

6-3. Wiener - Khintchine relations / Nyquist's theorem.

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③ notation

$y(t)$: a random function of time

$$y_0(t) \equiv \begin{cases} y(t) & \text{for } -T < t < T \\ 0 & \text{otherwise} \end{cases}, \quad T \rightarrow \infty$$

↓

Fourier transform.

$$C(\omega) = \frac{1}{2\pi} \int_{-T}^T y(t') e^{-i\omega t'} dt'$$

$$y_0(t) = \int_{-\infty}^{\infty} C(\omega) e^{i\omega t} d\omega$$

If $y(t)$ is real, $C^*(\omega) = C(-\omega)$.

④ ensemble average and time average.

$$\overline{y(t)} \equiv \langle y(t) \rangle \equiv \frac{1}{N} \sum_{k=1}^N y^{(k)}(t), \quad \text{ensemble average}$$

$$\{y^{(k)}(t)\} = \frac{1}{2T} \int_{-T}^T y^{(k)}(t+t') dt' \quad \text{time average.}$$

$$\langle \{y^{(k)}(t)\} \rangle = \sum \int = \int \sum = \{ \langle y(t) \rangle \}$$

Ergodic assumption : $y^{(k)}(t)$ will pass through all the accessible values in a sufficiently long time interval.

$$\Rightarrow \{y^{(k)}(t)\} = \{y\}, \text{ independent of } k.$$

for stationary system, $\langle y(t) \rangle = \langle y \rangle$ independent of t .

$$\Rightarrow \langle \{y^{(k)}(t)\} \rangle = \{y\}$$

$$\{ \langle y(t) \rangle \} = \langle y \rangle$$

$$\Rightarrow \boxed{\{y\} = \langle y \rangle}$$

for a stationary ergodic ensemble.

Wiener - Khintchine relations

random function $y(t)$ in a stationary situation

$$K(s) \equiv \langle y(t) y(t+s) \rangle$$

$\sim$ independent of t ($\because$ stationary)

$$K(s) = \int_{-\infty}^{\infty} J(\omega) e^{i\omega s} d\omega$$

$J(\omega) \sim$ spectral density of y .

$$J(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} K(s) e^{-i\omega s} ds$$

Wiener - Khintchine relations

Since $K(s) = K(-s)$ and $K(s) \sim$ real,

$$J^*(\omega) = J(\omega) \quad \text{and} \quad J(-\omega) = J(\omega)$$

Note:

$$\langle y^2 \rangle = K(0) = \frac{\pi}{(H)} \int_{-\infty}^{\infty} |C(\omega)|^2 d\omega \quad \text{for stationary and ergodic } y(t).$$

proof.

$$K(s) = \langle y(t) y(t+s) \rangle = \{ y(t) y(t+s) \}$$

$$= \frac{1}{2\pi} \int_{-\pi}^{\pi} dt' y(t') y(t'+s)$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} dt' y_{\Theta}(t') y_{\Theta}(t'+s)$$

$$y_{\Theta} = \begin{cases} y(t') & t' \in [\Theta, \Theta] \\ 0 & \text{otherwise} \end{cases}$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} dt' \int_{-\infty}^{\infty} d\omega' C(\omega') e^{i\omega' t'} \int_{-\infty}^{\infty} d\omega C(\omega) e^{i\omega(t'+s)}$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega \int_{-\infty}^{\infty} d\omega' C(\omega') C(\omega) e^{i\omega s} [2\pi \delta(\omega' + \omega)]$$

$$= \frac{7L}{(H)} \int_{-\infty}^{\infty} d\omega C(-\omega) C(\omega) e^{i\omega s}$$

$$\because C^*(-\omega) = C(\omega)$$

$$\Rightarrow K(0) = \frac{7L}{(H)} \int_{-\infty}^{\infty} d\omega |C(\omega)|^2$$

① Nyquist theorem.

$$\overline{J_{vt}}(\omega) = \frac{2}{\pi} kTR \quad \text{for } \omega \ll \frac{1}{\tau^*}$$

$J_{vt}(\omega)$: spectral density of fluctuating voltage

R : resistance

τ^* : correlation time of fluctuating voltage

Consider an electrical circuit.

$$\Rightarrow L \frac{dI}{dt} = -RI + V_{ext}.$$

$$\Rightarrow R = \frac{1}{2kT} \int_{-\infty}^{\infty} \langle V(\omega) V(s) \rangle_s ds$$

$$\Rightarrow R = \frac{1}{2kT} [2\pi J(\omega)] = \frac{\pi}{kT} J(0)$$

$$(\because K(s) = \int_{-\infty}^{\infty} J(\omega) e^{i\omega s} d\omega$$

$$\rightarrow K(0) = \int_{-\infty}^{\infty} J(\omega) d\omega = \int_0^{\infty} J_+(\omega) d\omega.$$

$$\underline{J_+(\omega) \equiv 2 J(\omega)} \quad (\because J(-\omega) = J(\omega))$$

↪ spectral density for positive frequency.

$$\Rightarrow R = \frac{\pi}{2kT} J_+(0)$$

$$\Rightarrow J_+(0) = \frac{2}{\pi} kTR.$$

Since $K(s)$ has a finite value for $s \approx \tau^*$,

$$\boxed{J_+(\omega) = \frac{2}{\pi} kTR \quad \text{for } \omega \ll \frac{1}{\tau^*}}$$

($\because J_+(\omega) = J_+(\omega)$)

Nyquist's theorem.

~ It represents a connection between fluctuation and dissipation.