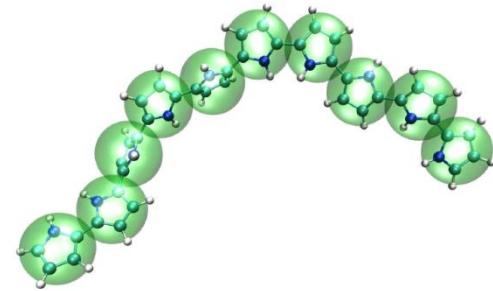
$$P(q) = Z^{-1} \exp [-\beta U(q)] , \quad (2.19)$$

where $Z = \int \exp [-\beta U(q)] dq$ is a partition function, $\beta = 1/k_B T$. Once $P(q)$ is known one can invert eq. 2.19 and obtain the coarse-grained potential, which in this case is a potential of mean force

$$U(q) = -k_B T \ln P(q) . \quad (2.20)$$

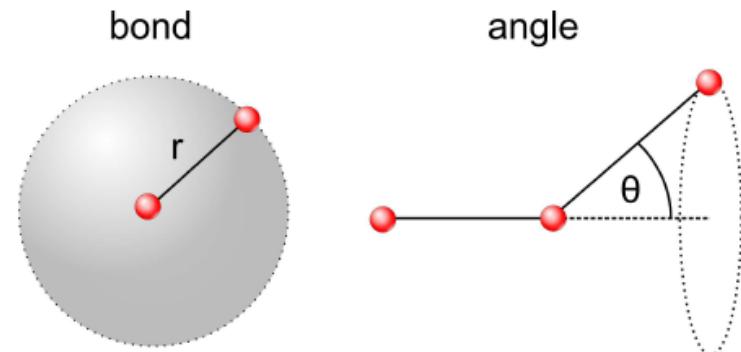
Apply this concept to $g(r)$, the radial distribution function

Basic standard methods



- Boltzmann inversion

$$\begin{aligned} P(r, \theta, \varphi) &= \exp [-\beta U(r, \theta, \varphi)] \\ P(r, \theta, \varphi) &= P_r(r) P_\theta(\theta) P_\varphi(\varphi) . \end{aligned}$$



$$U(r, \theta, \varphi) = U_r(r) + U_\theta(\theta) + U_\varphi(\varphi) ,$$

$$U_q(q) = -k_B T \ln P_q(q), \quad q = r, \theta, \varphi$$

Hendersen Theorem: unique correspondence between $U(r)$ and $g(r)$

Potential is state point dependent

Basic standard methods

- Iterative Boltzmann inversion

$$U^{(1)}(r) = -k_{\text{B}}T \ln g(r)$$

$$U^{(n+1)} = U^{(n)} + \Delta U^{(n)}$$

$$\Delta U^{(n)} = k_{\text{B}}T \ln \frac{P^{(n)}}{P_{\text{ref}}}$$

Iterative Boltzmann Inversion

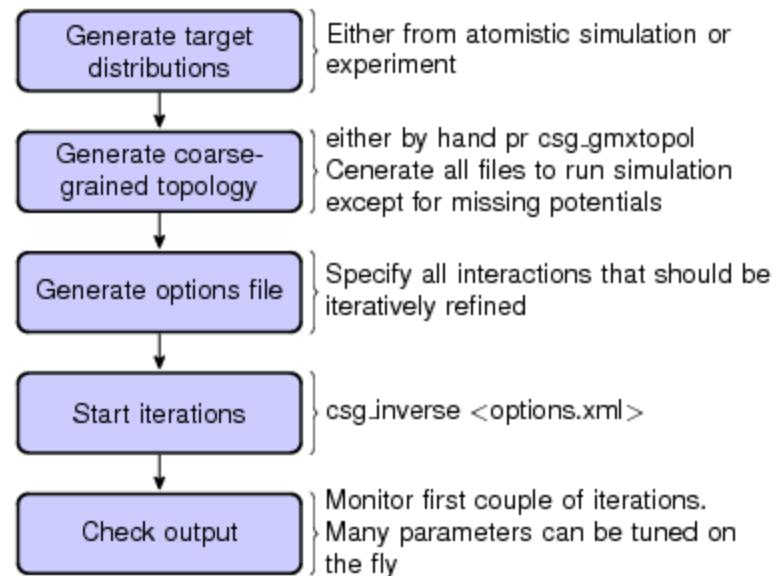
S_α are proportional to the radial distribution function

$$S_\alpha = \frac{N(N-1)}{2} \frac{4\pi r_\alpha^2 \Delta r}{V} g(r_\alpha)$$

$$A_{\alpha\gamma} = \beta (\langle S_\alpha \rangle \langle S_\gamma \rangle - \langle S_\alpha S_\gamma \rangle)$$

Ignore off-diagonal elements, and take into account that for an ideal gas

$$\Delta U(r) = k_B T \ln \frac{g(r)}{g^{ref}(r)}$$



IBI: local, no correlations. Robust, but it is not clear how well it performs for mixtures.

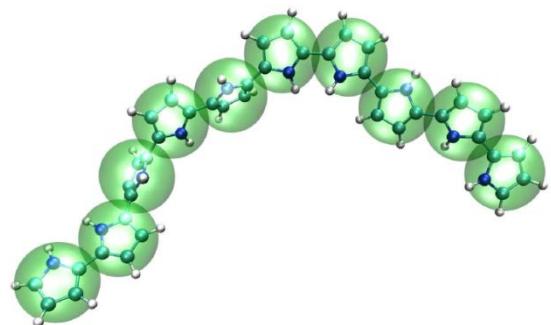
Basic standard methods

- **Inverse Monte Carlo (Lyubartsev)**

Similar concept as Boltzmann inversion, however sampling directly on the discretized potential $u(r, \dots)$

Automatically includes correlations (unlike factorization ansatz for Boltzmann inversion)

Numerically more demanding



Basic standard methods

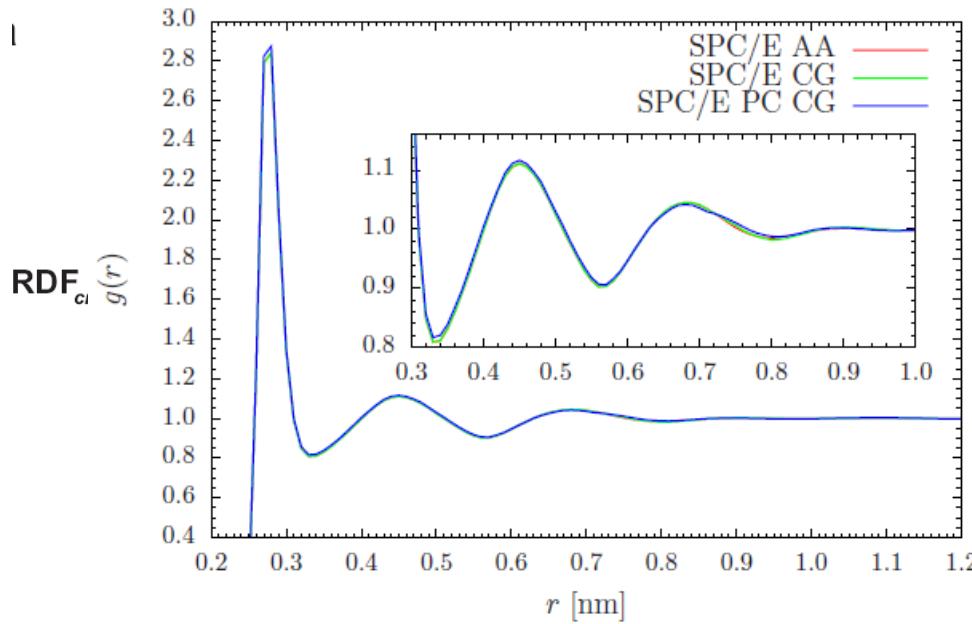
- Iterative Boltzmann inversion
- Inverse Monte Carlo (Lyubartsev)

Both structure based methods

Henderson Theorem: Unique relation between $g(r)$ and $U(r)$,
but potentials are **state point dependent**, not generally transferable

only $g(r)$ is reproduced and thus the compressibility,
but not the pressure (i.e. thermodynamic properties)!

Structure based CG - Compressibility



SPC/E water,
cg based on matching $g(r)$

Pressure

$$P = \rho k T - \frac{2}{3} \pi \rho^2 \int_0^\infty U'(r) g(r) r^3 dr$$

Compressibility

$$\kappa_T = \frac{1}{\rho k_B T} (1 + 4\pi\rho \int_0^\infty r^2 [g(r) - 1] dr) = \frac{1}{\rho k_B T} \frac{\langle N^2 \rangle - \langle N \rangle^2}{\langle N \rangle}$$

Basic standard methods

- Iterative Boltzmann inversion
- Inverse Monte Carlo (Lyubartsev)

Structure based
Coarse Graining:

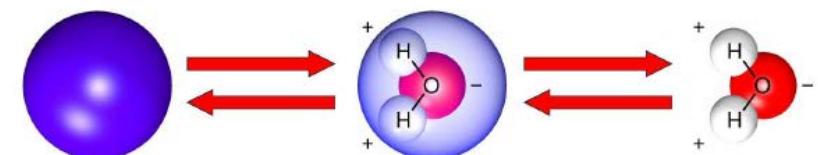
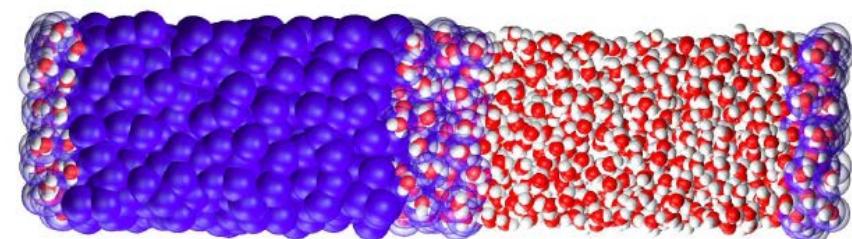
Perfect match of $g(r)$
all atom – coarse grained

$$K_{\text{atomistic}} = K_{\text{cg}}$$

BUT

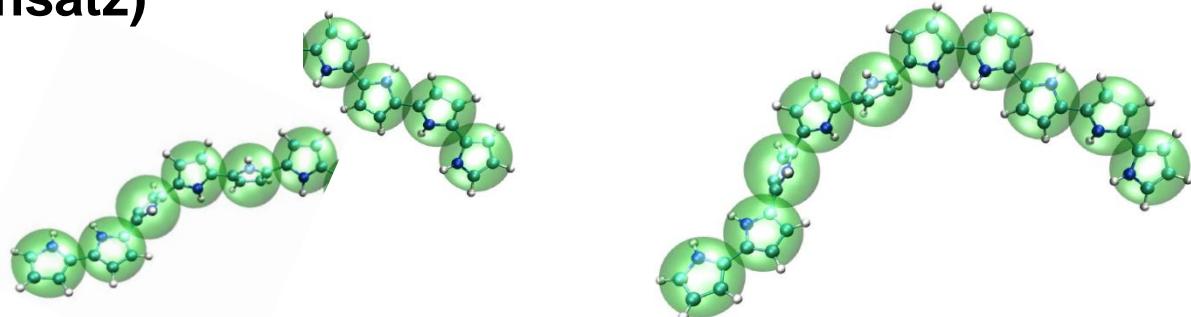
$$p_{\text{cg}} = 6200 p_{\text{atomistic}}$$

Coarse Grained Water: SPC/E



Basic standard methods: Force Matching

- Ercolessi&Adams 1993, Noid&Voth et al JCP 2008ff (Martini Force Field follows similar ansatz)



$$\vec{f}_i^{\text{ref}} = M_i \sum_{\alpha} \frac{w_{\alpha} \vec{f}_{\alpha}}{m_{\alpha}}$$

cg force

AA forces

$$\vec{R}_i = \sum_{\alpha} w_{\alpha} \vec{r}_{\alpha}$$

Position of cg particle

$$M_i = (\sum_{\alpha} w_{\alpha}^2 / m_{\alpha})^{-1}$$

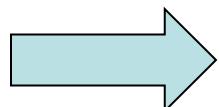
mass of cg particle

VOTCA: Versatile Objects for Coarse-graining

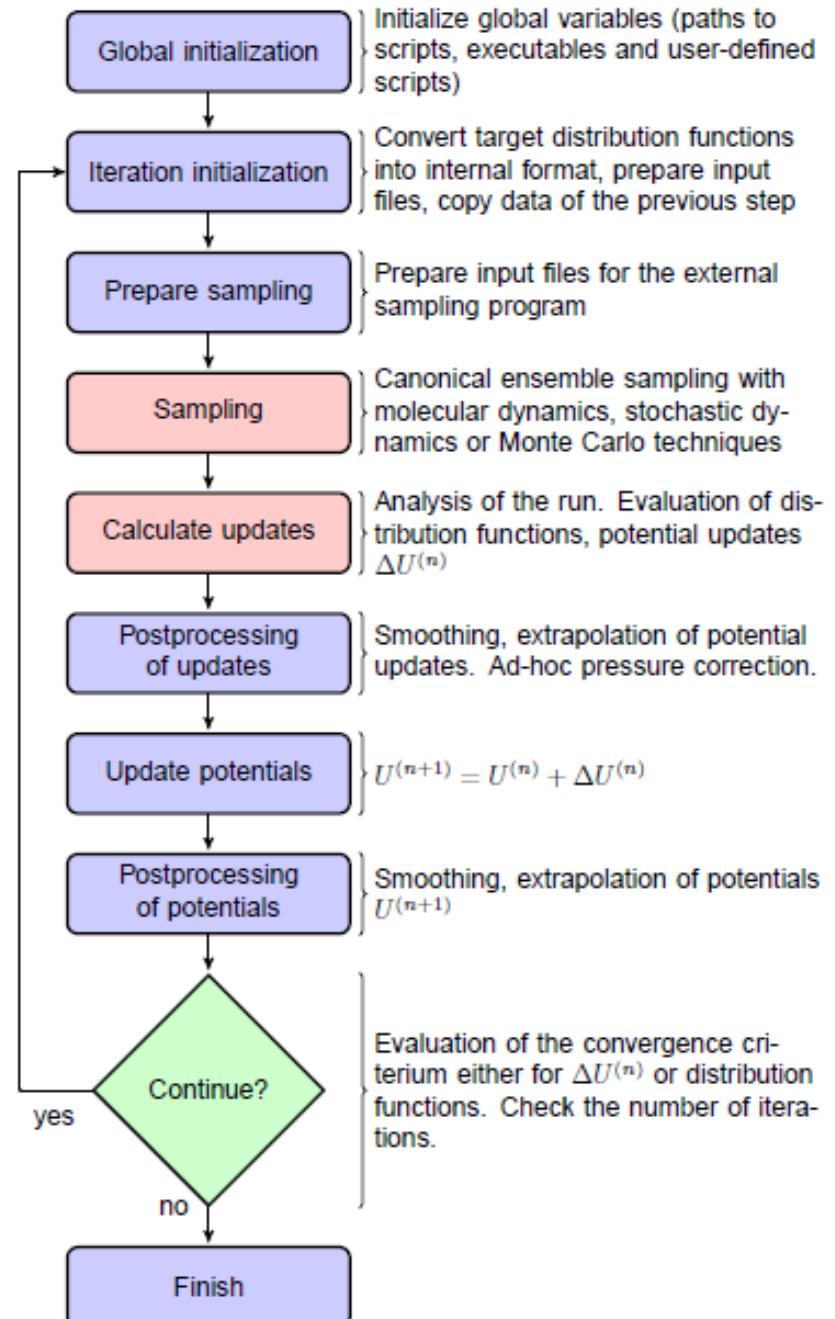
- General ansatz:

- n “atomistic”,
- N cg particles

$$\begin{aligned} \{r_i\}, i = 1, 2, \dots, n \\ \{R_j\}, j = 1, 2, \dots, N \end{aligned}$$



$$R = \hat{M}r$$





Structure Based Coarse Graining

Coarse Graining (CG) <=> Inverse Mapping

Micro-Meso Simulation

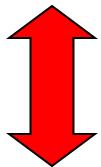
Interplay Energy \leftrightarrow Entropy
Free Energy Scale: $k_B T$

atomistic model

=> parameterize cg model

=> run/analyze large system (melt etc)

“measurements” $R^2(N)$...



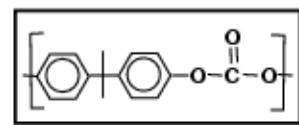
reintroduce details

=> run/analyze atomistic system

Structure Based Molecular Coarse-Graining: Bisphenol-A-Polycarbonate

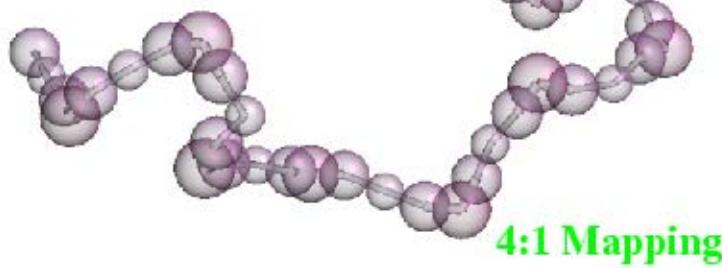
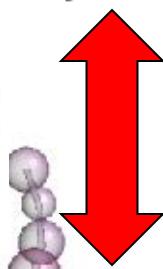
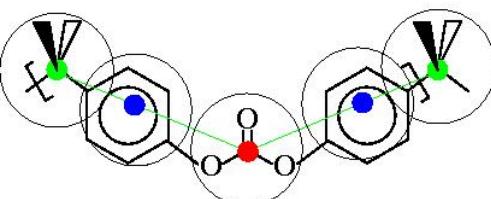
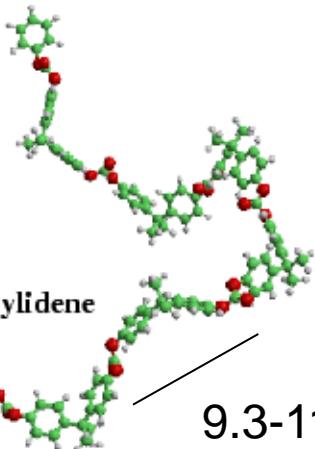


BPA-PC, N=10



phenylene
isopropylidene
carbonate

9.3-11.5 Å



4:1 Mapping

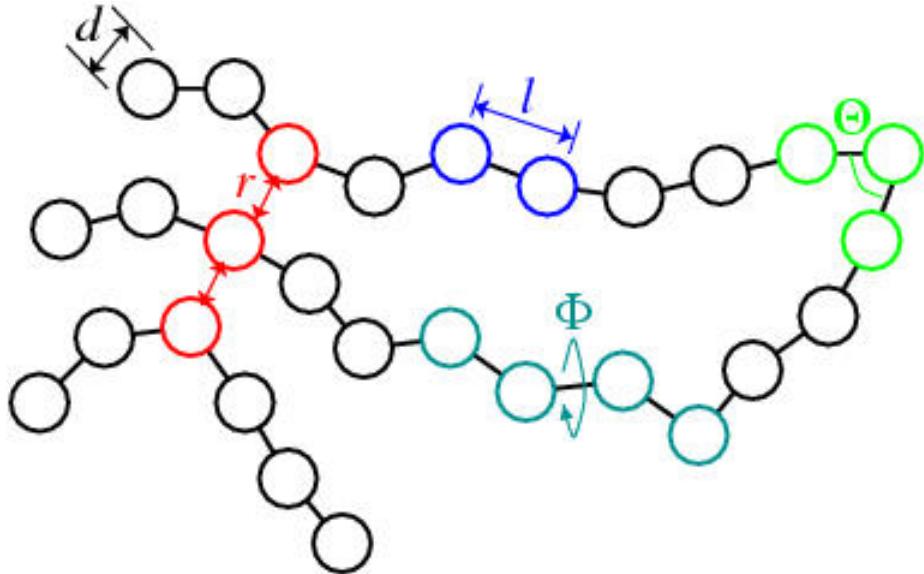
Coarse-graining:

map **bead-spring chain** over molecular structure.
=> Many fewer degrees of freedom

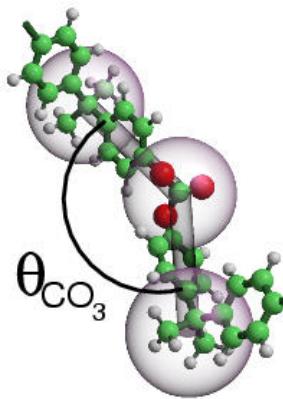
Inverse mapping: grow atomic structure on top of coarse-grained backbone

=> Large length-scale equilibration in an atomically resolved polymer

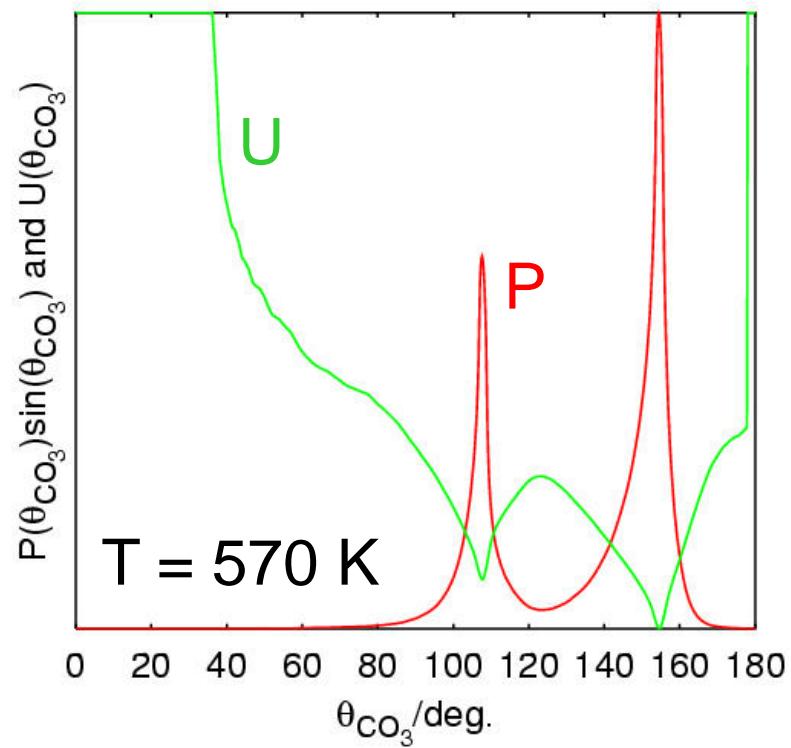
Interaction Energies in the Coarse-Grained Model



- Excluded volume
- Bonds
- Angles
- Torsions



Angle potentials are T-dependent Boltzmann inversions; e.g., at carbonate:



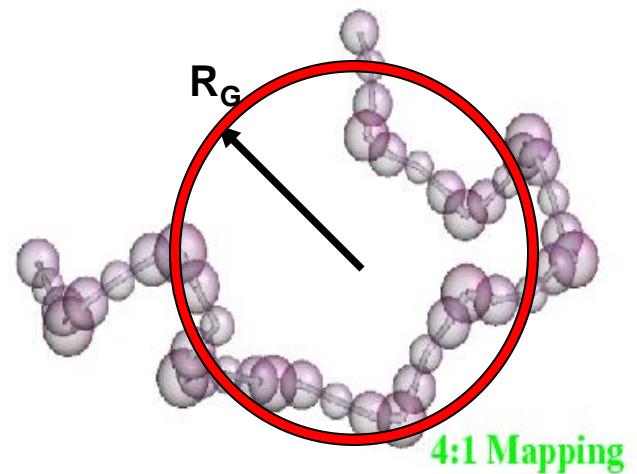


Results for Melts, N=20....120

- Molecular Coarse-Grained Melt

$$\left\langle R_G^2(N) \right\rangle / N \simeq 37 \text{ \AA}^2$$

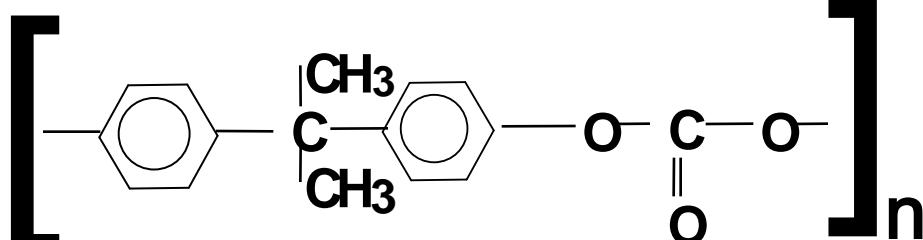
R_G of coarse grained simulations
agrees to n-scattering experiments!



- Reintroduce Details: Inverse Mapping

Speedup $\approx 10^4$!

Polycarbonate



BPA-PC

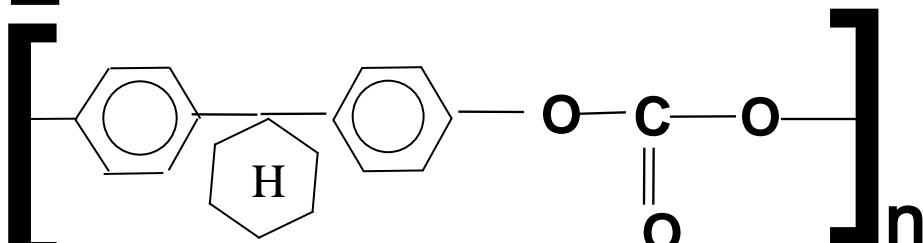
$$T_{VF} = 390 \text{ K}$$

$$T_G = 420 \text{ K}$$

ductile

$$N_e = 7$$

$$M_e = 1700$$



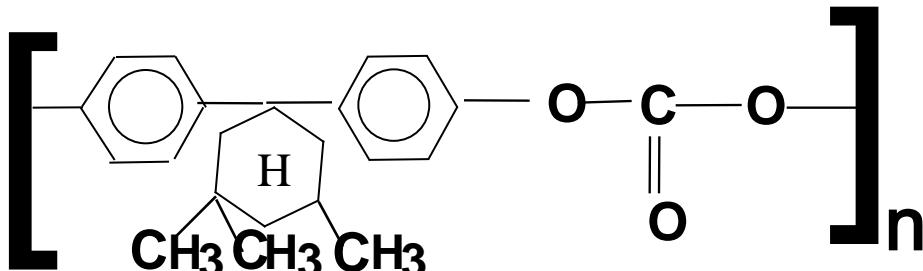
BPZ-PC

$$T_{VF} = 390 \text{ K}$$

$$T_G = 420 \text{ K}$$

$$N_e = 10$$

brittle



TMC-PC

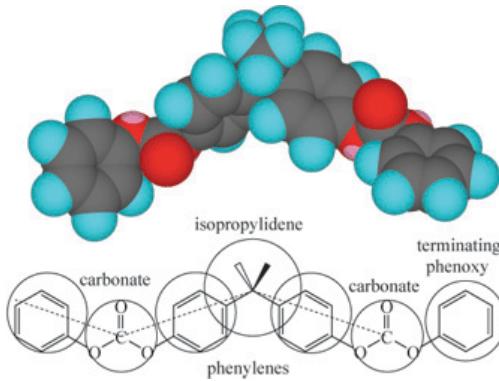
$$T_{VF} = 480 \text{ K}$$

$$T_G = 510 \text{ K}$$

$$N_e = 15$$

“somewhat ductile”

(M_e Entanglement molecular weight $M < M_e$, $T_{VF}(M) \approx T_{VF} - 50 \text{ K}$)



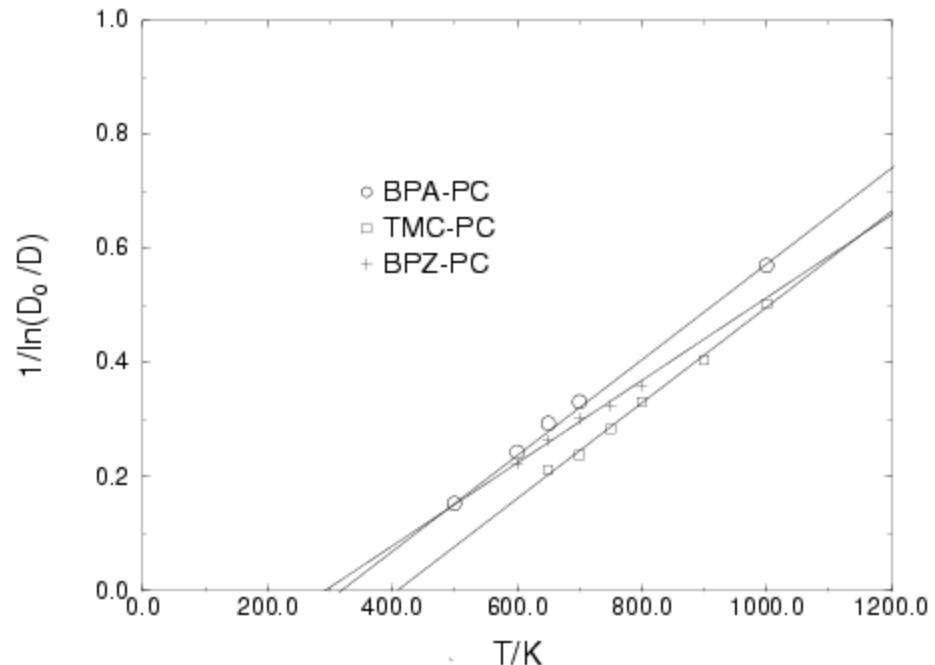
Dynamics: Vogel-Fulcher Law

Simulation vs Experiment

BPA	BPZ	TMC
323K	292K	407K
387K	392K	477K

Simulation: N=10
Experiment: N≈200

$$D = D_o \exp\left(-\frac{A_o}{T - T_{VF}}\right)$$



[W. Tschöp, K. Kremer, J. Batoulis, T. Bürger, O. Hahn, Acta Polym. **49**, 61 (1998); ibid. **49**, 75].

Viscosity => Time Mapping

$$\eta = \rho R_N^2 k_B T / s^3 (N - 1)^2 36 D$$

- Melt simulation
- Viscosity from chain diffusion coefficient (Rouse)
- Property of entire chains

$$\tau \approx 22 \text{ ps}$$

(T=500K, new data 2005)

[W. Tschöp, K. Kremer, J. Batoulis, T. Bürger, O. Hahn, Acta Polym. **49**, 61 (1998); ibid. **49**, 75].

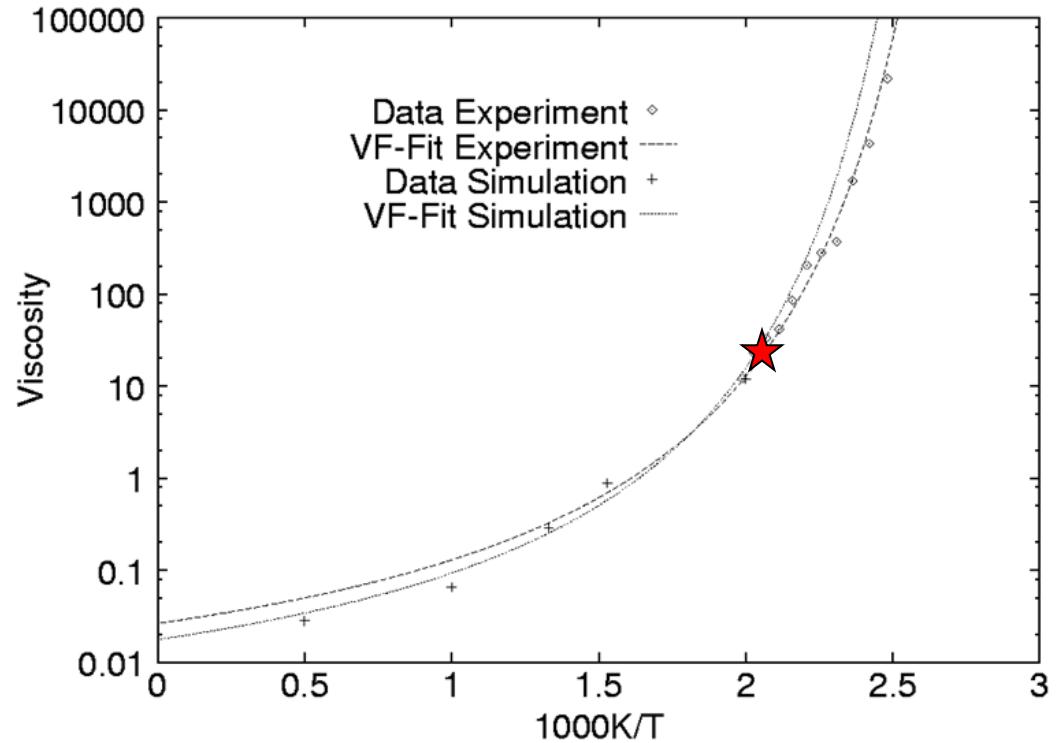
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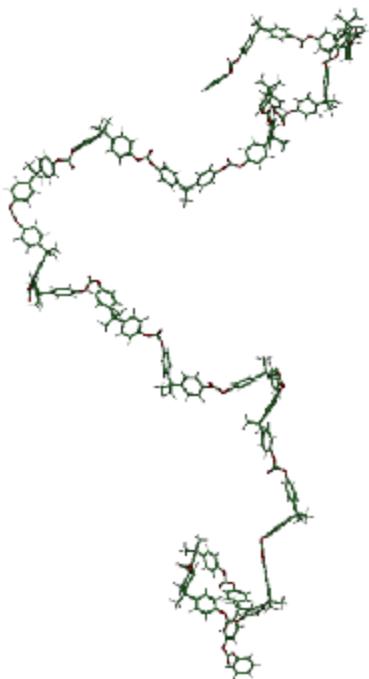


[W. Tschoep, K. Kremer, J. Batoulis, T. Burger, O. Hahn, Acta Polym. **49**, 61 (1998); ibid. **49**, 75].

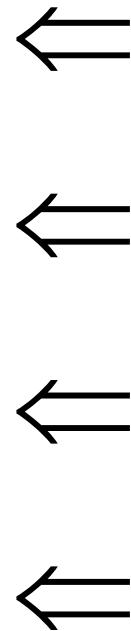


How good are generated conformations?

Inverse Mapping: Reintroduce Chemical Details



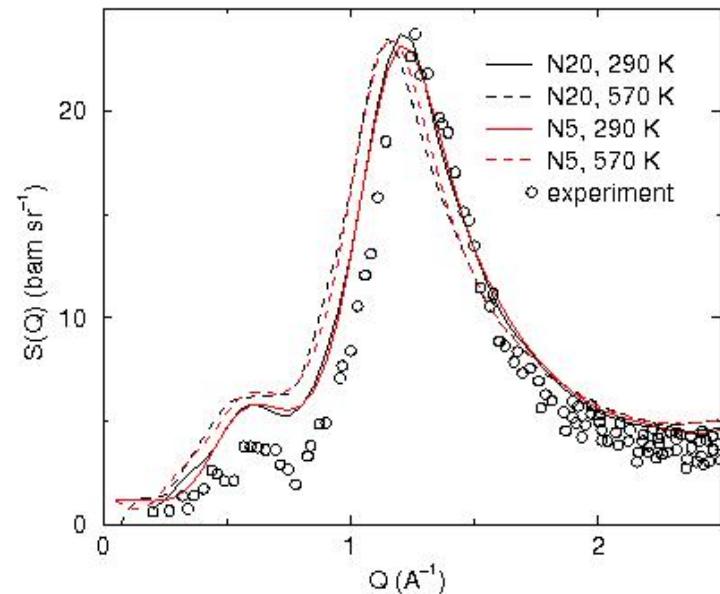
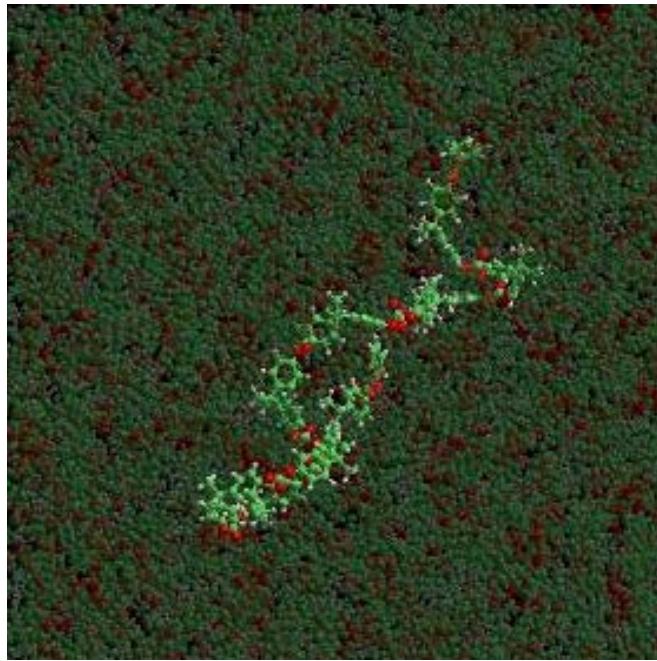
All atom model



Coarse grained
BPA-PC chain



Comparison: Simulation vs n-Scattering



$$\left\langle R_G^2(N) \right\rangle / N \cong 37 \text{ \AA}^2$$

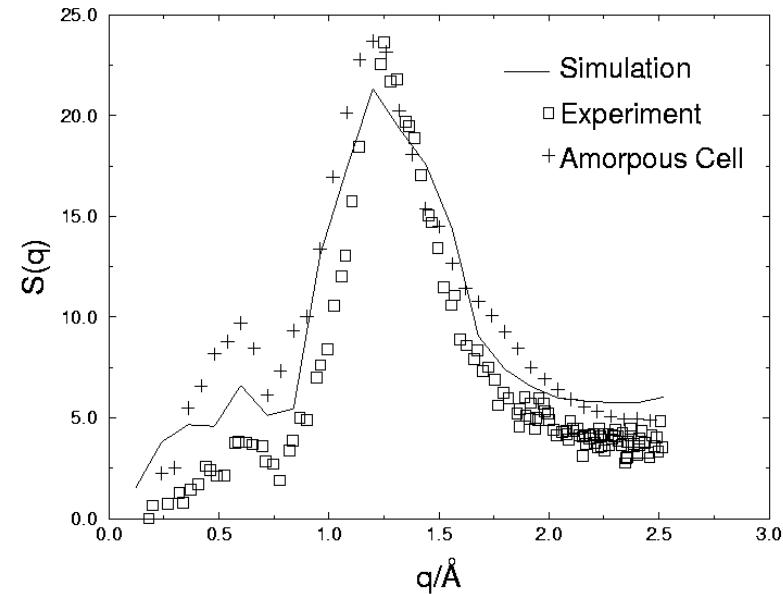
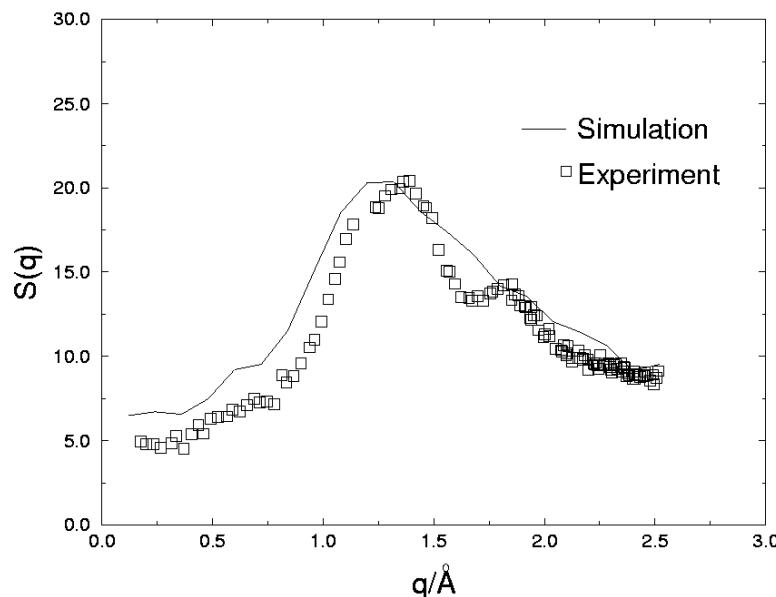
J. Eilhard et al, J. Chem. Phys.
110, 1819 (1999)
B. Hess et al, Soft Matter 2006

Well equilibrated coarse grained AND all atom melt configurations!

Comparison: Simulation n-Scattering

Structure factors of
(deuterated) BPA-PC

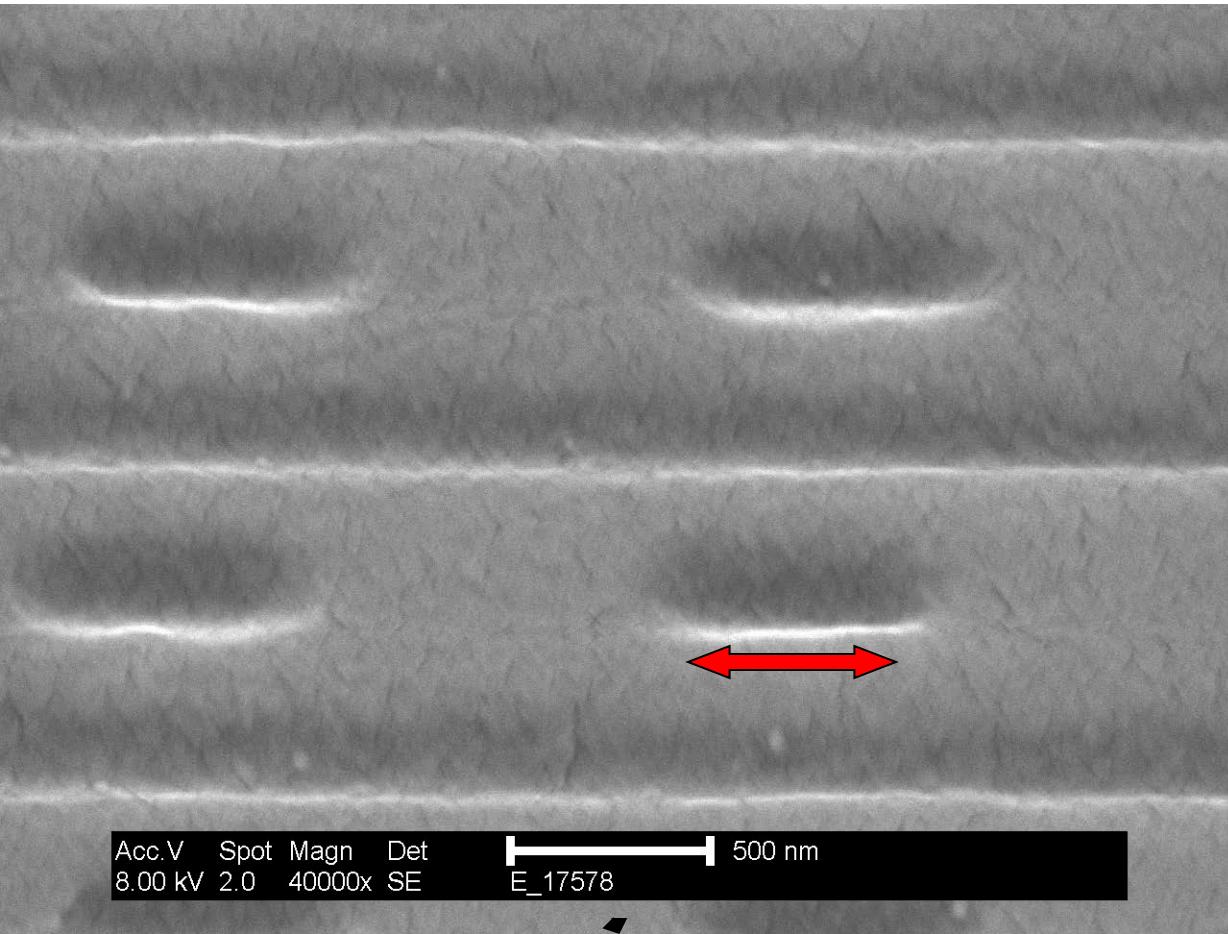
Right: standard BPA-PC
Bottom: fully deuterated BPA-PC



[J. Eilhard, A. Zirkel, W. Tschöp,
O. Hahn, K. K., O. Schärpf,
D. Richter, U. Buchenau,
J. Chem. Phys. **110**, 1819 (1999)]

Polycarbonate (PC) and the PC/Ni interface

Grooves and address pits of a die cast sample of polycarbonate for a high storage density (**blue laser**) optical disc



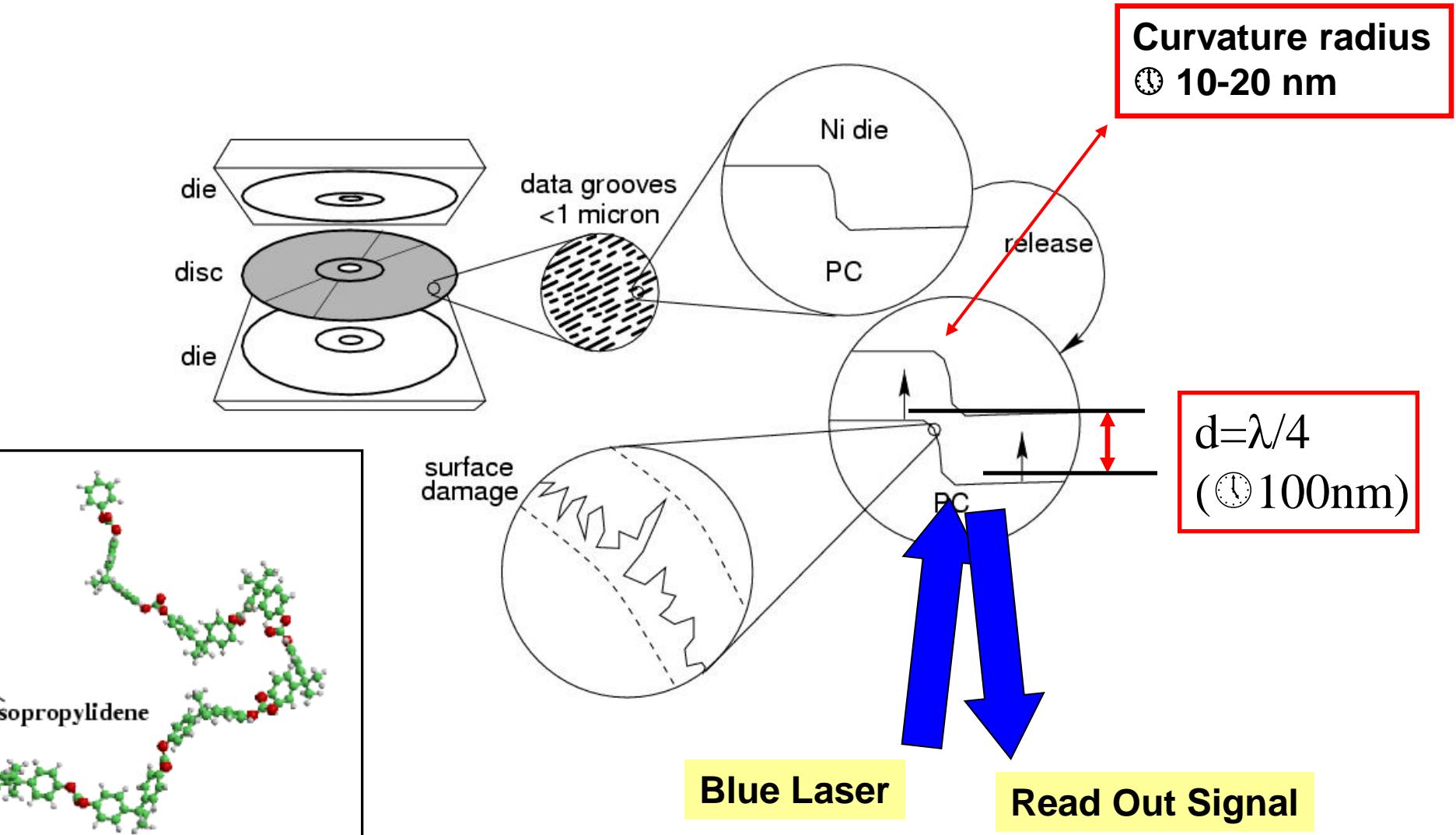
500 nm

Acc.V Spot Magn Det
8.00 kV 2.0 40000x SE

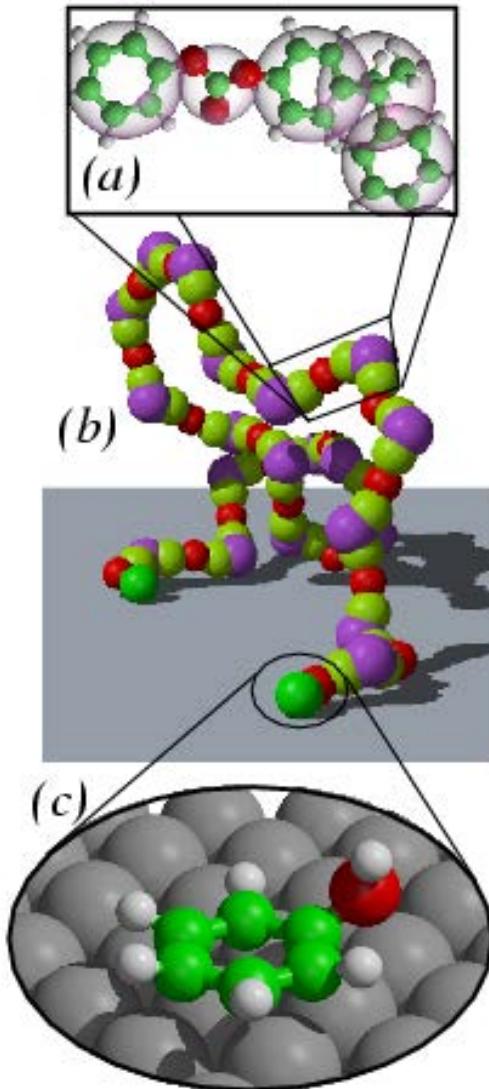
500 nm
E_17578

Bayer Materials

Why study Polycarbonate and the PC/Ni interface?



Simulating BPA-PC/Metal Interfaces

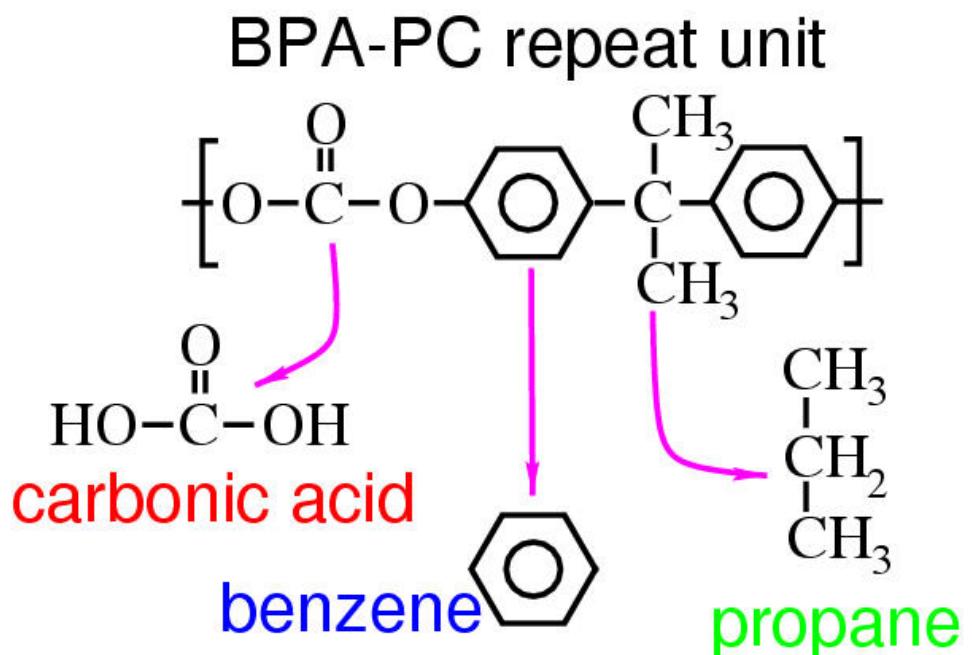
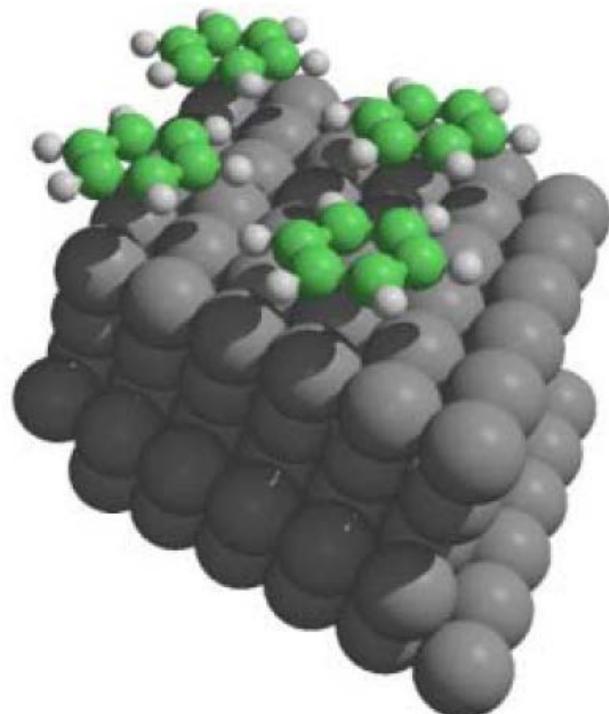


Molecular structure coarse-grained onto bead-spring chain

Simulation of coarse-grained BPA-PC liquids ($T = 570\text{K}$) next to metal surface

Specific surface interactions investigated via *ab initio* calculations

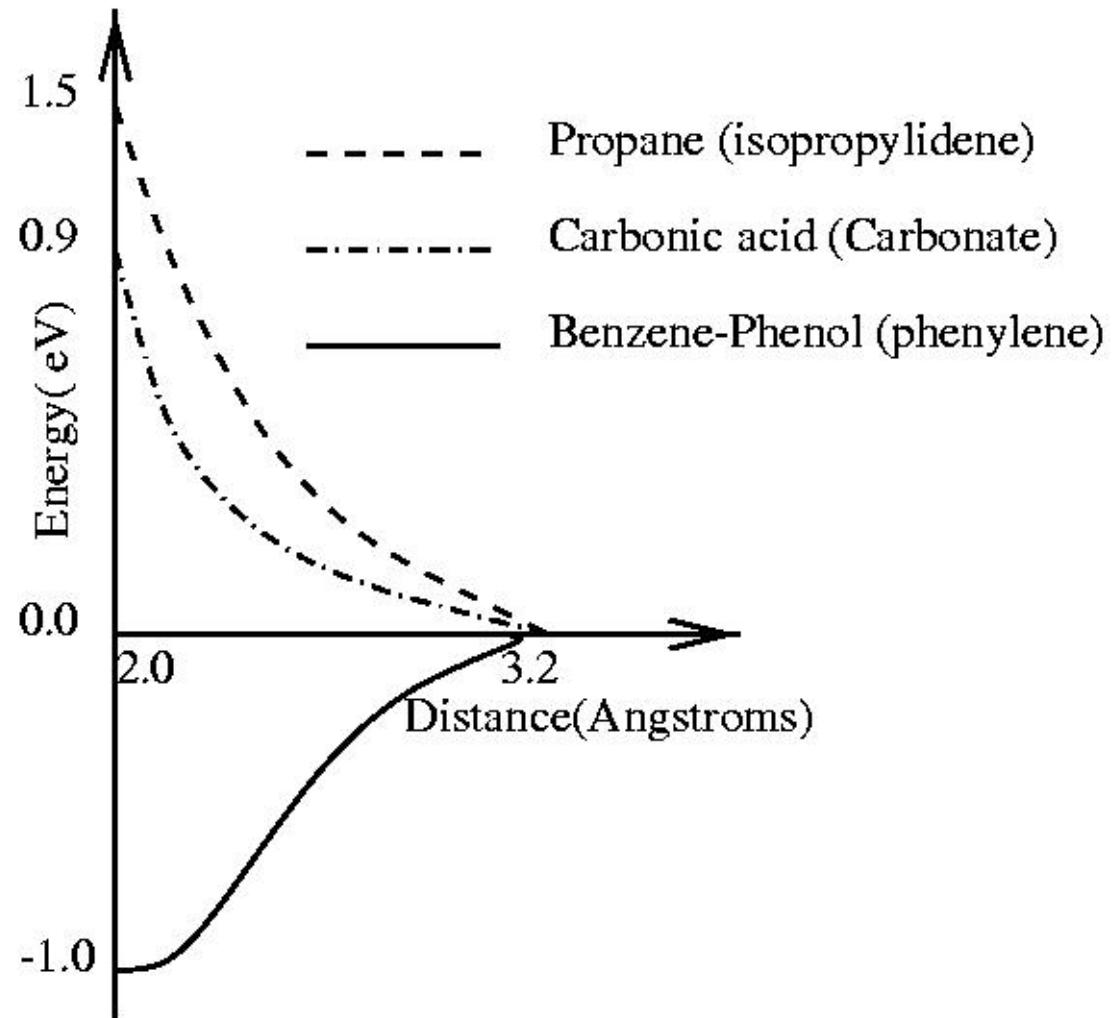
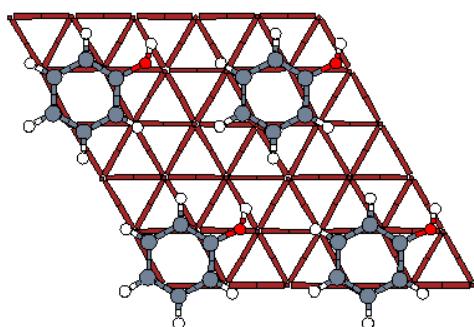
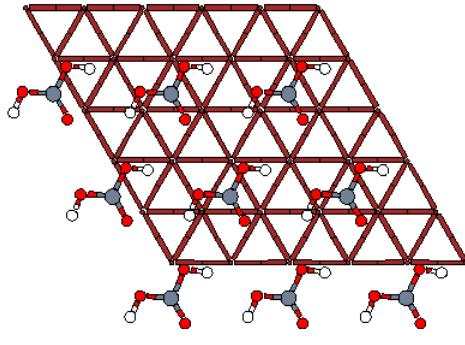
Ab initio Investigations of Comonomeric Analogues on Nickel



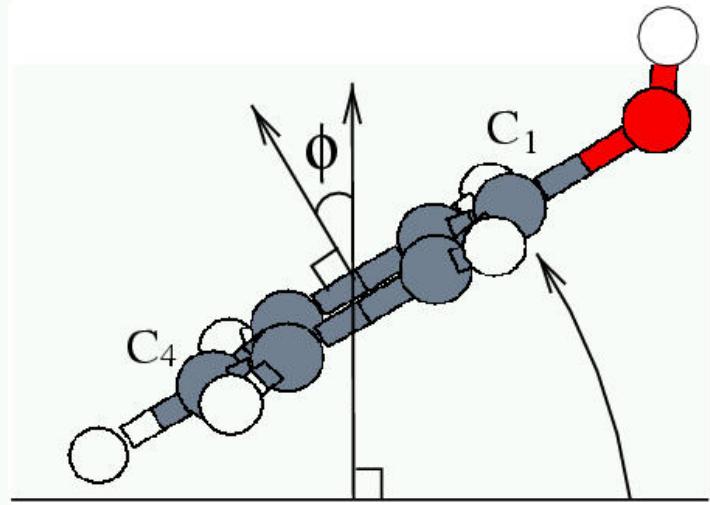
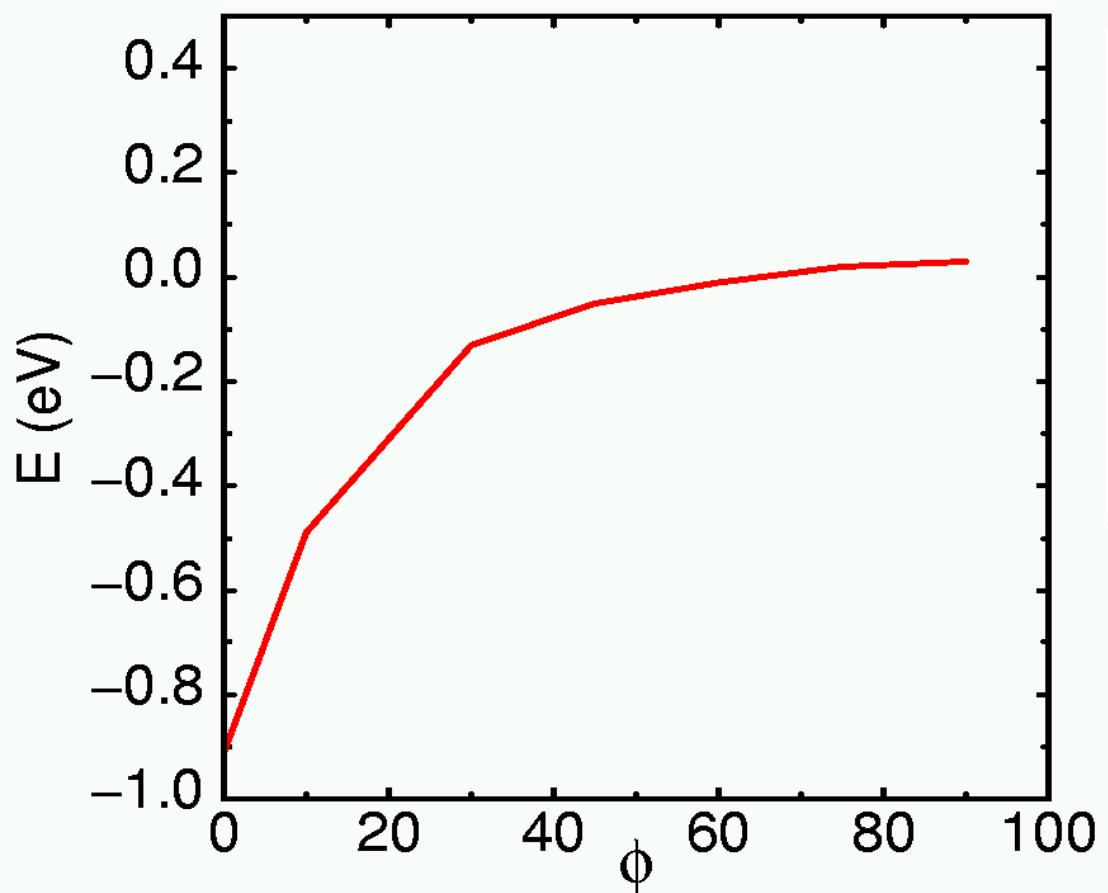
Snapshot: CPMD simulation of benzene chemisorbed on nickel (111)

(CPMD Program: M. Parrinello)

Adsorption energies of BPA-PC fragments



CPMD: Dependence of Phenol-Ni Interaction on Ring Orientation



Interaction very
sensitive to orientation!

Ab Initio MD: Conclusions

- Strong repulsion of propane and carbonic acid
 - + the strong orientational dependence
 - + short interaction range of phenol with Ni {111}

→

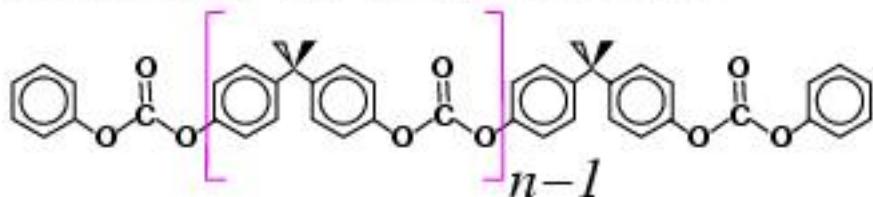
Internal phenylene comonomers in BPA-PC are sterically hindered from adsorbing on Ni {111}.

→

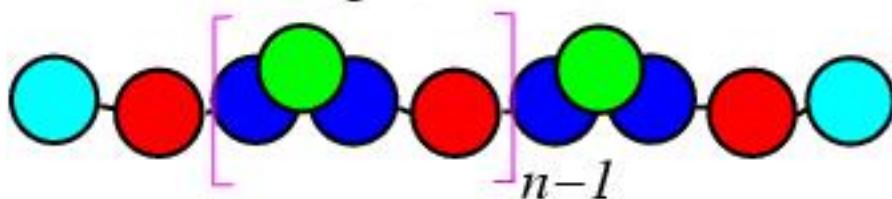
Torsional freedom in carbonate group allows for terminal phenoxy groups to adsorb

Coarse-Grained BPA-PC with End-Group Resolution (*Dual Scale MD*)

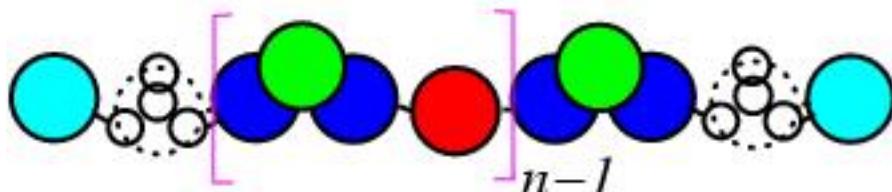
Bisphenol-A-polycarbonate



4:1 Coarse-grained



4:1 with end-group resolution

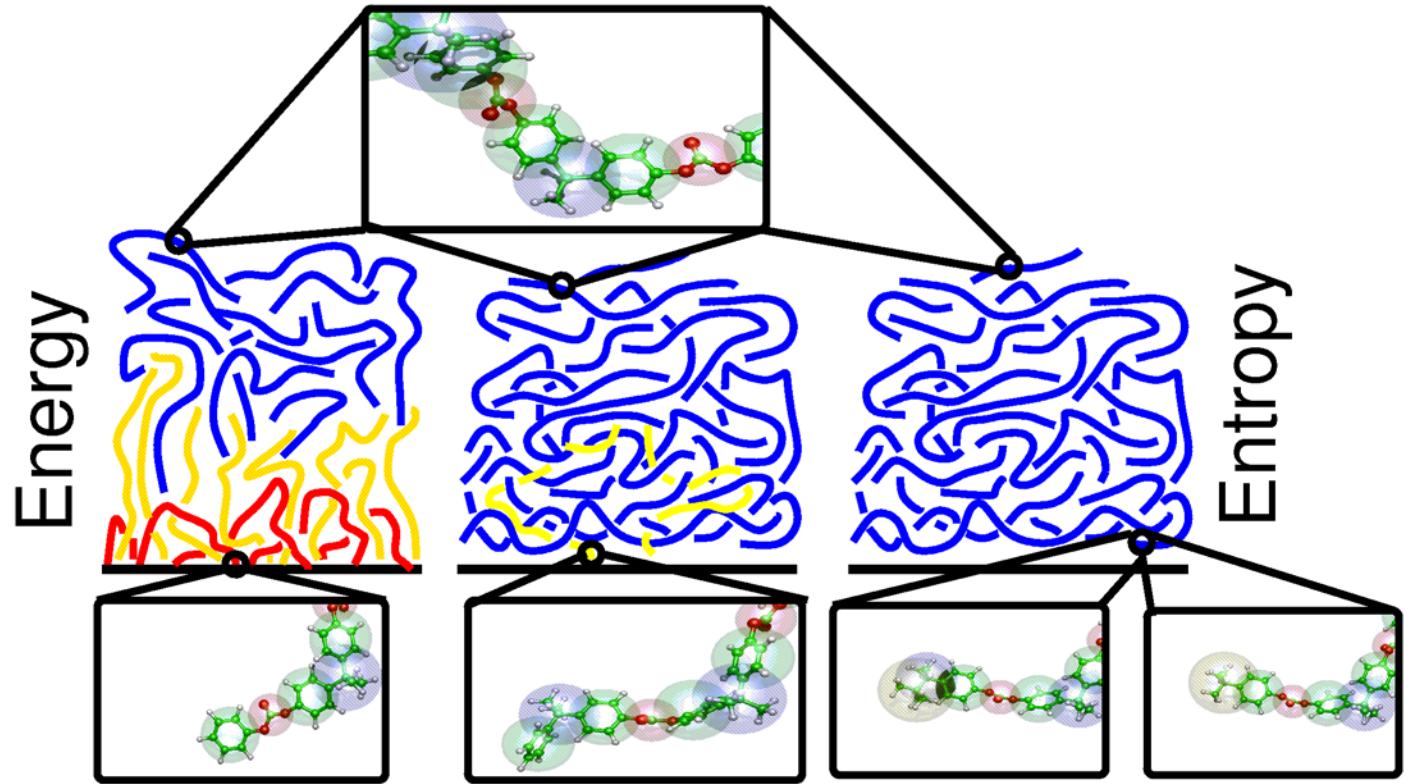


- Phenol-Ni interaction strongly dependent on C₁-C₄ phenol orientation

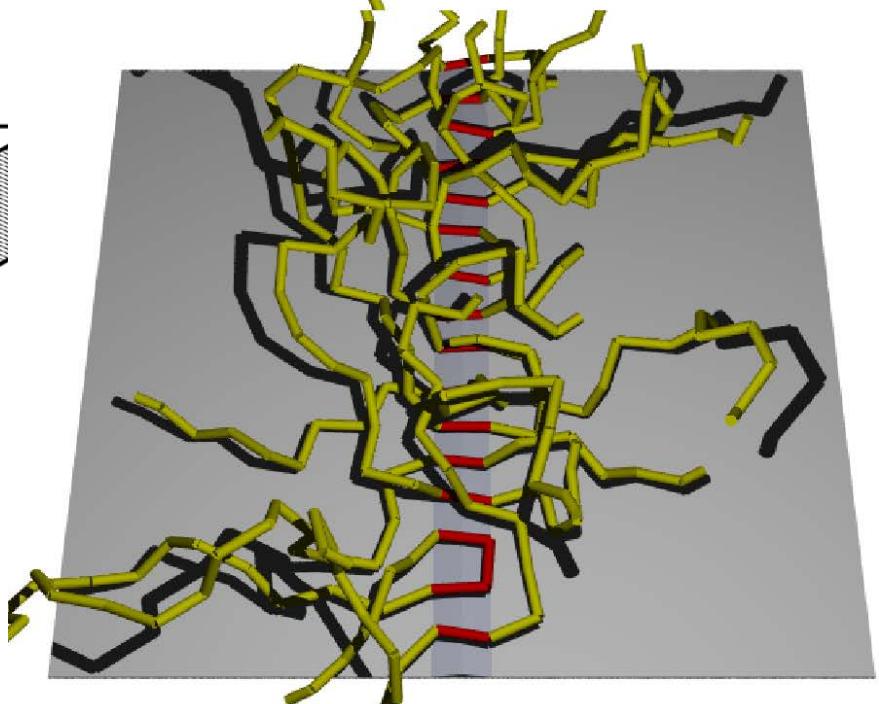
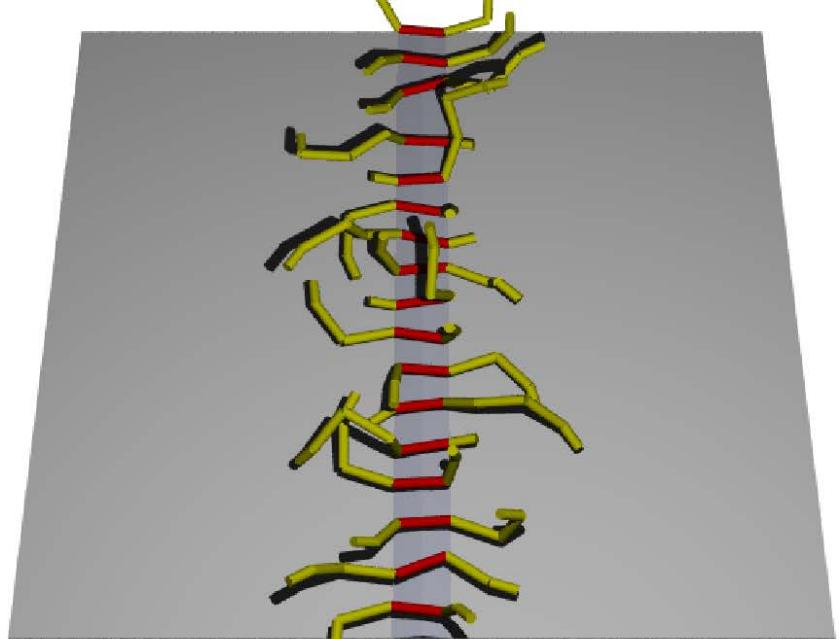
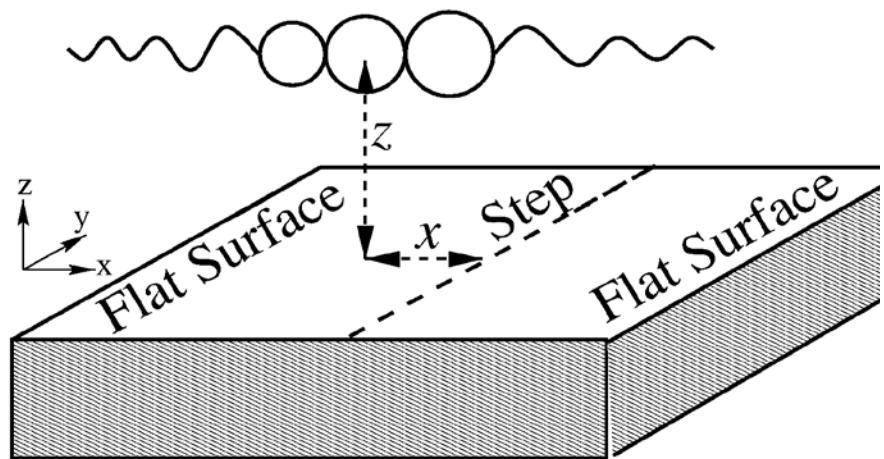
- In standard 4:1 model, phenoxy end orientation not strictly accounted for

- Resolving only the *terminal* carbonates specifies 1-4 orientation and is inexpensive

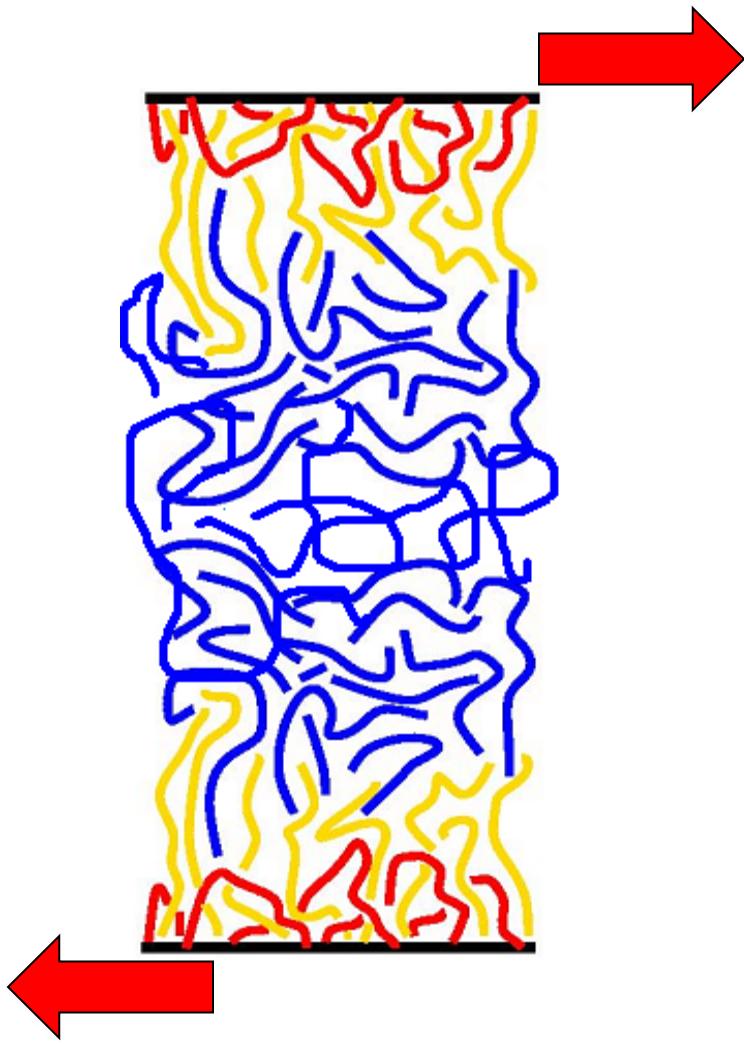
Competition Energy-Entropy Other Chain Ends



Line Defect Induced Ordering

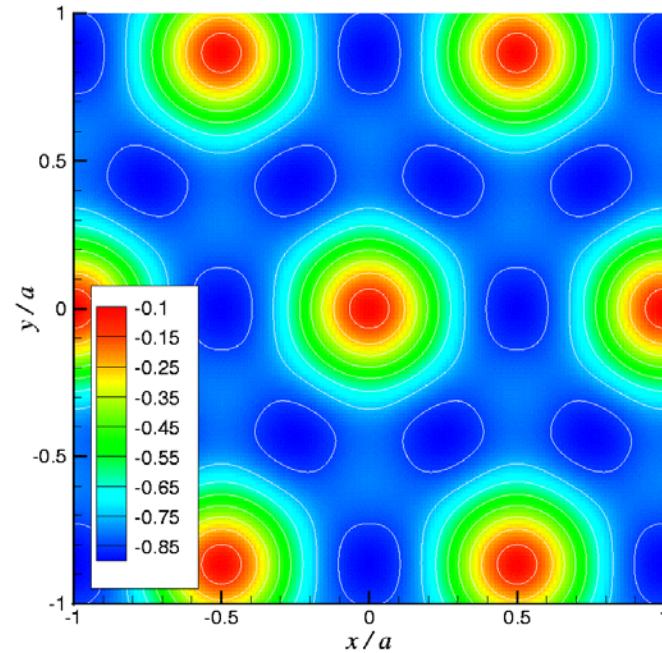


Shearing a Melt



end adsorption energy
dominated case:
phenolic chain ends

Surface Potential for Ends



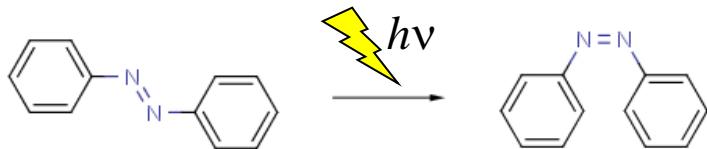


Photoswitchable Azobenzene Structure Developing Liquid

VW Stiftung, C. Peter, L. Delle Site, KK - D. Marx (Bochum)

QM scale

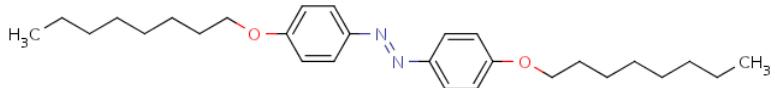
Photoisomerization of azobenzene



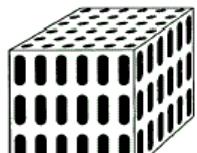
isomerization mechanism ?!?

(particularly in polymer or anisotropic environment)

trans azobenzene as a mesogen 8AB8

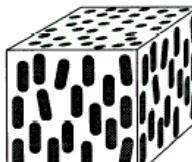


light induced phase transition



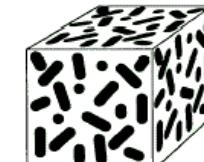
Crystal

$99\text{ }^{\circ}\text{C}$



Nematic

$112\text{ }^{\circ}\text{C}$



Isotropic

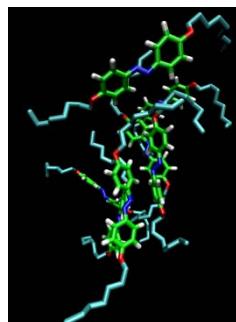
Ikeda & Tsutsumi, Science 1995

mesoscale

phenomena on wide range of time and length scales mutually influencing each other
⇒ multiscale approach

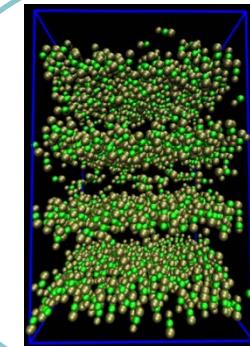
Example LC phase transitions

Disordered @ 480 K

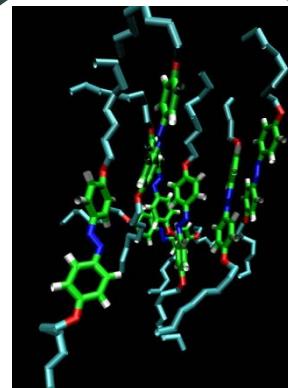


CG model should
reproduce this transition !

SmA @ 460 K



Super-cooled fluid @ 460 K



overall isotropic

Parametrise HERE !!

residual
nematic order

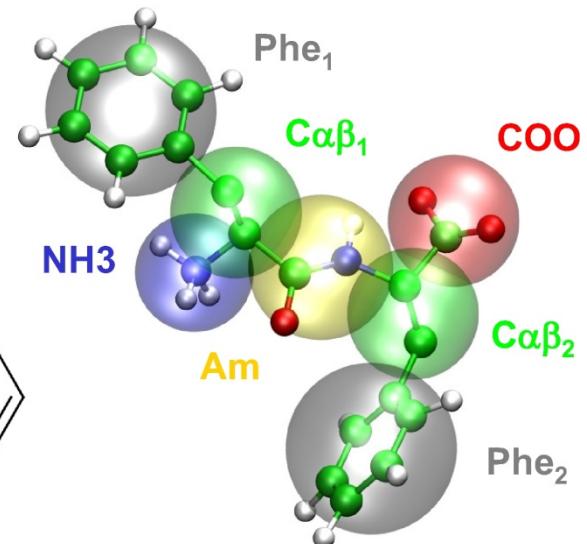
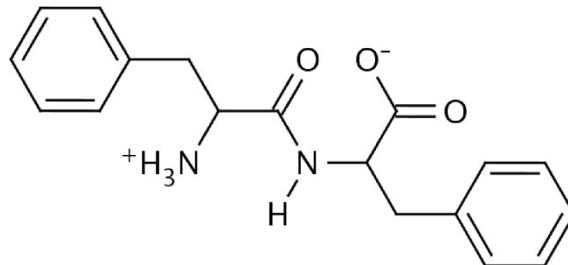
CG model should
reproduce structure

Parametrisation
impossible here:
(layering built
into potentials)

CG model for di-phenylalanine with/without explicit water

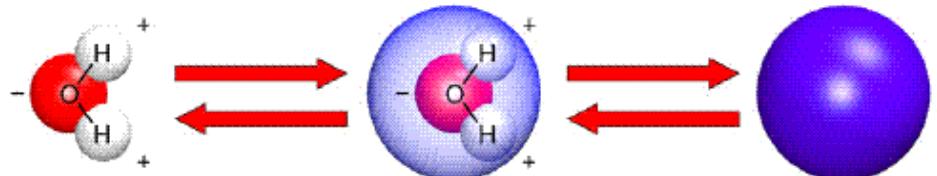


⇒ Mapping

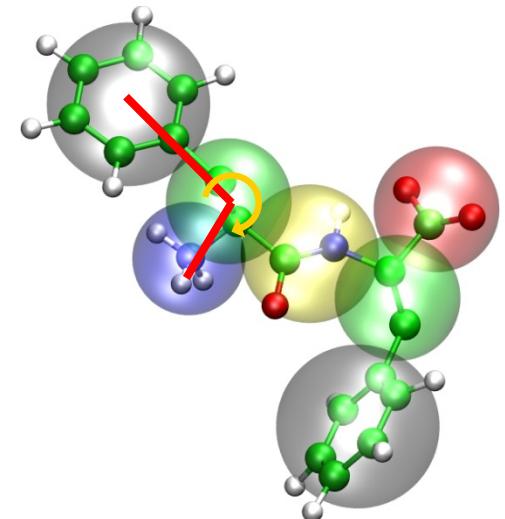
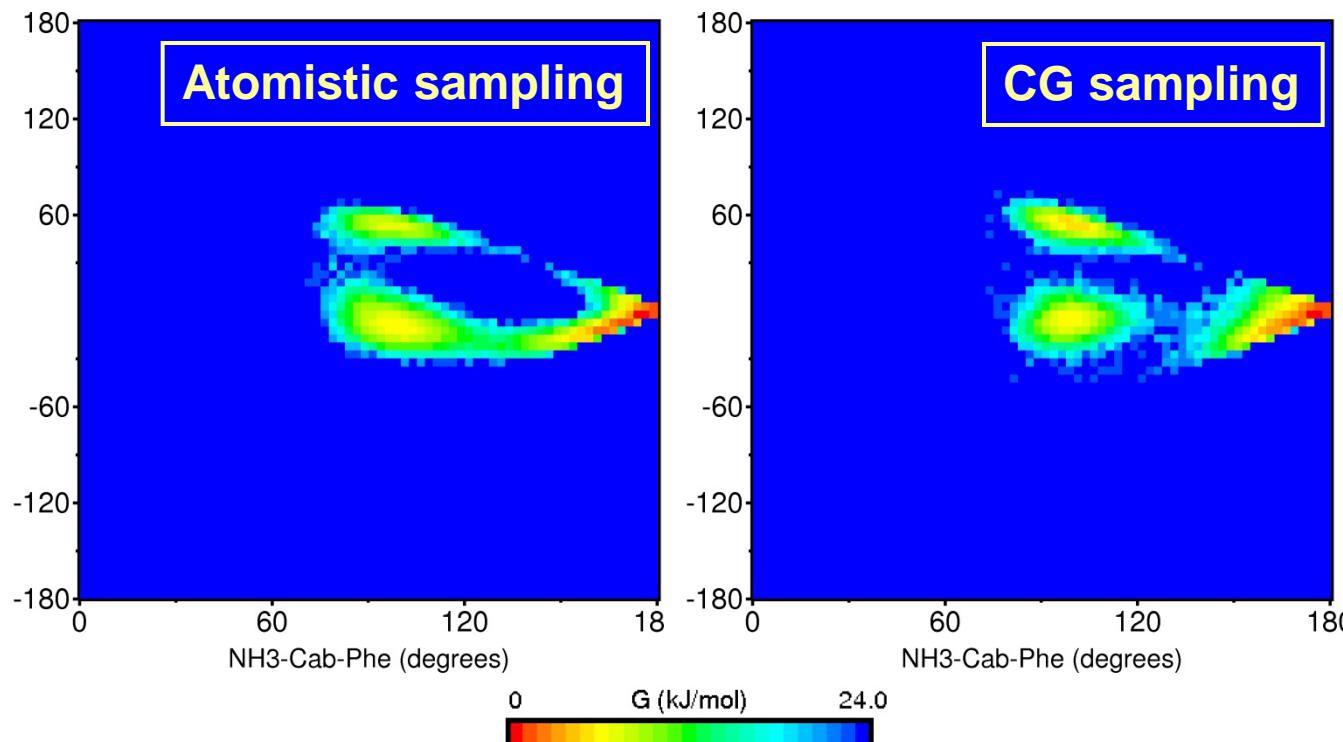


⇒ Intramolecular degrees of freedom
⇒ Nonbonded interactions

1. solvent free model
2. model w. explicit (CG) solvent representation

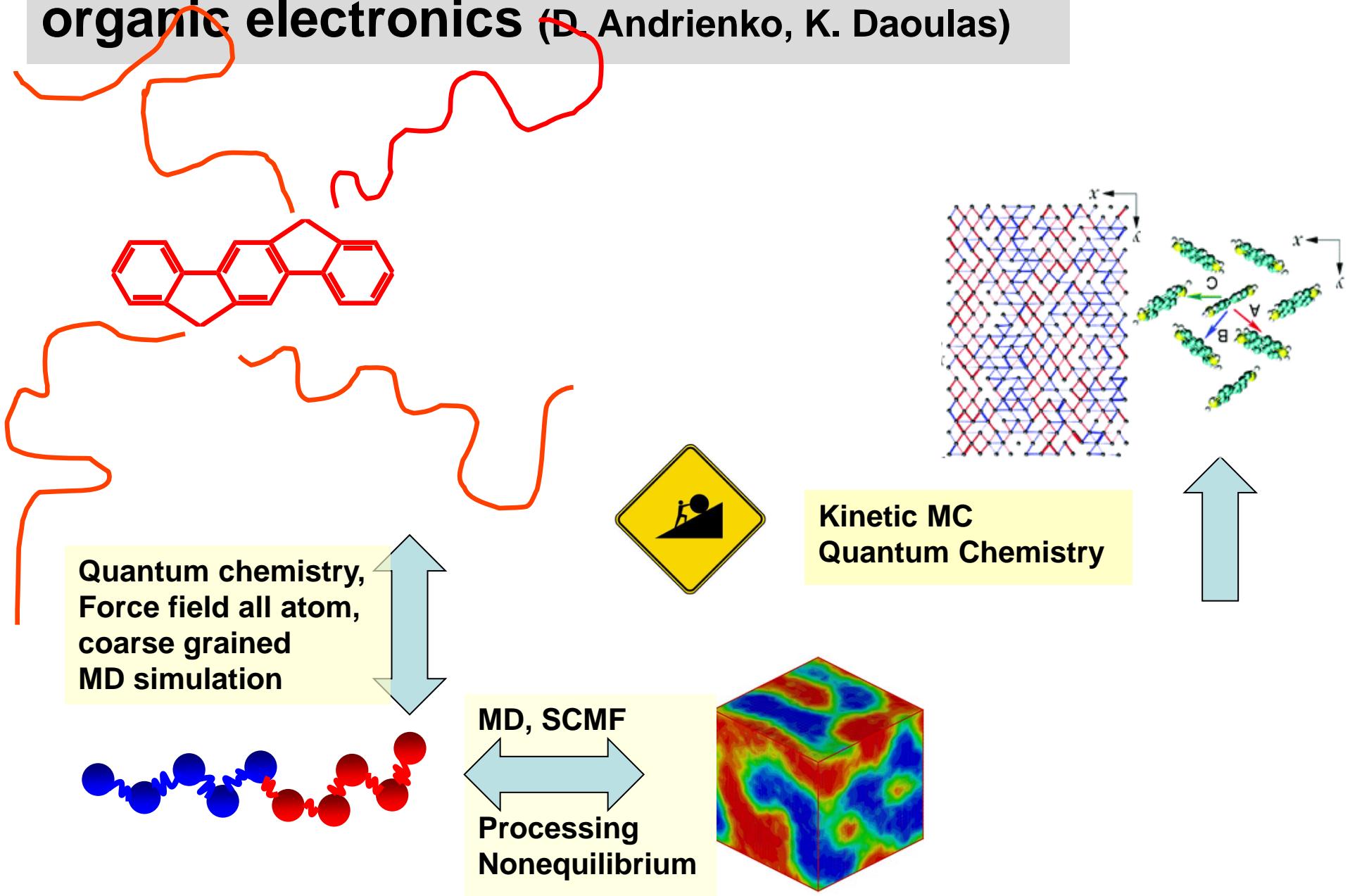


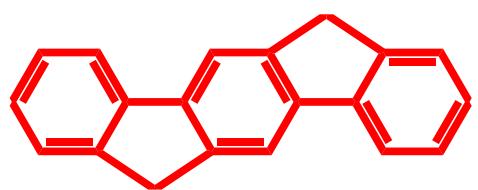
Intramolecular (bonded) degrees of freedom



⇒ Correlations between degrees of freedom are well represented in the CG description

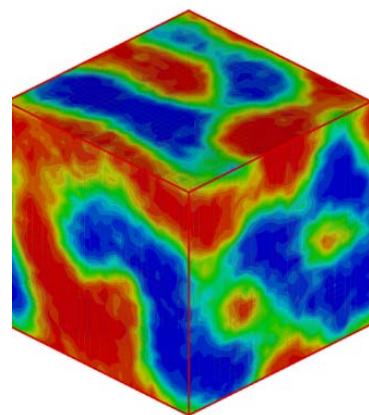
Simulation and Modeling: (ideal) Work Flow organic electronics (D. Andrienko, K. Daoulas)





MD, SCMF

**Processing
Nonequilibrium**

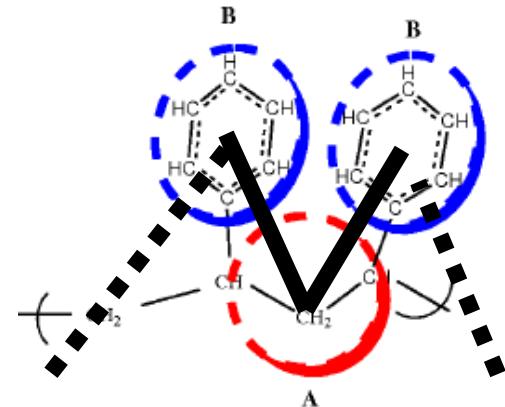


Structure Based Coarse Graining



- Dynamics
 - Melts, Rheology
 - Diffusion of Additives

Time scale, PS, BPA-PC



$$M_{\text{mon}} = 104$$

$$T=463\text{K}$$

$$M_{\text{sphere}} \approx 52$$

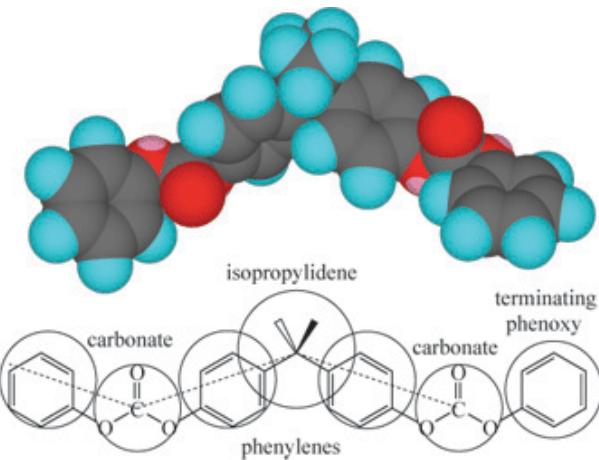
$$\sigma \approx 3 \text{ \AA}$$

$$\Rightarrow \tau_{\text{LJ}} \approx 1 \text{ ps}$$

$$N_{\text{e}}^{\text{primpath}} \approx 150$$

$$\tau_{\text{LJ}} = (M_{\text{sphere}} \sigma^2 / \epsilon)^{1/2}$$

Time scaling: $\tau^{\text{cg}}_{\text{LJ}} = \text{physical time} \times \text{constant}$



$$M_{\text{mon}} = 254$$

$$T=570\text{K}$$

$$M_{\text{sphere}} \approx 63$$

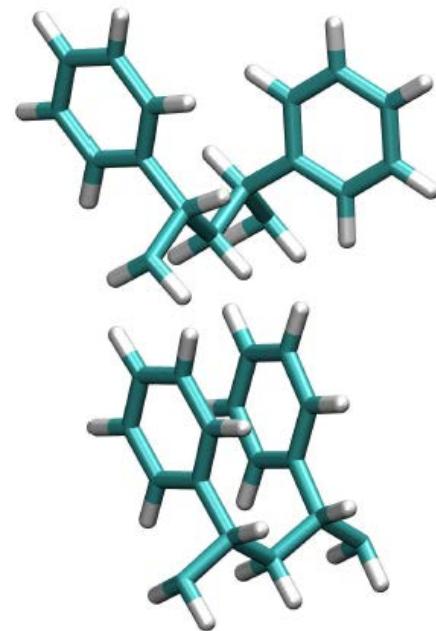
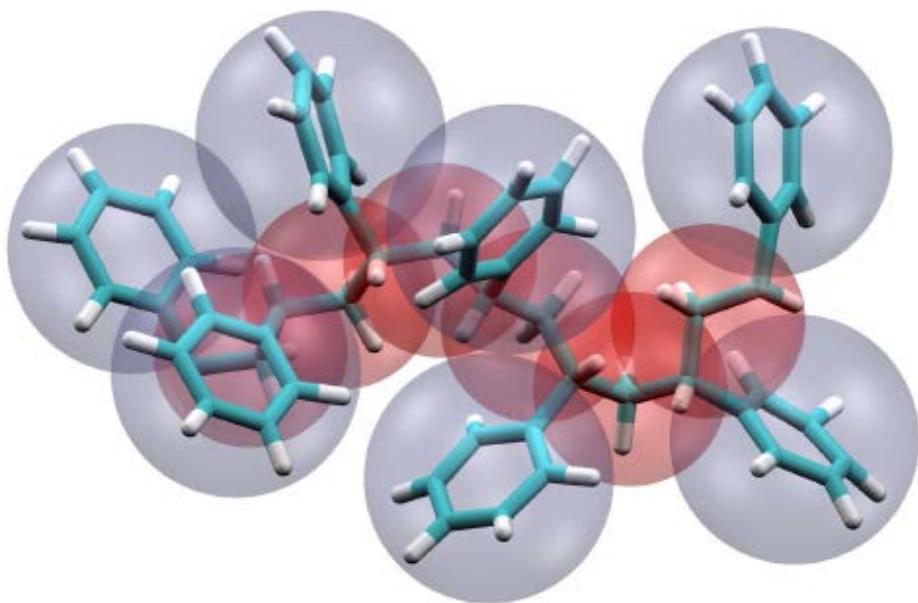
$$\sigma \approx 5 \text{ \AA}$$

$$\Rightarrow \tau_{\text{LJ}} \approx 1.7 \text{ ps}$$

$$N_{\text{e}}^{\text{primpath}} \approx 5.5$$

PS: Polystyrene

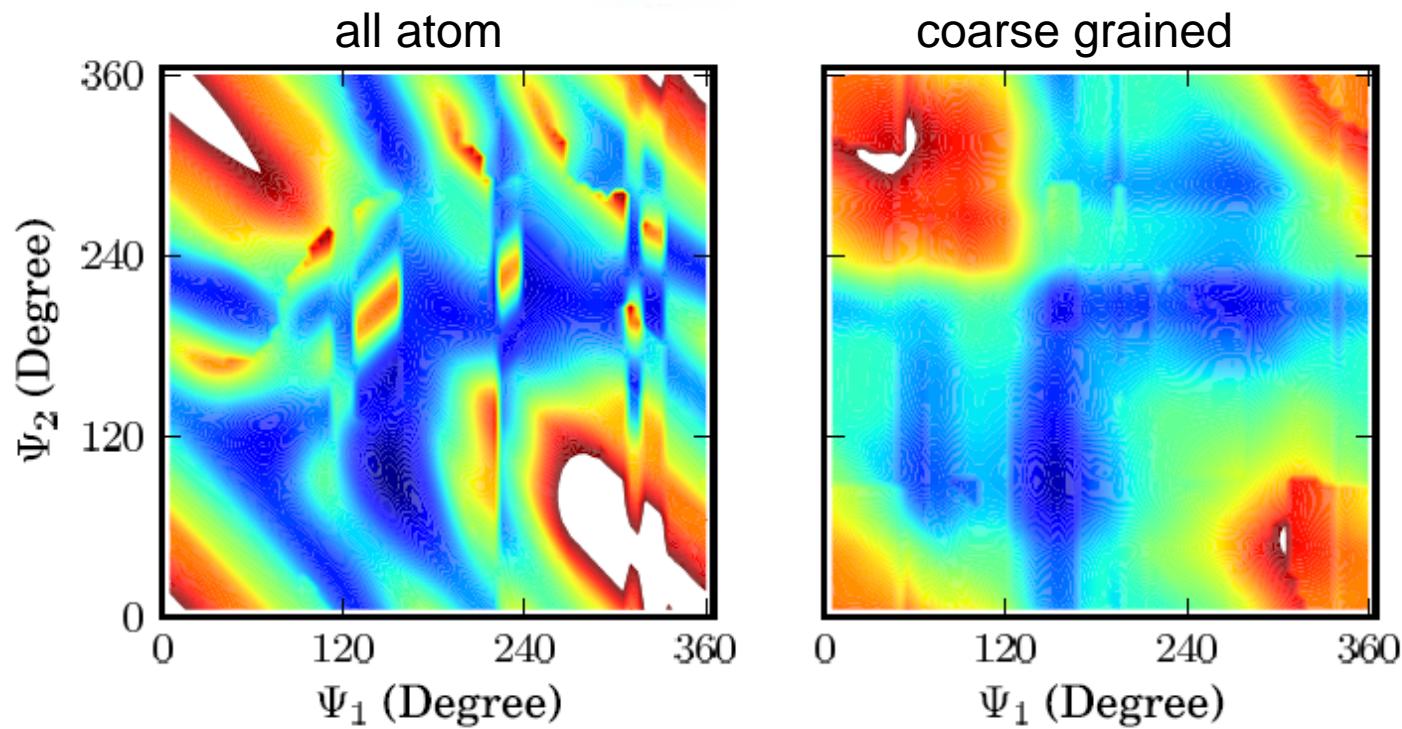
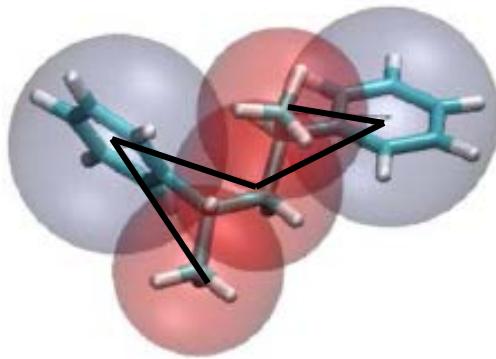
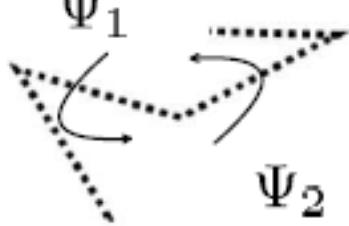
D. Fritz, V. Harmandaris, G. Floudas, N. Van der Vegt





Dynamics

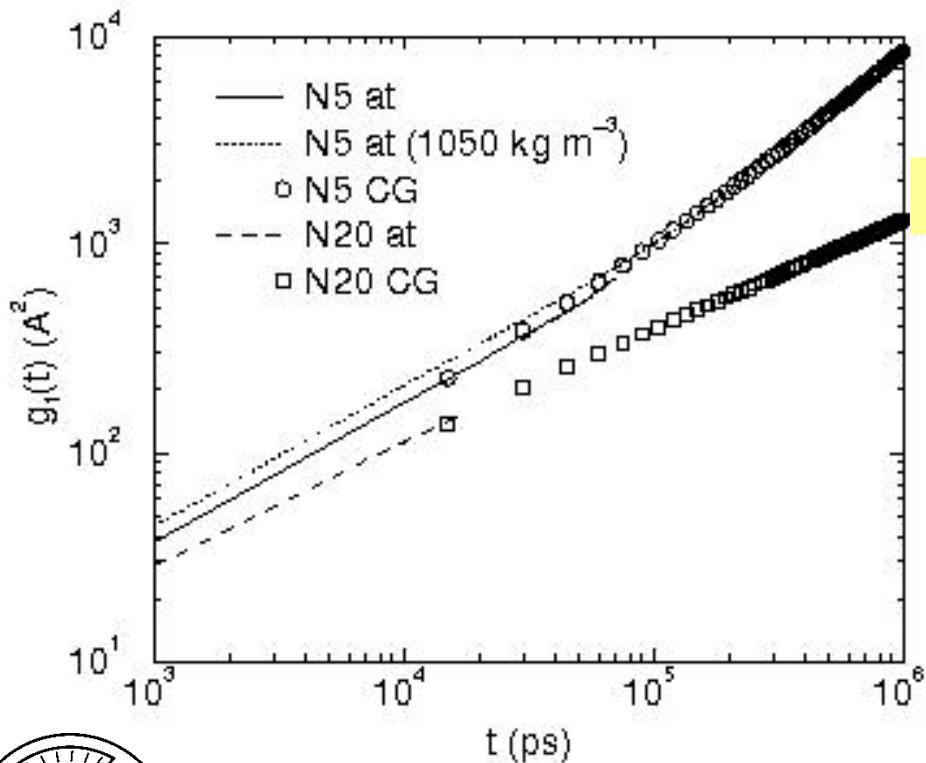
- Too complicated already for PS!



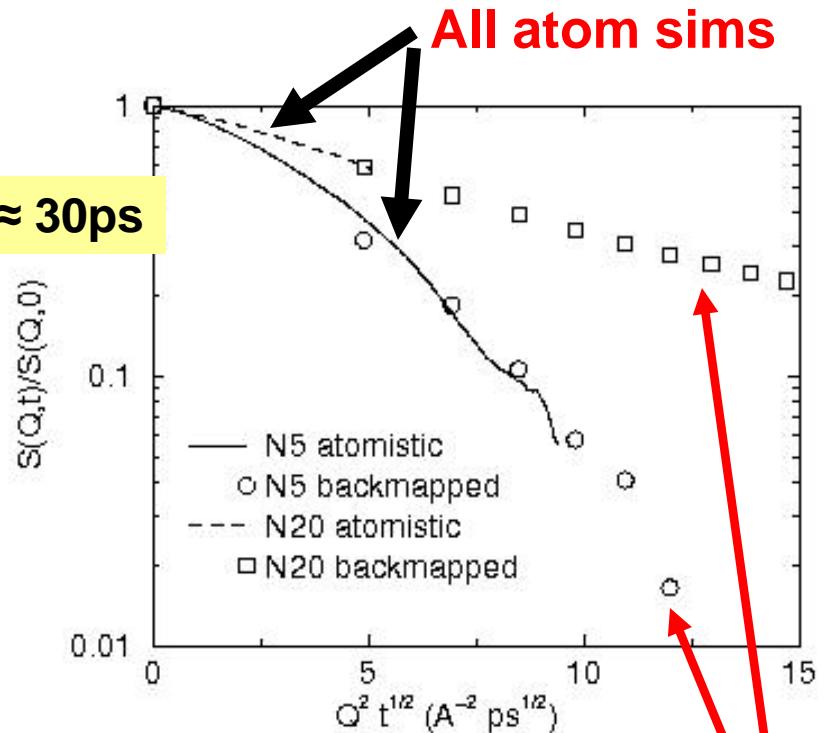
Atomistic vs cg Simulation: Time Mapping



mean square displacements
atomistic vs coarse grained



BPA-PC



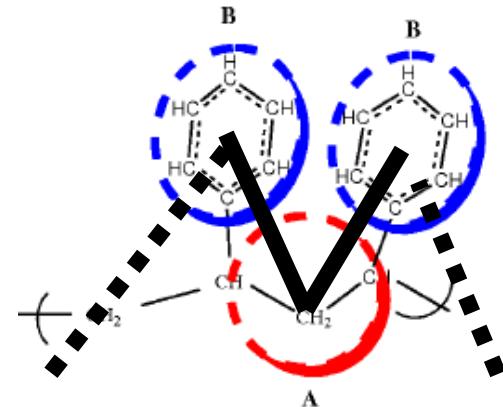
$1\tau \approx 30\text{ps}$



Leon et al, Macrom. 2005, Hess et al Soft Matter, 2006

CG sims
Back mapped

Time scale, PS, BPA-PC



$$M_{\text{mon}} = 104$$

$$T=463\text{K}$$

$$M_{\text{sphere}} \approx 52$$

$$\sigma \approx 3 \text{ \AA}$$

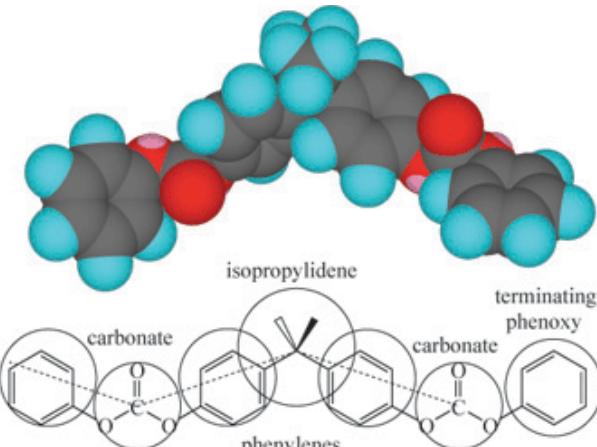
$$\implies \tau_{\text{LJ}} \approx 1 \text{ ps}$$

$$N_{\text{e}}^{\text{primpath}} \approx 150$$

$$\boxed{\tau_{\text{sim}} \approx 400 \text{ ps}}$$

$$\tau_{\text{LJ}} = (M_{\text{sphere}} \sigma^2 / \epsilon)^{1/2}$$

Time scaling: $\tau^{\text{cg}}_{\text{LJ}} = \text{physical time} \times \text{constant}$



$$M_{\text{mon}} = 254$$

$$T=570\text{K}$$

$$M_{\text{sphere}} \approx 63$$

$$\sigma \approx 5 \text{ \AA}$$

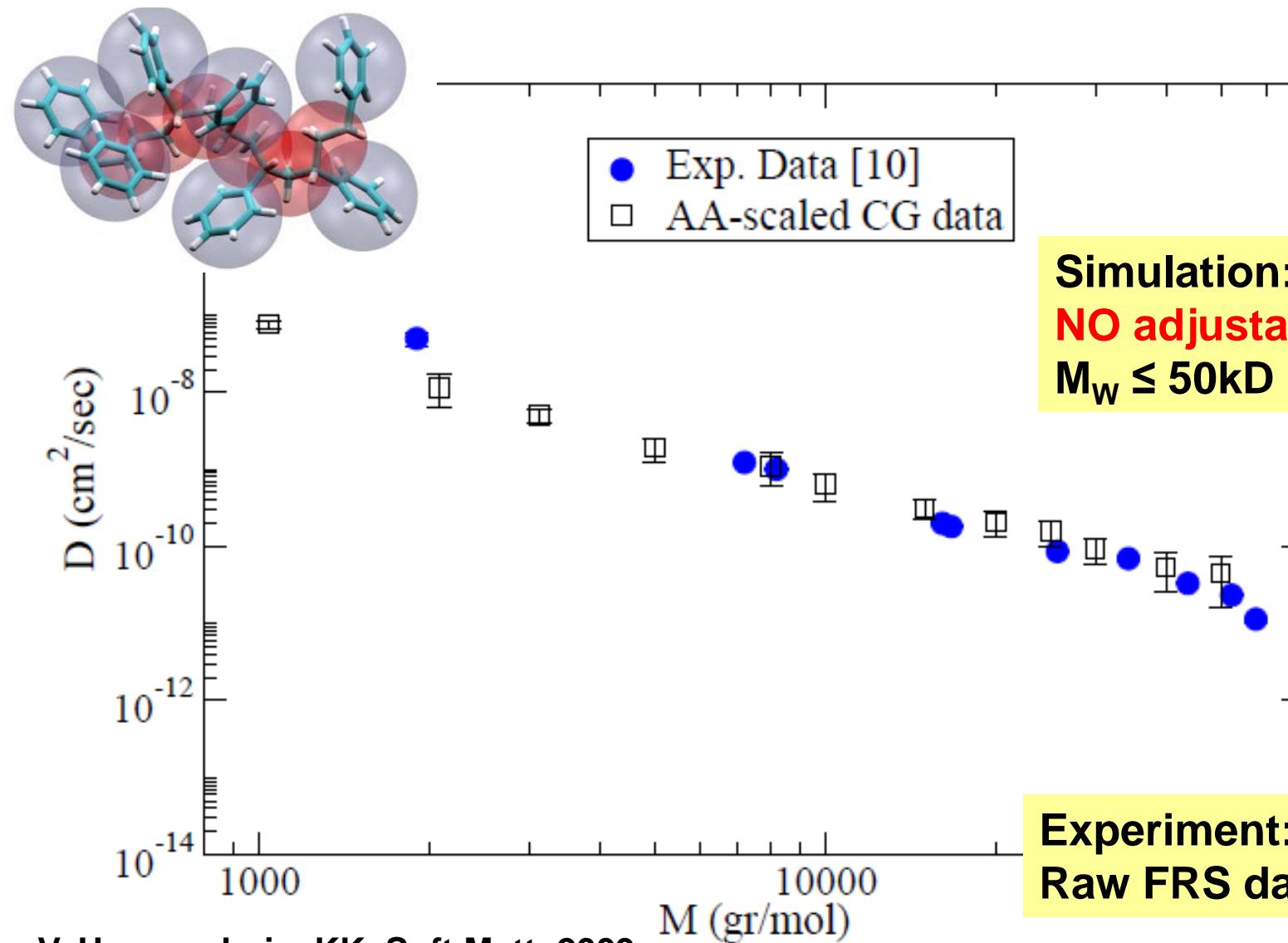
$$\implies \tau_{\text{LJ}} \approx 1.7 \text{ ps}$$

$$N_{\text{e}}^{\text{primpath}} \approx 5.5$$

$$\boxed{\tau_{\text{sim}} \approx 30 \text{ ps}}$$

Application: Diffusion Constant of PS

(two step approach AA->UA->CG)



Application: NMR data of PS

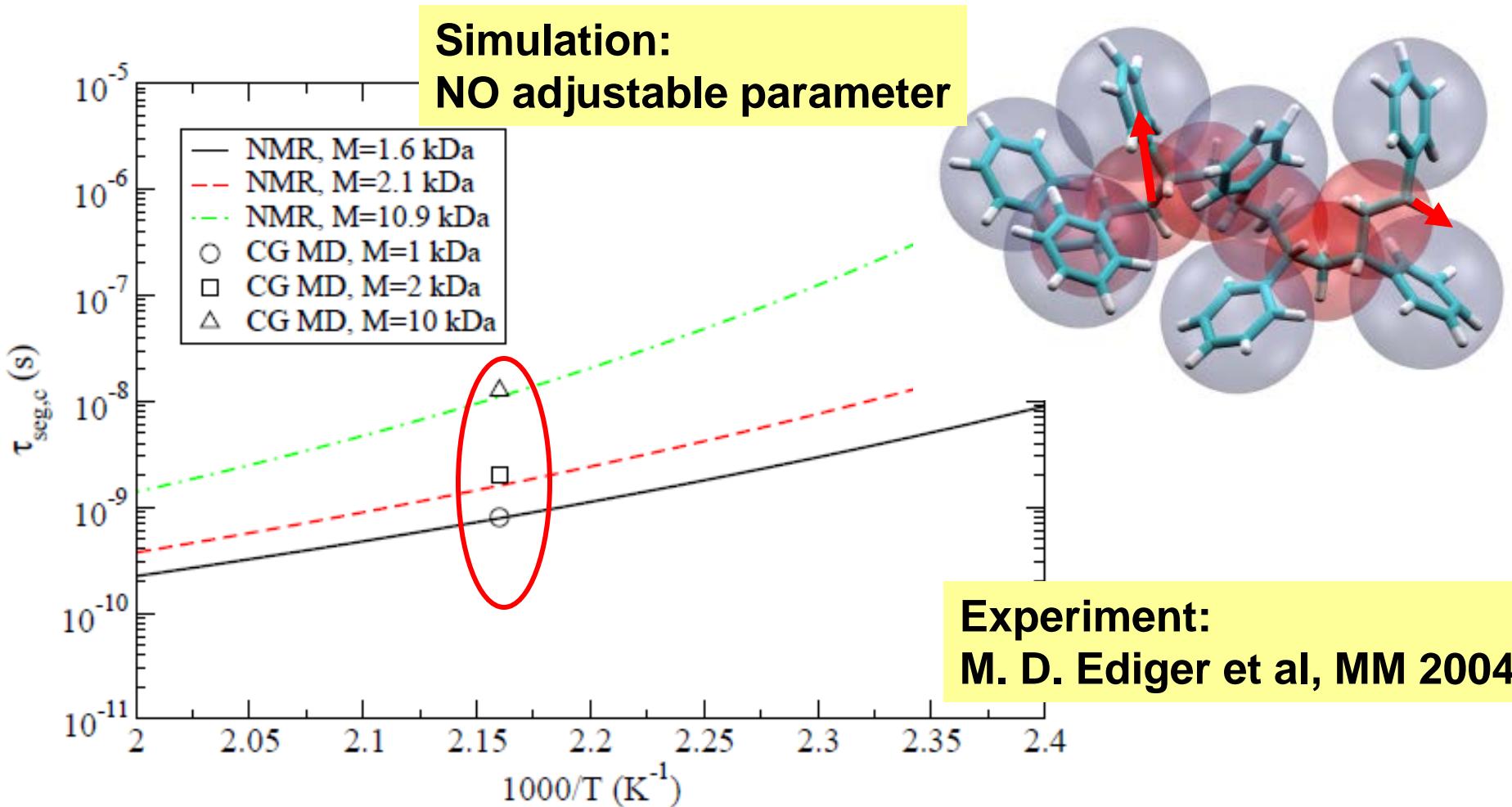


Figure 5: Segmental correlation times from NMR experiments [30] and CG MD simulations for three different molecular weights.



Conclusion preliminary

- **Soft and Nanostructured Matter**
 - Finite size effects crucial
- **Coarse Graining**
 - Bonded, nonbonded interactions, transferability
 - Inverse mapping => link to details
 - Dynamics



Force Field Bias in Protein Folding Simulations

Peter L. Freddolino,[†] Sanghyun Park,[‡] Benoît Roux,[§] and Klaus Schulten^{†*}

[†]Center for Biophysics and Computational Biology, University of Illinois at Urbana-Champaign, Urbana, Illinois; [‡]Mathematics and Computer Science Division, Argonne National Laboratory, Argonne, Illinois; and [§]Department of Biochemistry and Molecular Biology, University of Chicago, Chicago, Illinois

Force field errors can be very significant for large molecules

For proteins easily sum up to size of basin off attraction

Can CG be used to refine underlying force fields?



Conclusion preliminary

- **Soft and Nanostructured Matter**
 - Finite size effects crucial
- **Coarse Graining**
 - Bonded, nonbonded interactions, transferability
 - Inverse mapping => link to details
 - Dynamics
- **Adaptive Resolution Simulation: AdResS**
- **Conclusion/Outlook**