

Vacuum fluctuation of inflaton $\delta\phi$

↓
classical perturbation $\delta\phi$

setting the initial condition → density perturbation $\delta\rho$
(curvature perturbation)
(scalar) (vector) (tensor)

Reheating

primordial values.

RD

(R-M Equality)

Evolution of perturbation

MD

Recombination

$\frac{\delta T}{T}$

CMBA

structure formation. $\frac{\delta\rho}{\rho}$

Observation.

graviton

δh

↓
 δh

Cosmological perturbation theory.

- The unperturbed universe

$$ds^2 = a^2(\tau) \left(\underbrace{-d\tau^2}_{\substack{\uparrow \\ \text{conformal time}}} + \underbrace{dx^2 + dy^2 + dz^2}_{\text{flat}} \right)$$

$$T^{\mu\nu} = (\rho + p) u^\mu u^\nu + p g^{\mu\nu}, \quad u^\mu = (-1, 0, 0, 0) \\ \text{rest in the comoving frame}$$

- Perturbed metric

- No uniquely preferred coordinate system
- Requirement: reduced to the unperturbed metric in the limit of zero perturbation

Most general first-order perturbation

$$ds^2 = a^2(\tau) \left\{ \underbrace{-(1+2A)d\tau^2}_{\text{lapse function}} - \underbrace{B_i d\tau dx^i}_{\text{shift function}} + \underbrace{[(1+2D)\delta_{ij} + 2E_{ij}] dx^i dx^j}_{\text{traceless}} \right\}$$

$A, B_i, D, E_{ij} : 10$

- perturbed stress energy tensor

Global coordinate system $x^\mu = (\tau, x^i)$

At each spacetime point, a locally orthonormal coordinate system (t, r_i) with the following property

$$dt = a(1+A)d\tau$$

$$dr_i = a dx^i$$

Let $v^i = \frac{dr_i}{dt}$: the fluid velocity in the locally orthonormal frame.

Fluid velocity in the global coordinate

$$u^\mu = \frac{dx^\mu}{d\tau} \quad u^0 = \\ = \frac{dx^\mu}{a d\tau} \quad u^i =$$

density perturbation

$$T^0_0 = -(\rho + \delta\rho)$$

fluid velocity in locally orthonormal frame.

$$T^0_i = (\rho + p)(v_i - B_i)$$

$$T^i_0 = -(\rho + p)v^i$$

anisotropic stress (traceless)

$$T^i_j = (\rho + \delta\rho)\delta^i_j + \Sigma^i_j$$

pressure perturbation

$$* \Pi_{ij} \equiv \frac{\Sigma_{ij}}{p}$$

$\delta\rho, \delta p, v^i, \Sigma^i_j : 10$

- Scalar-vector-tensor decomposition

$$\delta G_{\mu\nu} = \kappa^2 \delta T_{\mu\nu}$$

metric perturbation stress-energy perturbation

$\Rightarrow$ evolution equation
constraint equation

o U_i decomposition

$$U_i = U_i^S + U_i^V \quad (\text{scalar} + \text{vector})$$

$\partial_i S$ (derivative of some scalar function)

In Fourier space

$$U_i^S = -\frac{ik_i}{k} V, \quad k_i U_i^V = 0$$

parallel to $\vec{k}$ perpendicular to $\vec{k}$

o Π_{ij} decomposition

$$\Pi_{ij} = \underbrace{\Pi_{ij}^S}_{\partial_i \partial_j S} + \underbrace{\Pi_{ij}^V}_{\partial_i V_j} + \Pi_{ij}^T$$

traceless

$$\Pi_{ij}^S = \left(-\frac{k_i k_j}{k^2} - \frac{1}{3} \delta_{ij} \right) \Pi$$

scalar ①

$$\Pi_{ij}^V = \frac{-i}{2k} (k_i \Pi_j - k_j \Pi_i) \text{ with } k_i \Pi_i = 0.$$

vector ②

$$k_i \Pi_{ij}^T = 0. \quad ③$$

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$$\vec{k} = (0, 0, k) : \Pi_{ij}^S = \begin{pmatrix} \Pi & 0 & 0 \\ 0 & \Pi & 0 \\ 0 & 0 & -2\Pi \end{pmatrix}, \Pi_{ij}^V = -\frac{i}{2} \begin{pmatrix} 0 & 0 & \Pi_1 \\ 0 & 0 & \Pi_2 \\ \Pi_1 & \Pi_2 & 0 \end{pmatrix}, \Pi_{ij}^T = \begin{pmatrix} \Pi^x \Pi^+ & 0 \\ \Pi^+ \Pi^x & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

• B_i Decomposition

$$B_i = B_i^S + B_i^V : B_i^S = -\frac{ik_i}{k} \textcircled{B}, \quad k_i \textcircled{B_i^V} = 0$$

• E_{ij} decomposition

$$E_{ij} = E_{ij}^S + E_{ij}^V + E_{ij}^T ; \quad E_{ij}^S = \left(-\frac{k_i k_j}{k^2} + \frac{1}{3} \delta_{ij}\right) \textcircled{E}$$

$$E_{ij}^V = -\frac{i}{2k} (k_i \textcircled{E_j} + k_j \textcircled{E_i})$$

$$\textcircled{B_i E_{ij}^T} = 0.$$

* Equations break down into three independent sets.

	Metric	Stress-energy
Scalar	A, B, D, E	$\delta\phi, \delta p, V, \Pi$
Vector	B_i^V, E_i	v_i^V, Π_i
tensor	$E_{ij}^T \equiv a^{-2} h_{ij}^T$	Π_{ij}^T

* Physical significance

scalar perturbation - generated by the vacuum fluctuation of the inflaton field
"Density perturbation" included.

vector perturbation - relativistic generalization of purely rotational fluid flow.

tensor perturbation - gravitational waves
(generated during inflation)

- Gauge transformations

$$\tilde{x}^0 = x^0 + \delta\tau(x^\mu), \quad \tilde{x}^i = x^i + \delta x^i(x^\mu)$$

tensor : gauge invariant

vector : not interested

scalar :

Metric

$$\tilde{A} = A - \dot{\delta\tau} - aH\delta\tau$$

$$\tilde{B} = B + (\dot{\delta x}) + k\delta\tau$$

$$\tilde{D} = D - \frac{k}{3} \delta x - aH\delta\tau$$

$$\tilde{E} = E + k\delta x$$

stress-energy

$$\tilde{V} = V + (\dot{\delta x})$$

$$\tilde{\delta} = \delta + 3(1+w)aH\delta\tau$$

$$\tilde{\delta p} = \delta p - \dot{p}\delta\tau$$

$$\tilde{\Pi} = \Pi$$

- spatial curvature for scalar perturbations

spatial curvature of the slices of fixed τ

$$R^{(3)} = 4 \frac{k^2}{a^2} (D + \frac{1}{3} E)$$

under the gauge transformation

$$(\tilde{D} + \frac{1}{3} \tilde{E}) = (D + \frac{1}{3} E) - aH \delta\tau \quad \text{--- indep. of } \delta x$$

curvature perturbation

$$R = D + \frac{1}{3} E \quad (\text{in comoving slice})$$

$$R^{(3)} = 4 \frac{k^2}{a^2} R$$

- Evolution of the perturbations

o Tensor perturbations

$$ds^2 = a^2(\tau) [-d\tau^2 + (d_{ij} + 2\underbrace{E_{ij}^T}_{\text{tensor}}) dx^i dx^j]$$

$$\ddot{E}_{ij}^T + 2aH \dot{E}_{ij}^T + k^2 E_{ij}^T = 8\pi G a^2 p \Pi_{ij}^T$$

After horizon entry ($k \gtrsim aH$): oscillates $\rightarrow$ gravitational wave

$$h_{ij} \equiv a E_{ij}^T$$

$$\frac{\partial^2 h_{ij}}{\partial t^2} + 3H \frac{\partial h_{ij}}{\partial t} + \left(\frac{k}{a}\right)^2 h_{ij} = 0 \quad \text{if } \underline{\Pi_{ij}^T = 0}$$

same as massless scalar field.

o Scalar perturbation

Choice of gauge $\leftarrow \begin{matrix} D_\mu T^{\mu\nu} = 0 \\ \mu=0 \end{matrix}$ $\begin{matrix} \mu=1 \\ \text{Euler equation} \end{matrix}$
the use of conservation eq & Euler equation

conformal Newtonian gauge

$$ds^2 = a^2(\tau) [-(1+2\Psi)d\tau^2 + (1-2\Phi)\delta_{ij}dx^i dx^j]$$

Evolution equations

$$\dot{\delta}_N = -(1+w)(kV_N - 3\dot{\Phi}) + 3aHw\delta_N - 3aH\frac{\delta p_N}{\rho}$$

$$\dot{V}_N = -aH(1-3w)V_N - \frac{\dot{w}}{1+w}V_N + k\frac{\delta p_N}{\rho+p} - \frac{2}{3}k\frac{w}{1+w}\Pi + k\Phi$$

constraint equation

$$k^2\Phi = -4\pi G a^2 \rho \left[\delta_N + 3\frac{aH}{k}(1+w)V_N \right]$$

$$k^2(\Psi - \Phi) = -8\pi G a^2 p \Pi$$

Total matter gauge : make a gauge transformation of
from conformal Newtonian gauge

$$\delta\chi = 0, \quad \delta\tau = \frac{V}{k}$$

stress-energy tensor perturbation

$$V = V_N$$

$$\delta = \delta_N + 3\frac{aH}{k}(1+w)V_N$$

$$\delta p = \delta p_N - \frac{\dot{p}V_N}{k}$$

Rewrite the equation (Φ, Ψ in conformal Newtonian gauge
 $V, \delta, \delta p$ in total matter gauge

$$\dot{\delta} - 3w a H \delta = -(1+w)kV - 2aHw\Pi$$

$$\dot{V} + aH V = k\frac{\delta p}{\rho+p} - \frac{2}{3}\frac{w}{1+w}k\Pi + k\Phi$$

$$k^2\Phi = -4\pi G a^2 \rho \delta$$

$$k^2(\Psi - \Phi) = -8\pi G a^2 p \Pi$$

- Four fluid : CDM, B, γ , ν

- Matter equation

o CDM : Non-interacting, $\omega=0$.

$$\dot{\delta}_c^N = -kV_c + 3\dot{\Phi}$$

$$\dot{V}_c = -aH V_c + k\Psi$$

o Baryons

pressure, baryon (electron)-photon interaction

equation of state $\frac{dP_b}{d\rho_b} = C_s^2 \rightarrow \delta P_b = C_s^2 \delta \rho_b$

$C_s^2 = ?$ $P_b = nm$, $P_b = nT_b$, $n \propto a^{-3}$

$$d\rho_b = -3nm \frac{da}{a} = -3nm d \ln a$$

$$dP_b = -3nT_b d \ln a + nT_b d \ln T_b$$

$$C_s^2 = \frac{dP_b}{d\rho_b} = \frac{T_b}{m} \left(1 - \frac{1}{3} \frac{d \ln T_b}{d \ln a} \right)$$

baryon temperature evolution

$$\frac{dT_b}{dt} = -2aHT_b + \frac{2}{3} \frac{\mu}{m_e} \frac{P_r}{P_b} a n_e \sigma_T (T_r - T_b)$$

$$\dot{\delta}_b^N = -kV_b + 3\dot{\Phi}$$

$$\dot{V}_b = -aH V_b + k\Psi + \underbrace{k C_s^2 \delta_b^N}_{\text{pressure}} + \underbrace{\frac{4P_r}{3P_b} a n_e \sigma_T (V_r - V_b)}_{\text{Coupling to photons}}$$

Free electron density

~~coupling~~ Coupling to photons

Gas dynamics in flat spacetime

- Distribution function

$$dN = g_i f(t, \vec{r}, \vec{p}) d^3\vec{r} \frac{d^3\vec{p}}{(2\pi)^3}$$

stress-energy tensor

$$T^{\mu\nu} = g_i \int f p^\mu p^\nu \frac{d^3\vec{p}}{(2\pi)^3 E}$$